Application of RPA and Beyond-RPA methods to the A24 and S22 Test Sets

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The adiabatic connection fluctuation dissipation theorem (ACFDT) is a powerful theoretical framework to determine the Exchange-Correlation (xc) energy functional of DFT. Within the context of the ACFDT, it was shown that the dielectric matrix formulation, which is usually applied to calculate the random phase approximation (RPA) correlation energy, can be used for alternative RPA expressions including exchange effects. By using an approximate electron-hole exchange kernel two approximations were derived: An adiabatic connection analog of the second order screened exchange (AC-SOSEX) and the electron-hole time-dependent Hartree-Fock (eh-TDHF)[1]. It has been shown that this formalism is suitable for an efficient implementation within the plane-wave basis set[2].

This work introduces a new version of the code that optimizes the number of basis vectors involved in the response functions calculations. This is achieved by Gram-Schmidt orthogonalization and elimination of the linearly dependent vectors. The A24 test set, composed mostly of small dimers, is used to test the numerical accuracy of this procedure against the original implementation. This new efficient approach with Gram-Schmidt orthogonalization will then be used to compute the binding energies in the S22 test set, which involves weakly bound dimers with up to about 100 valence electrons.

Correlation Energies within ACFDT[2]

A24 Test Set [3]

Within the adiabatic connection fluctuation dissipation theorem,

$$E_{c}^{ACFDT} = -\frac{1}{2} \int_{0}^{1} d\lambda \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \operatorname{Tr} \left[\widetilde{\Pi}^{\lambda}(i\omega) \widetilde{K} - \widetilde{\Pi}^{0}(i\omega) \widetilde{K} \right]$$
$$\widetilde{\Pi}_{RPA}^{\lambda}(i\omega) = \left(\widetilde{I} - \lambda \widetilde{\Pi}^{0}(i\omega) \widetilde{K} \right)^{-1} \widetilde{\Pi}^{0}(i\omega),$$

Approximations for $\Pi^{\lambda}_{\text{RPA}}$:

• Random Phase Approximation (RPA) is the most used

$$\mathbb{F} = \mathbb{G} = \mathbb{V} = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix}, \qquad K_{vc,v'c'} = 2\langle \phi_v \phi_c | V | \phi_{v'} \phi_{c'} \rangle$$
$$E_c^{\text{RPA}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \operatorname{tr} \{ \log(\mathbf{I} - \mathbf{\Pi}_0(i\omega)\mathbf{K}) + \mathbf{\Pi}_0\mathbf{K} \}$$

Approximations beyond the RPA can be obtained by introducing an approximation in the exchange kernel [1, 2]:

$$\widetilde{\Xi} = \widetilde{B} + (\widetilde{\Xi} - \widetilde{B}) = \left(egin{array}{ccc} B & B \\ B & B \end{array}
ight) + \left(egin{array}{ccc} A - B & B \\ B & A - B \end{array}
ight) pprox \left(egin{array}{ccc} B & B \\ B & B \end{array}
ight)$$

 $A_{vc,v'c'} = 2\langle \phi_v \phi_c | V | \phi_{v'} \phi_{c'} \rangle - \langle \phi_v \phi_{v'} | V | \phi_c \phi_{c'} \rangle$ $B_{vc,v'c'} = 2\langle \phi_v \phi_c | V | \phi_{v'} \phi_{c'} \rangle - \langle \phi_v \phi_{c'} | V | \phi_{v'} \phi_c \rangle.$

• Adiabatic Connection Second Order Screened Exchange (AC-SOSEX)

$$\mathbb{F} = \mathbb{V} \text{ and } \mathbb{G} = \mathbb{W} \qquad \Rightarrow \qquad E_c^{\text{AC-SOSEX}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{\mathrm{d}\omega}{2\pi} \operatorname{tr} \left\{ \log(\boldsymbol{I} - \boldsymbol{\Pi}_0(i\omega)\boldsymbol{K})\boldsymbol{K}^{-1}\boldsymbol{B} + \boldsymbol{\Pi}_0\boldsymbol{B} \right\}$$

• Electron-Hole Time-Dependent Hartree-Fock (eh-TDHF)

Version with Re-orthogonalization

Original Version



 \rightarrow Good agreement with CCSDT(Q) for both versions

 \rightarrow MAE (kcal/mol) between two versions: 0.047 (RPA), 0.033 (eh-TDHF), 0.029 (AC-SOSEX)

$$\mathbb{F} = \mathbb{W} \text{ and } \mathbb{G} = \mathbb{V} \qquad \Rightarrow \qquad E_c^{\text{eh}-\text{TDHF}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{\mathrm{d}\omega}{2\pi} \operatorname{tr} \left\{ \log(\boldsymbol{I} - \boldsymbol{\Pi}_0(i\omega)\boldsymbol{B})\boldsymbol{B}^{-1}\boldsymbol{K} + \boldsymbol{\Pi}_0\boldsymbol{K} \right\}$$

Implementation Details

Using an auxiliary basis set $\Phi_Q(\mathbf{r})$ to more compactly represent the dielectric matrix, the RPA correlation energy requires evaluation of matrix elements of the type:

$$C_{QR}(i\omega) = 4 \sum_{v} \langle U_Q^v | \Re \left(i\omega + \epsilon_v - \hat{H} \right)^{-1} | U_R^v |$$

with $| U_G^v \rangle = \hat{P} V^{\frac{1}{2}} \Phi_G | \phi_v \rangle$

Beyond-RPA methods require additional terms such as,

 $J_{QG}^{vv'}(i\omega) = 2\langle U_Q^v | \Re(i\omega + \epsilon_v - \hat{H})^{-1} | U_G^{v'} \rangle$

Inclusion of different occupied states v and v' are the reason for the higher numerical complexity of Beyond-RPA methods.

To overcome this difficulties

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• Reduce the large number of right and left vectors by applying a Gram-Schmidt orthogonalization procedure

S22 Test Set 4

System		CCSD(T)	RPA	eh-TDHF	AC-SOSEX
Hydrogen-bonded systems					
1	ammonia dimer	-3.17	-2.54	-2.72	-2.71
2	water dimer	-5.02	-4.20	-4.58	-4.65
3	formic acid dimer	-18.80	-17.05	-18.24	-18.60
4	formamide dimer	-16.12	-14.75	-15.57	-15.71
5	uracil dimer (c2h)	-20.69	-18.78	-20.32	-21.02
6	2-pyridoxine 2 -aminopyridine	-17.00	-15.78	-16.46	-16.61
7	adeninethymine (WC)	-16.74	in progress	in progress	$in \ progress$
Mixed-type systems					
8	methane dimer	-0.53	-0.47	-0.52	-0.48
9	ethene dimer	-1.50	-1.15	-1.33	-1.31
10	benzenemethane	-1.45	in progress	$in \ progress$	$in \ progress$
11	benzene dimer	-2.62	in progress	$in \ progress$	$in \ progress$
12	pyrazine dimer	-4.20	-3.40	-3.39	-3.24
13	uracil dimer	-9.74	-8.61	-9.11	-9.26
14	indolebenzene (stack)	-4.59	-3.41	-3.35	-3.19
15	adeninethymine (stack)	-11.66	-10.36	-10.59	-10.62
Dispersion-dominated systems					
16	etheneethine	-1.51	-1.29	-1.50	-1.57
17	benzenewater	-3.29	-2.89	-3.10	-3.10
18	benzeneammonia	-2.32	-1.91	-2.07	-2.06
19	benzenehydrogen cyanide	-4.55	-4.03	-4.49	-4.61
20	benzene dimer	-2.71	-2.31	-2.45	-2.42
21	indolebenzene (T-shape)	-5.62	in progress	in progress	$in \ progress$
22	phenol dimer	-7.09	-6.31	-6.76	-6.79
	MAE		0.847	0.441	0.437
	ME		-0.06	-0.01	-0.05
	MAX		1.91	1.24	1.4

$$|O_1\rangle = |U_1^1\rangle/N_1$$

$$|O_2\rangle = \left[|U_1^2\rangle - \langle U_1^2|O_1\rangle |O_1\rangle \right] / N_2$$

$$O_{N_{aux} \times N_{v}} \rangle = \left[|U_{N_{aux}}^{N_{v}} \rangle - \sum_{i=1}^{N_{aux} \times N_{v}} \langle U_{N_{aux}}^{N_{v}} | O_{i} \rangle | O_{i} \rangle \right] / N_{N_{aux} \times N_{v}} \rangle$$

• If a certain $|O_n\rangle$ is linearly dependent to all previous ones $(|O_1\rangle, \ldots, |O_{n-1}\rangle)$ then it does not need to be included in the orthonormalized basis set.



Good agreement of Beyond-RPA methods with CCSD(T) methods.

Further Reading

B. Mussard, D. Rocca, G. Jansen, and J. Ángyán. JCTC, 12:2191, 2016.

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