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Dispersion Interactions from the Exchange-Hole Dipole Moment.

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XDM	I, basics					

$$E_{disp} = -\frac{1}{2} \sum_{ij} \frac{C_6 f_6(R_{ij})}{R_{ij}^6} + \left[\frac{C_8 f_8(R_{ij})}{R_{ij}^8} + \frac{C_{10} f_6(R_{ij})}{R_{ij}^{10}} + \dots \right]$$

comes from perturbation theory:

$$E^{(2)} = \frac{\langle \hat{V}_{\rm int}^2 \rangle}{\Delta E}$$



where:

- Interaction between neutral fragments (classical electrostatic interactions already captured at semilocal level).
- Asymptotic expression.

Johnson, E. R. and Becke, A. D., J. Chem. Phys. **123** (2005) 024101 Becke, A. D. and Johnson, E. R., J. Chem. Phys. **122** (2005) 154104, **127** (2007) 154108

XDM

The e	exchange	hole model				
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The exchange hole:

$$h_{x\sigma}(\mathbf{1},\mathbf{2}) = -rac{|
ho_{1\sigma}(\mathbf{1},\mathbf{2})|^2}{
ho_{1\sigma}(\mathbf{1})}$$

- Probability of exclusion of same-spin electron.
- On-top depth condition: $h_{x\sigma}(\mathbf{1},\mathbf{1}) = -\rho_{1\sigma}(\mathbf{1})$
- Normalization: $\int h_{x\sigma}(1,2)d2 = -1$ for all 1.
- $\rho_{1\sigma}(\mathbf{1}, \mathbf{2}) = \sum_{i}^{\sigma} \psi_{i}^{*}(\mathbf{1}) \psi_{i}(\mathbf{2})$



The exchange-hole model



- Model for dispersion: interaction of electron-hole dipoles.
- Dipole: $d_{x\sigma}(\mathbf{r}) = \int \mathbf{r}' h_{x\sigma}(\mathbf{r}, \mathbf{r}') d\mathbf{r}' \mathbf{r}$
- Assumption: dipole oriented to nearest nucleus.

$$\langle M_l^2 \rangle_i = \sum_{\sigma} \int \omega_i(\mathbf{r}) \rho_{\sigma}(\mathbf{r}) [r_i^l - (r_i - d_{X\sigma})^l]^2 d\mathbf{r} .$$

Becke, A. D. and Johnson, E. R., J. Chem. Phys. 122 (2005) 154104

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The Becke-Roussel model of exchange-hole

• Becke-Roussel model of *h_x*. (PRA **39** (1989) 3761)

Parameters (A,a,b) obtained:

- Normalization
- Value at reference point.
- Curvature at reference point (reqs. kinetic energy density).

Advantages:

- Semilocal model of the dipole $(d_x = b)$.
- **2** XDM dispersion model \longrightarrow meta-GGA.
- Better performance than exact hole (HF) version in molecules.

Becke, A. D. and Roussel, M. R., Phys. Rev. A 39 (1989) 3761



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The XDM equations: interaction coefficients

Multipole moments

$$\langle M_l^2 \rangle_i = \sum_{\sigma} \int \omega_i(\mathbf{r}) \rho_{\sigma}(\mathbf{r}) [r_i^l - (r_i - d_{X\sigma})^l]^2 d\mathbf{r}$$

use Hirshfeld atomic partition:

$$\omega_i(\mathbf{r}) = rac{
ho_i^{
m at}(\mathbf{r})}{\sum_j
ho_j^{
m at}(\mathbf{r})}$$

Non-empirical dispersion coefficients. *n*-body and any order. For instance:

$$C_{6,ij} = rac{lpha_i lpha_j \langle M_1^2
angle \langle M_1^2
angle_j}{\langle M_1^2
angle lpha_j + \langle M_1^2
angle_j lpha_i}$$

We include: two-body terms C_6 , C_8 and C_{10} .

Johnson, E. R. and Becke, A. D., J. Chem. Phys. 124 (2006) 174104

Imple	ementatio	n for molecu	امد			
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XDM implemented post-**Gaussian 09** using the **postg** program. Also **nwchem** (available upon request).

From the wfn file, **postg** gives:

- XDM dispersion coefficients, volumes, polarizabilities
- XDM dispersion energy
- forces for geometry optimization (fixed coefficients)
- second derivatives for frequencies
- Hirshfeld charges

Download **postg** from the XDM page at: http://faculty1.ucmerced.edu/ejohnson29

Kannemann, F. O. and Becke, A. D., *J. Chem. Theory Comput.* **6** (2010) 1081 Otero-de-la-Roza, A. and Johnson, E. R., *J. Chem. Phys.* **138** (2013) 204109

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Implementation for solids

- PS/PW (Quantum ESPRESSO)
- Solids Uniform 3D grid:
 - $d_{x\sigma}$, valence τ , ρ .
 - ω_i , all-electron ρ , ρ_{at} .
- Computational cost.
 - Comparable to DFT-D.
 - E_{disp} fast compared to E_{DFT} .

• Optimization: atomic forces and stresses.



Insensitive to grid density (CO ₂)							
$n_{\rm grid} =$	64	80	120				
<i>C</i> ₆ (C-C)	22.300	22.425	22.426				
<i>C</i> ₆ (O-O)	11.580	11.627	11.627				
$E_{\rm disp}$ (Ry)	-0.062965	-0.063374	-0.063374				

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Damping function parametrization

$$E_{disp} = -\frac{1}{2} \sum_{ij} \frac{C_6 f_6(R_{ij})}{R_{ij}^6} + \left[\frac{C_8 f_8(R_{ij})}{R_{ij}^8} + \frac{C_{10} f_6(R_{ij})}{R_{ij}^{10}} + \dots \right]$$
$$f_n(R_{ij}) = \frac{R_{ij}^n}{R_{ij}^n + (\mathbf{a_1} R_{ij,c} + \mathbf{a_2})^n}$$



Supercell calculations.



Domo		000000	000	000	00	0
Para	metrizatio	n set				

49 gas-phase dimers from Kannemann and Becke; JCTC 6 (2010) 1081.



Statis	stics of the	e fit (solids)				
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	Training	set (KB-	49)					
B86bPBE PW86PBE BLYP								
a_1	0.684	0.407	0.934	0.267	0.774			
$a_2(\AA)$	1.368	2.415	0.965	2.227	0.839			
MAE (kcal/mol)	0.41	0.46	0.42	0.41	0.31			
MAPE	11.3	13.8	11.8	14.3	<i>9</i> .8			
<u>S22</u>								
MAE (kcal/mol)	0.43	0.46	0.35	0.32	0.22			
MAPE	7.00	8.12	5.92	8.24	4.85			

Statis	tics of the	e fit (molecul	es)			
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XDM with aug-cc-pVTZ; mean absolute errors in kcal/mol.

Pure functionals:

Quantity	BLYP	PW86	PBE
MAE	0.31	0.40	0.50
MA%E	<i>9</i> .8	11.8	14.3

Hybrid and range-separated functionals:

Quantity	B3LYP	BH&HLYP	PBE0	CAM-B3LYP	LC - ωPBE
MAE	0.28	0.37	0.41	0.39	0.28
MA%E	6.7	7.8	10.2	<i>8.3</i>	7.8

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Role	of exchan	ge				

- The exact exchange potential decays as -1/r far from a molecule.
- In terms of the exchange hole, h_X remains on the molecule as the reference point moves away from it.
- The -1/r asymptotic dependence was used to design the B88 exchange functional.
- Functionals based on B88 or range-separated hybrids with the full exact-exchange limit (LC- ω PBE) give more accurate intermolecular exchange contributions.

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Benc	hmark set	S				

Mean absolute errors in XDM binding energies with aug-cc-pVTZ

(kcal/mol)	BLYP	PW86	B3LYP	LC - ωPBE
<i>S22</i>	0.22	0.35	0.31	0.31
S66	0.22	0.29	0.25	0.20
HSG	0.20	0.17	0.12	0.23

S22 and HSG reference data: Marshall *et al.* JCP **135** (2011) 194102.
S66 reference data: Rezac *et al.* JCTC **7** (2011) 2427.
See JCP **138** (2013) 204109 for additional data.

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Prediction of sublimation enthalpies

Benchmark:

- No reference wave-function data.
- Experimental sublimation enthalpies not directly comparable.

Phase Transition Enthalpy Measurements of Organic and Organometallic Compounds. Sublimation, Vaporization and Fusion Enthalpies From 1880 to 2010

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(Received 14 January 2010; accepted 15 January 2010; published online 4 October 2010)

- 21 crystals, small systems, low polymorphism.
- Well known sublimation enthalpies at or below room temperature.
- Different intermolecular interactions.

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 $\Delta H_{\rm sub}$: zero-point and thermal correction

$$\Delta H_{sub}(V,T) = E_{el}^{mol} + E_{trans} + E_{rot} + E_{vib}^{mol} + pV - \left(E_{el}^{crys} + E_{vib}^{crys}\right)$$

- $E_{\rm el}^{\rm crys} \longrightarrow \rm DFT+dispersion$
- $E_{\rm el}^{\rm mol} \longrightarrow \rm DFT+dispersion$, supercell
- $E_{\text{trans}} + E_{\text{rot}} + pV \longrightarrow 4RT (7/2RT)$
- Rigid molecule approximation $E_{vib}^{mol} = E_{vib}^{crys}$ for intramolecular
- Intermolecular $E_{\text{vib}}^{\text{crys}} \longrightarrow \text{Dulong-Petit } 6RT (5RT)$
- Zero-point vibrational contributions neglected
- Approximations tested for CO_2 crystal. Average experimental accuracy ≈ 1 kcal/mol.

Subl	imation e	nthalpies				
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Prediction of crystal structures

• Vibrational Helmholtz free energy:

$$F_{\text{vib}}(V,T) = \sum_{j=1}^{3n} \left[\frac{\omega_j}{2} + k_B T \ln \left(1 - e^{-\omega_j/k_B T} \right) \right]$$



Relax the crystal under negative pressure p_{th}

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Crystal structures



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Enantiomeric excess of amino-acids



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Enantiomeric excess of amino-acids

A simple model

- Same solvation energies.
- Same crystal temperature effects.
- $\Delta E = E_{dl} E_l$
- Predicted ee:

$$ee = \frac{\beta^2 - 1}{\beta^2 + 1} \times 100$$

 $\beta = e^{-\Delta E/RT}$

Amino acid	DFT	Expt.
Serine	100.0	100.0
Histidine	93.5	93.7
Leucine	92.2	87.9
Alanine	67.1	60.4
Cysteine	69.2	58.4
Tyrosine	70.6	51.7
Valine	62.3	44.1
Proline	0.0	39.7
Aspartic acid	0.0	0.0
Glutamic acid	0.0	0.0

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Enantiomeric excess



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Elec	trides					

An electride is an ionic substance in which a localized electron acts as an anion.

Existing electrides require a cage like structure to stabilise the cation: crown ethers and cryptands.

High magnetic susceptibilities, variable conductivities, very strong reducing agents.



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Elect	rides					

Use the NCI index to visualize the electrons - JACS 132 (2010) 6498.

Plots regions with low electron density and reduced density gradient.



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Elect	rides					

J. L. Dye used van der Waals radii to generate approximate channels and vacancies of electrons - JACS 1996, **118** (1996) 7329.



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Graphite step edges



Ye, Zhijiang et al. Appl. Phys. Lett., (2013) (in press).

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Summary						

- XDM implemented for gas-phase and solid-state.
- Excellent benchmarking results.
- Very accurate lattice energies and crystal geometries.
- Accurate enough to predict ee in solution.
- More: electrides, tribology,...

Download **postg**, QE+XDM, and CRITIC2 from:

http://faculty1.ucmerced.edu/ejohnson29