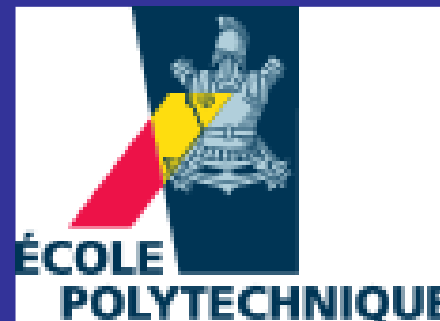


Overview of RPA from the condensed matter perspective

Lucia Reining
Palaiseau Theoretical Spectroscopy Group



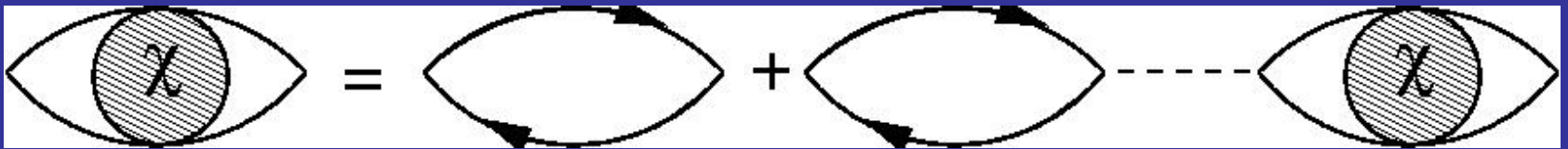
Outline

1. History
2. Look at the response
3. Look at a spectral function
4. Look at the total energy

1. History and Definition

Bohm and Pines 1952: response of an electron gas to an effective single particle potential

Equivalence (*Ehrenreich and Cohen 1959*) to the linearized time-dependent Hartree Approximation (*Lindhard 1954*)



RPA ~ bubble approximation (*Hubbard 1957*)

No exchange contribution

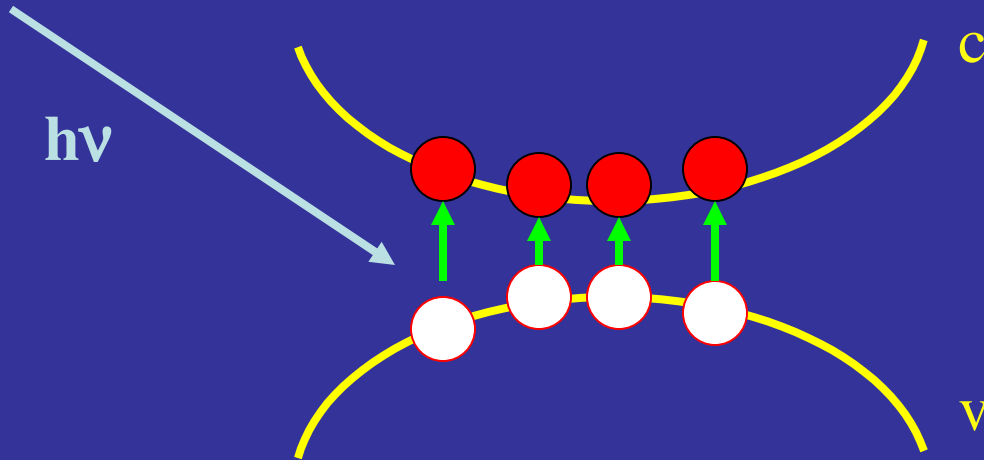
2. Look at the response

a. Connection to experiment

b. Long range and short range

c. Failures

Independent electrons and transitions

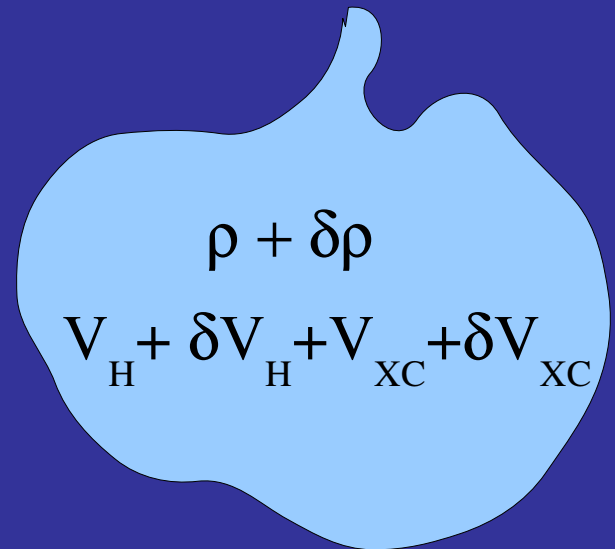
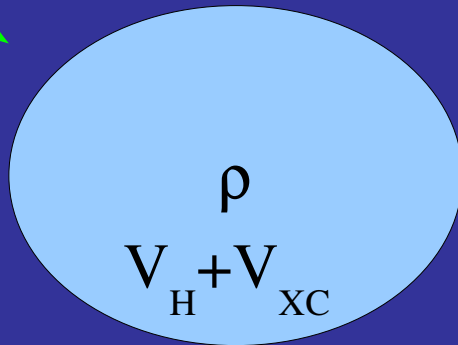


$$\text{Im} [\chi_{00}] \sim - \sum_{vc} |\langle v | D | c \rangle|^2 \delta (E_c - E_v - \omega)$$

Partial DOS \rightarrow EELS \sim XAS

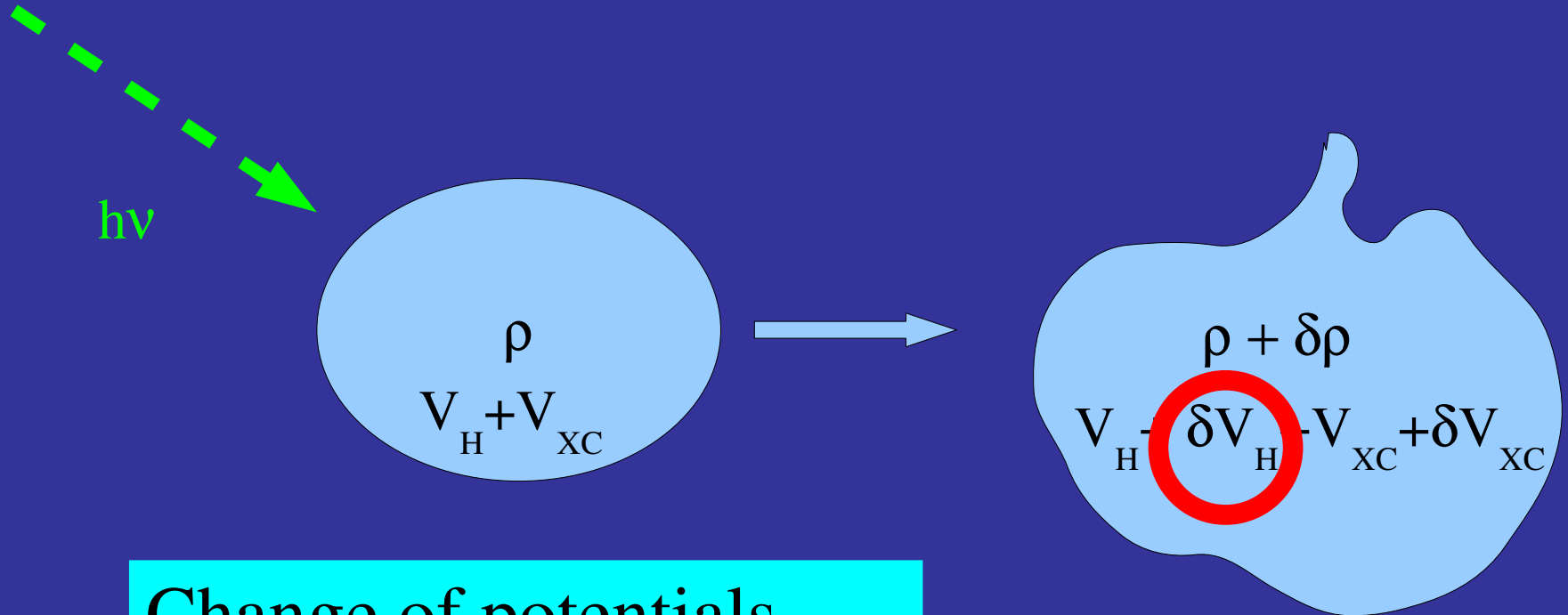
Excitation ?

$h\nu$



Change of potentials

Excitation ?



Change of potentials

Induced Hartree: long-range and local field effects

RPA

RPA:

$$\epsilon = 1 - v\chi_0$$

$$V_{\text{ext}} = \epsilon V_{\text{tot}}$$

$$\epsilon^{-1} = (1 - v\chi_0)^{-1} = 1 + v\chi_0 = 1 + v\chi_0(1 - v\chi_0)^{-1}$$

$$\chi = \chi_0 + \chi_0 v\chi$$

$$\bar{\chi} = \chi_0 + \chi_0 \bar{v} \bar{\chi}$$

$$\delta V_H / \delta \rho$$

$$\delta V_H / \delta \rho \quad (\text{micro})$$

$$V_{\text{ind}} = v\chi V_{\text{ext}}$$

$$V_{\text{ind}} = v\chi (V_{\text{ext}} + V_{\text{ind}}^0)$$

$$\text{Loss: } -\text{Im}(\chi)$$

$$\text{Abs: } \text{Im}(1/\epsilon^{-1}) = v\text{Im}(\bar{\chi})$$

$$\delta v_H / \delta \rho$$

$$\text{TDDFT-RPA: } \hat{U} \chi = \chi_0 + \chi_0 [v] \chi$$

$$v(q+G) \sim 1/|q+G|^2$$

$G=0 \rightarrow$ plasmon

$G \neq 0 \rightarrow$ crystal local field effects

RPA:

$$\epsilon = 1 - v\chi_0$$

$$V_{\text{ext}} = \epsilon V_{\text{tot}}$$

$$\epsilon^{-1} = (1 - v\chi_0)^{-1} = 1 + v\chi_0 = 1 + v\chi_0(1 - v\chi_0)^{-1}$$

$$\chi = \chi_0 + \chi_0 v\chi$$

$$\bar{\chi} = \chi_0 + \chi_0 \bar{v} \bar{\chi}$$

$$\delta V_H / \delta \rho$$

$$\delta V_H / \delta \rho \quad (\text{micro})$$

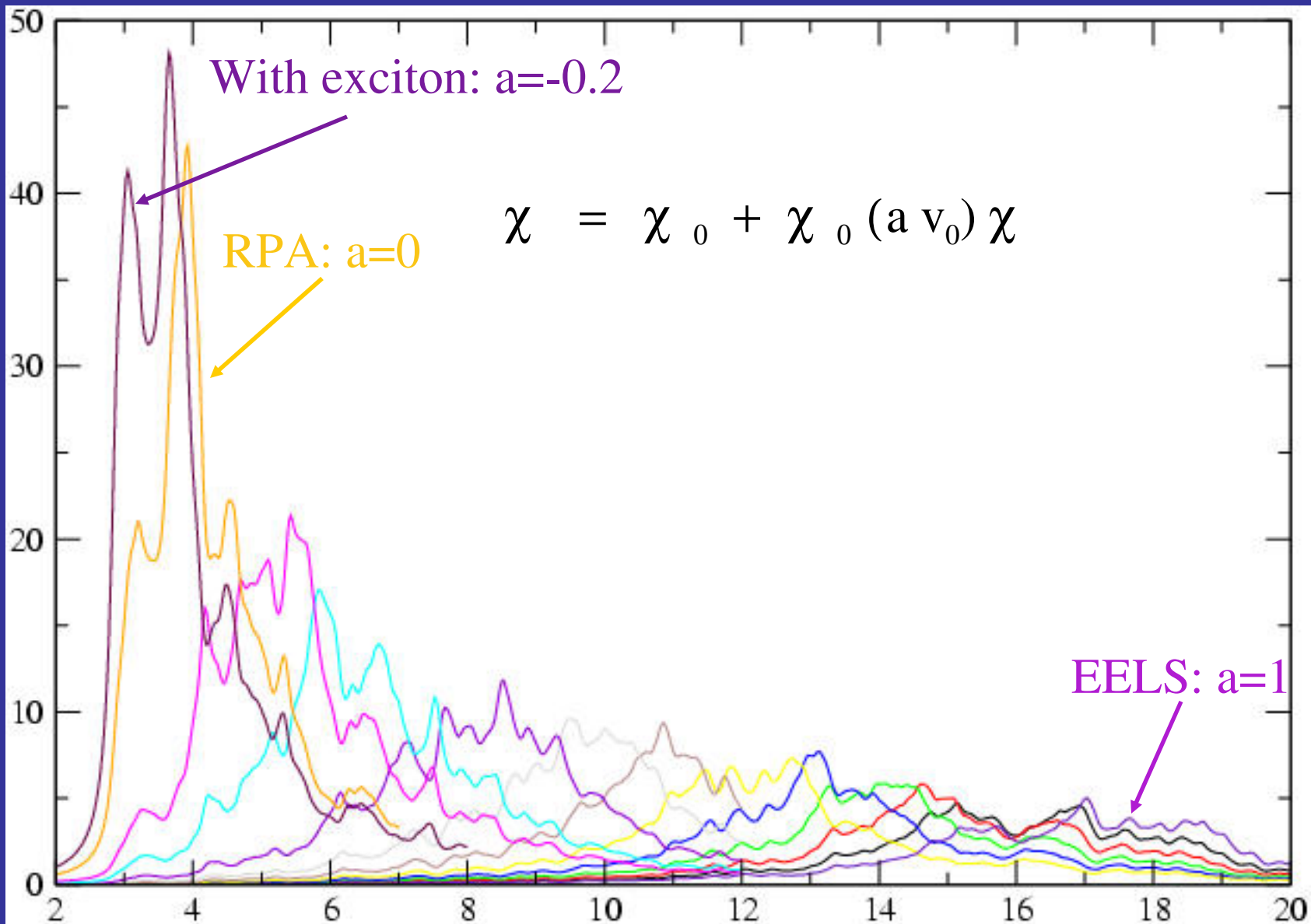
$$V_{\text{ind}} = v\chi V_{\text{ext}}$$

$$V_{\text{ind}} = v\chi (V_{\text{ext}} + V_{\text{ind}}^0)$$

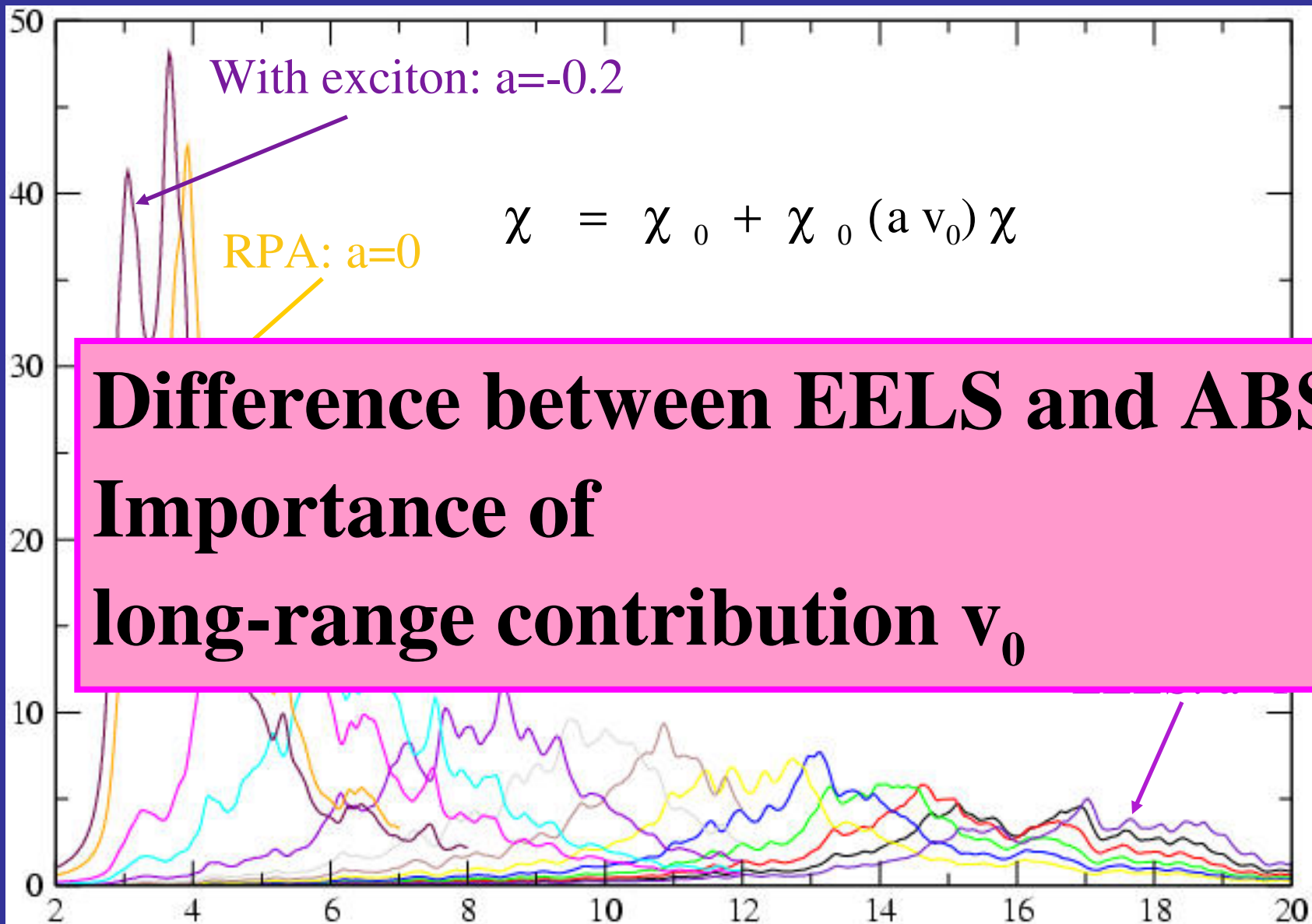
$$\text{Loss: } -\text{Im}(\chi)$$

$$1/q^2$$

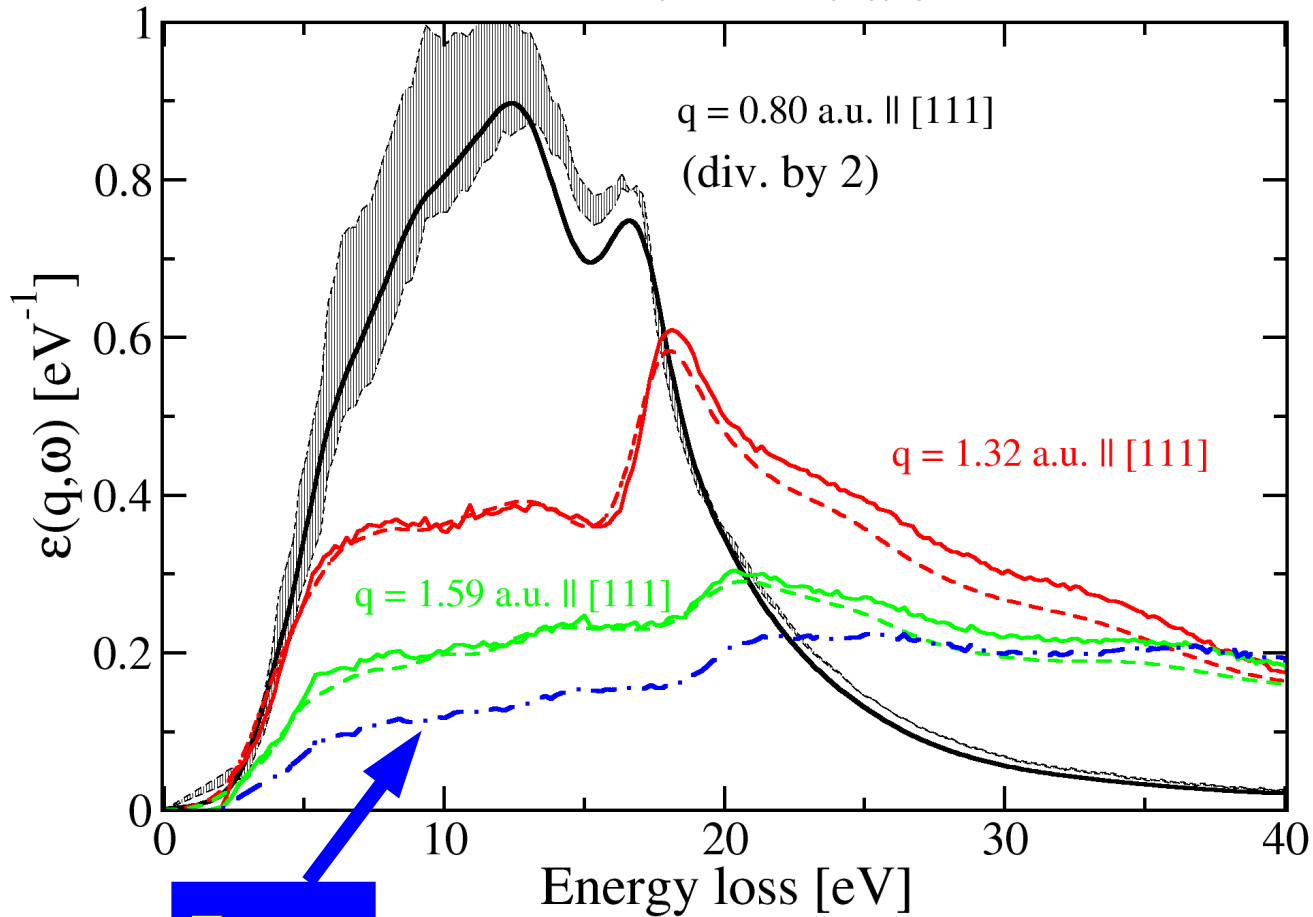
$$\text{Abs: } \text{Im}(1/\epsilon^{-1}) = v\text{Im}(\bar{\chi})$$



Sottile et al, Int. J. Quant. Chem. (2002)

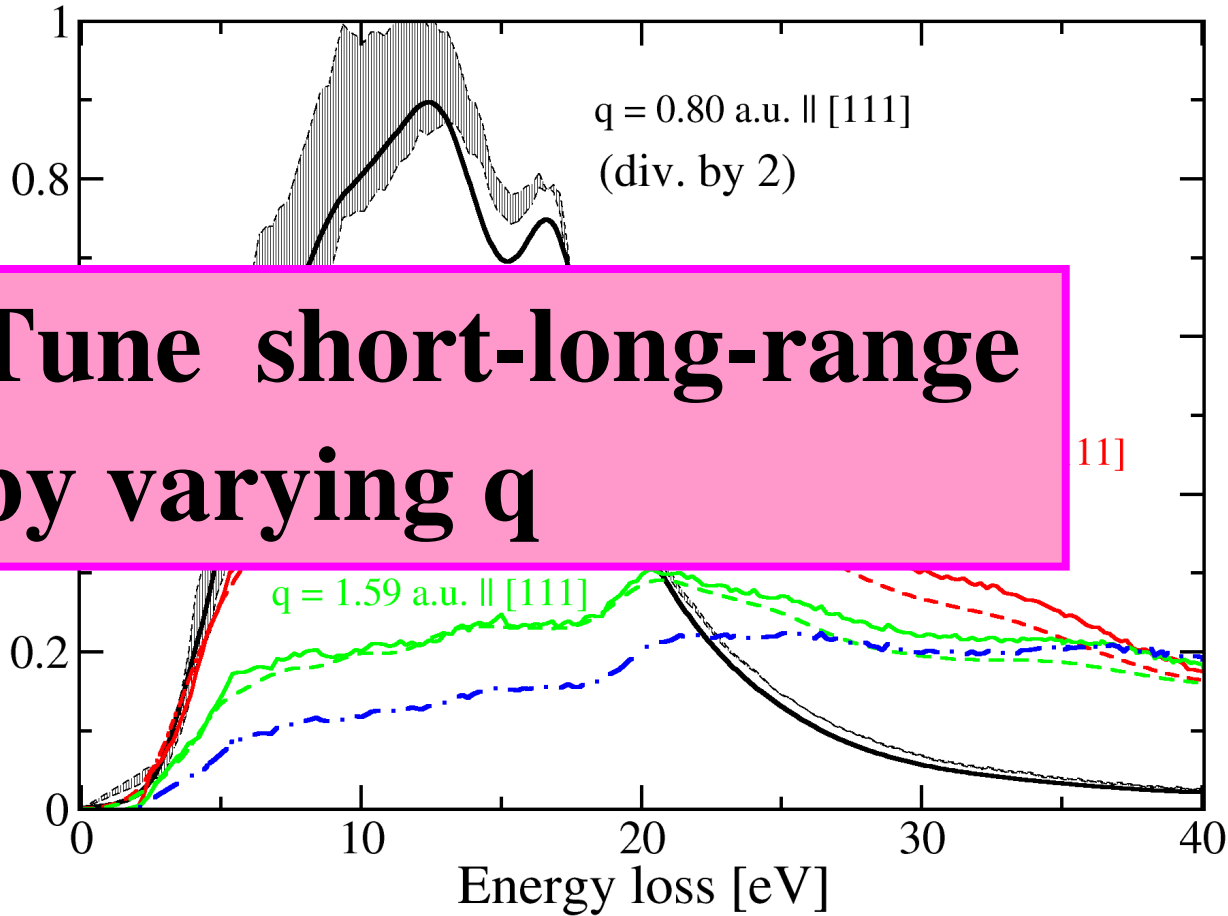


Sottile et al, Int. J. Quant. Chem. (2002)



Loss

/users/c10/weissker/Dpsrc452_it_gh/Test_2048/eps_2_paper.agr



RPA including crystal local field effects: \bar{v}

$$\epsilon = 1 - v\chi_0$$

$$V_{\text{ext}} = \epsilon V_{\text{tot}}$$

$$\epsilon^{-1} = (1 - v\chi_0)^{-1} = 1 + v\chi_0 = 1 + v\chi_0(1 - v\chi_0)^{-1}$$

$$\chi = \chi_0 + \chi_0 v\chi$$

$$\bar{\chi} = \chi_0 + \chi_0 \bar{v} \bar{\chi}$$

$$\delta V_H / \delta \rho$$

$$\delta V_H / \delta \rho \quad (\text{micro})$$

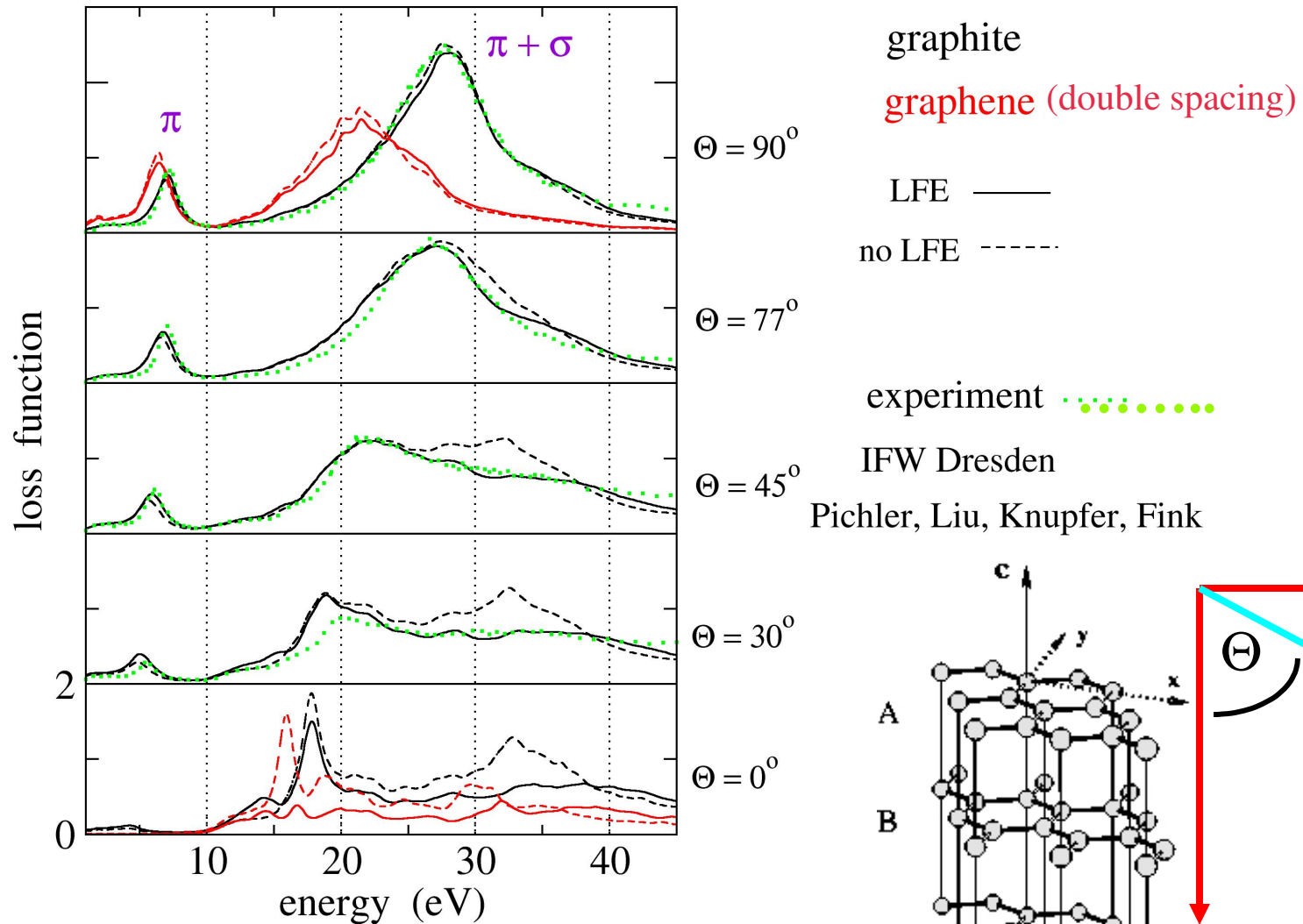
$$V_{\text{ind}} = v\chi V_{\text{ext}}$$

$$V_{\text{ind}} = v\chi (V_{\text{ext}} + V_{\text{ind}}^0)$$

$$\text{Loss: } -\text{Im}(\chi)$$

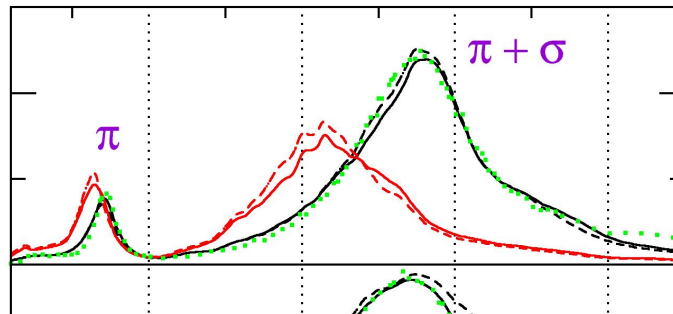
$$\text{Abs: } \text{Im}(1/\epsilon^{-1}) = v\text{Im}(\bar{\chi})$$

loss function versus momentum orientation



Marinopoulos, Reining, Olevano et al. **PRL** 89, 076402 (2002)

loss function versus momentum orientation



graphite

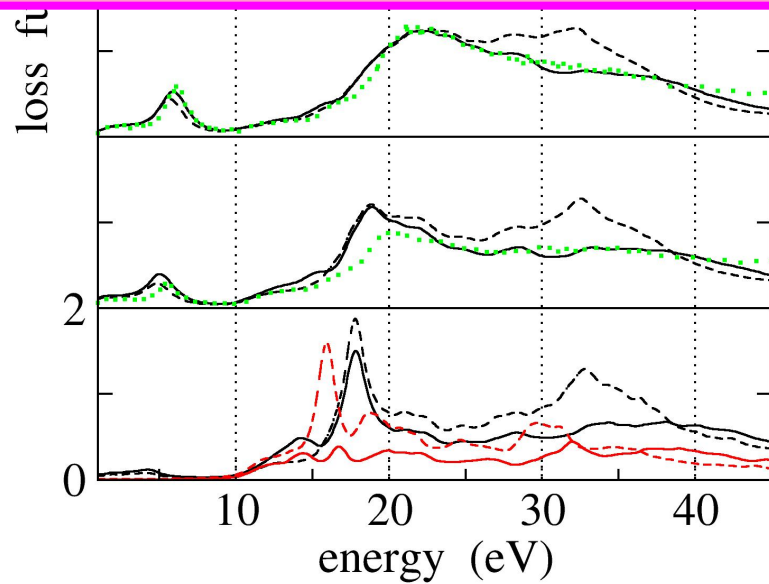
graphene (double spacing)

LFE ———

no LFE - - - -

$\Theta = 90^\circ$

LFE tell us about polarizable objects



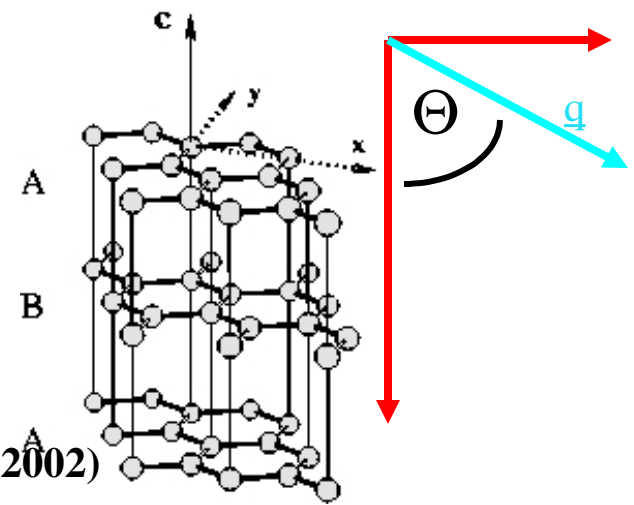
$\Theta = 45^\circ$

$\Theta = 30^\circ$

$\Theta = 0^\circ$

IFW Dresden

Pichler, Liu, Knupfer, Fink

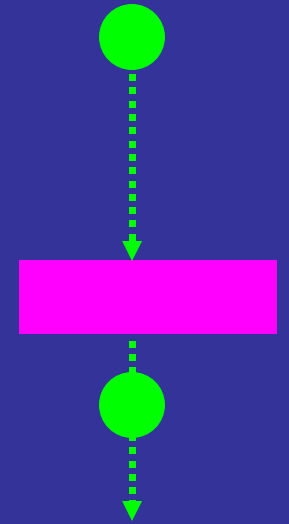
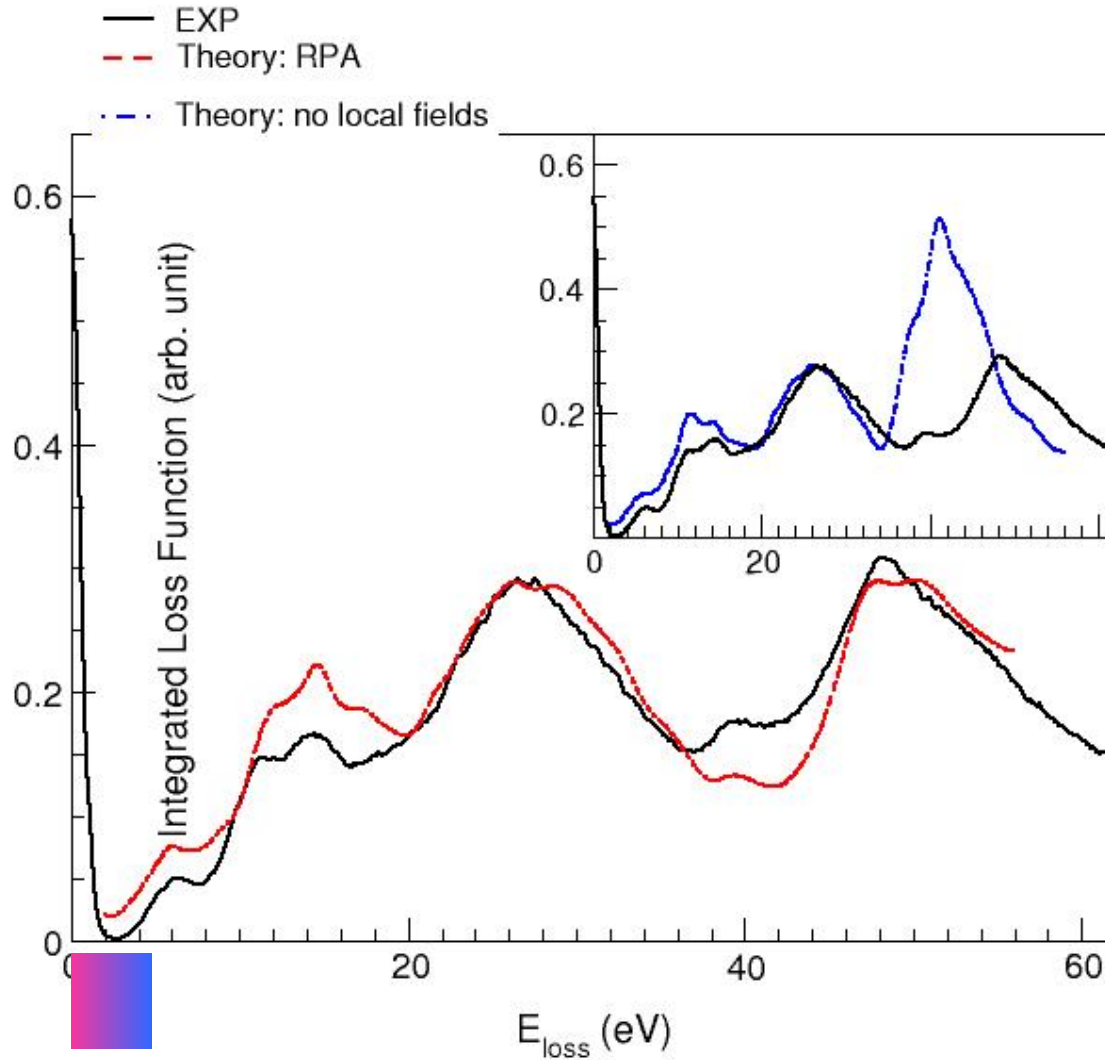


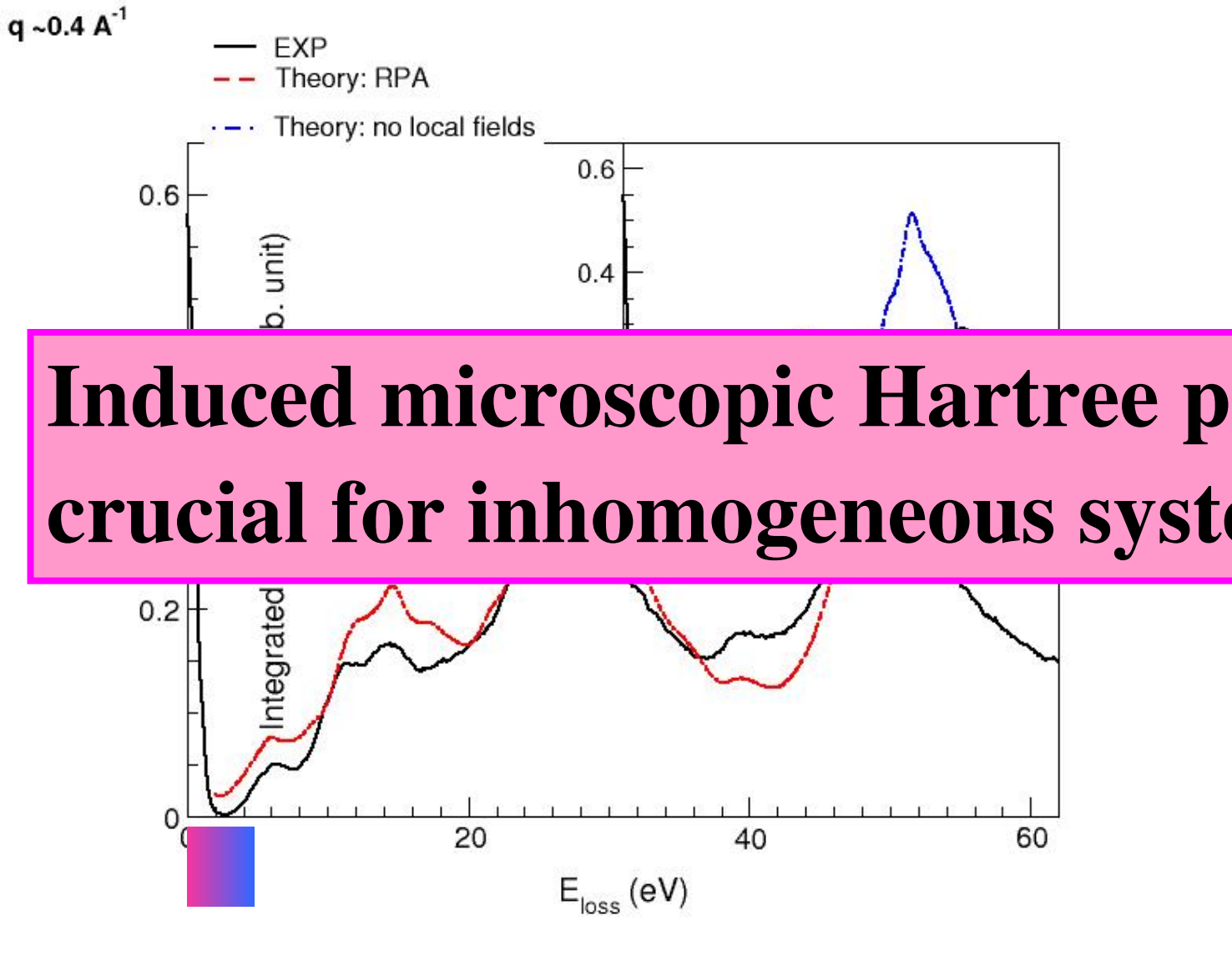
Marinopoulos, Reining, Olevano et al. *PRL* **89**, 076402 (2002)

Rutile TiO₂

Titanium Dioxide

$q \sim 0.4 \text{ \AA}^{-1}$





**Induced microscopic Hartree potential
crucial for inhomogeneous system**

$$\delta V_H / \delta \rho$$

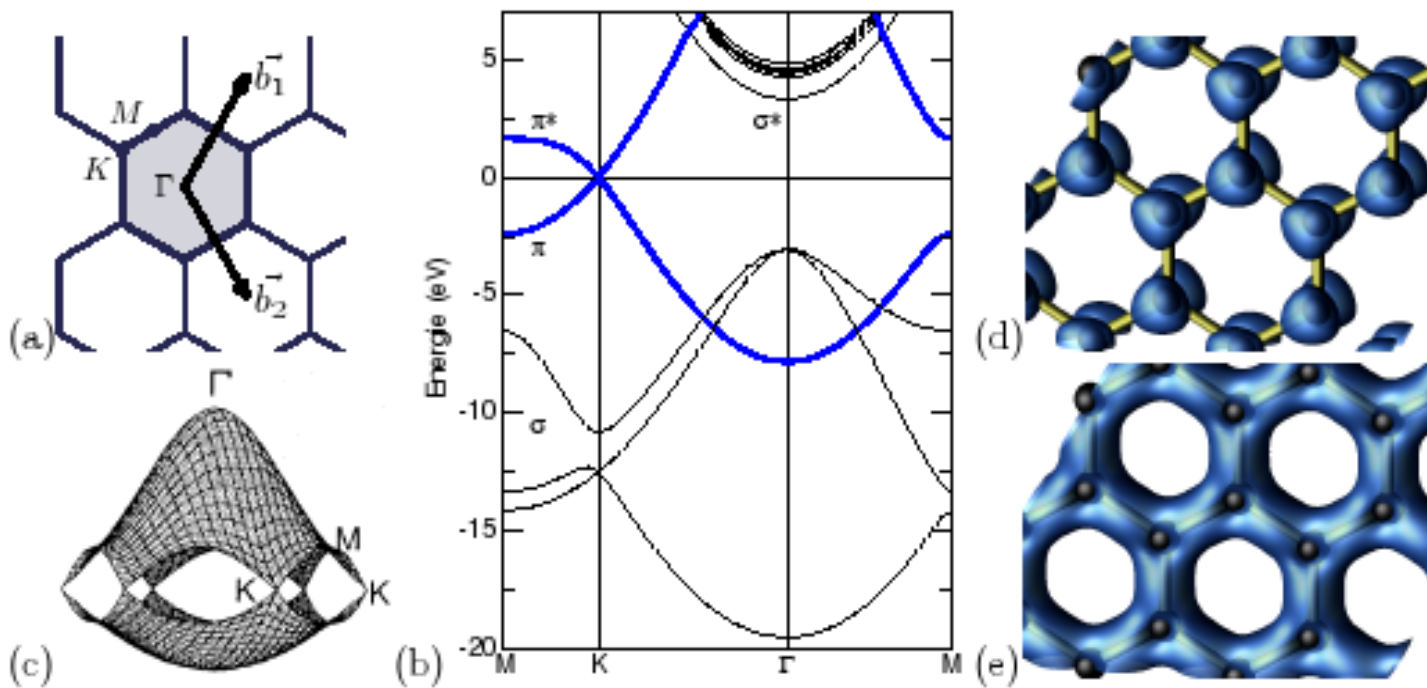
$$\text{RPA:} \quad \chi = \chi_0 + \chi_0 [v] \chi$$

$$v(q+G) \sim 1/|q+G|^2$$

$G=0 \rightarrow$ plasmon

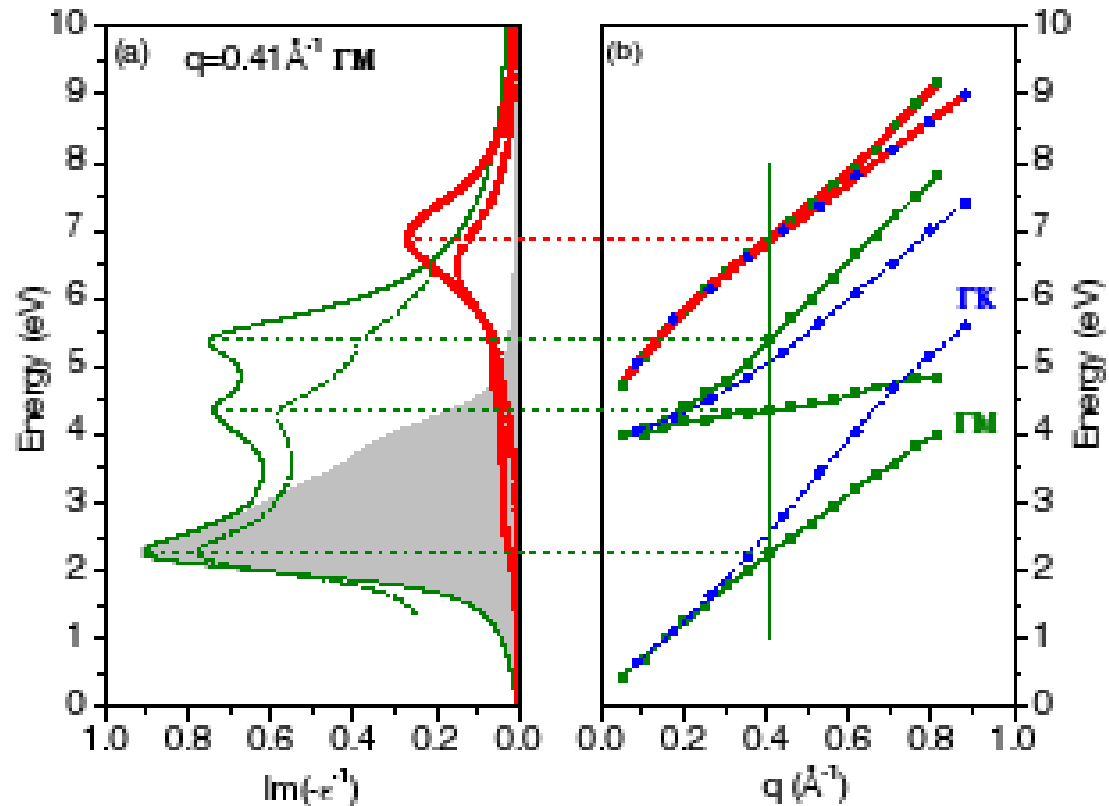
$G \neq 0 \rightarrow$ crystal local field effects

Graphene, π plasmon



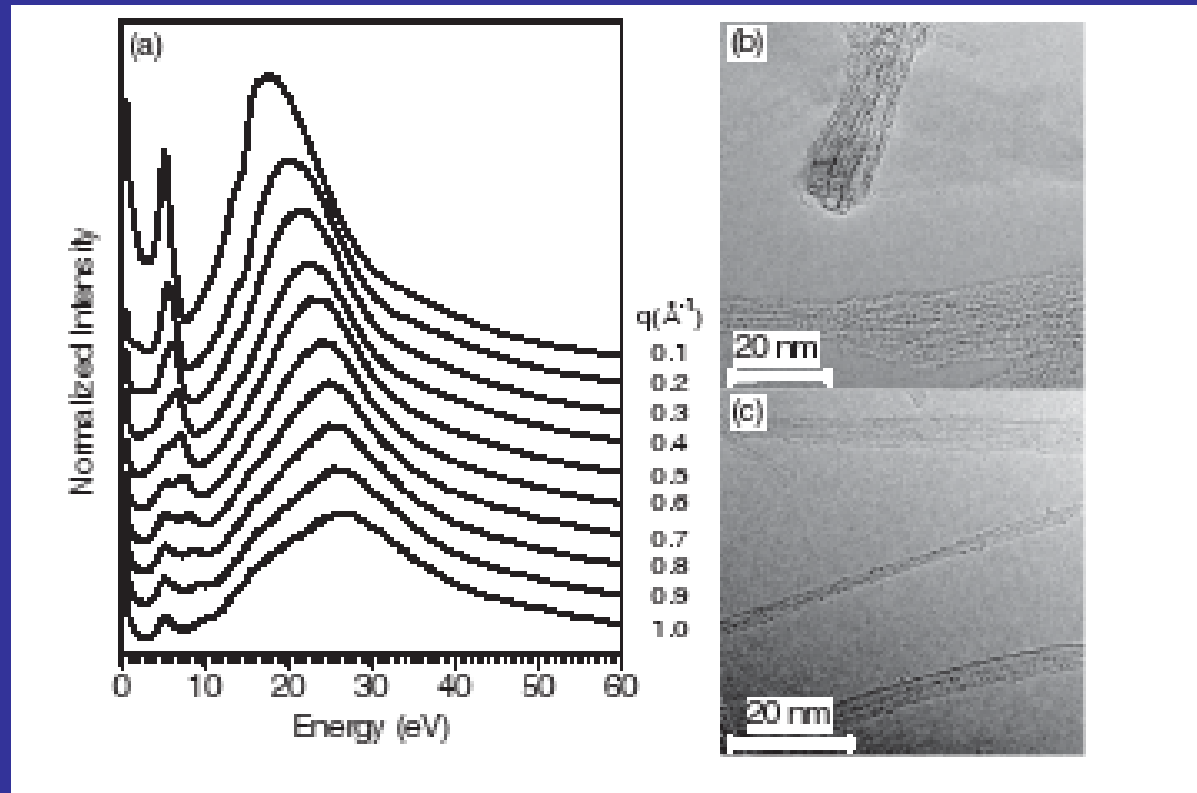
R. Hambach, Diplomarbeit

Graphene, π plasmon



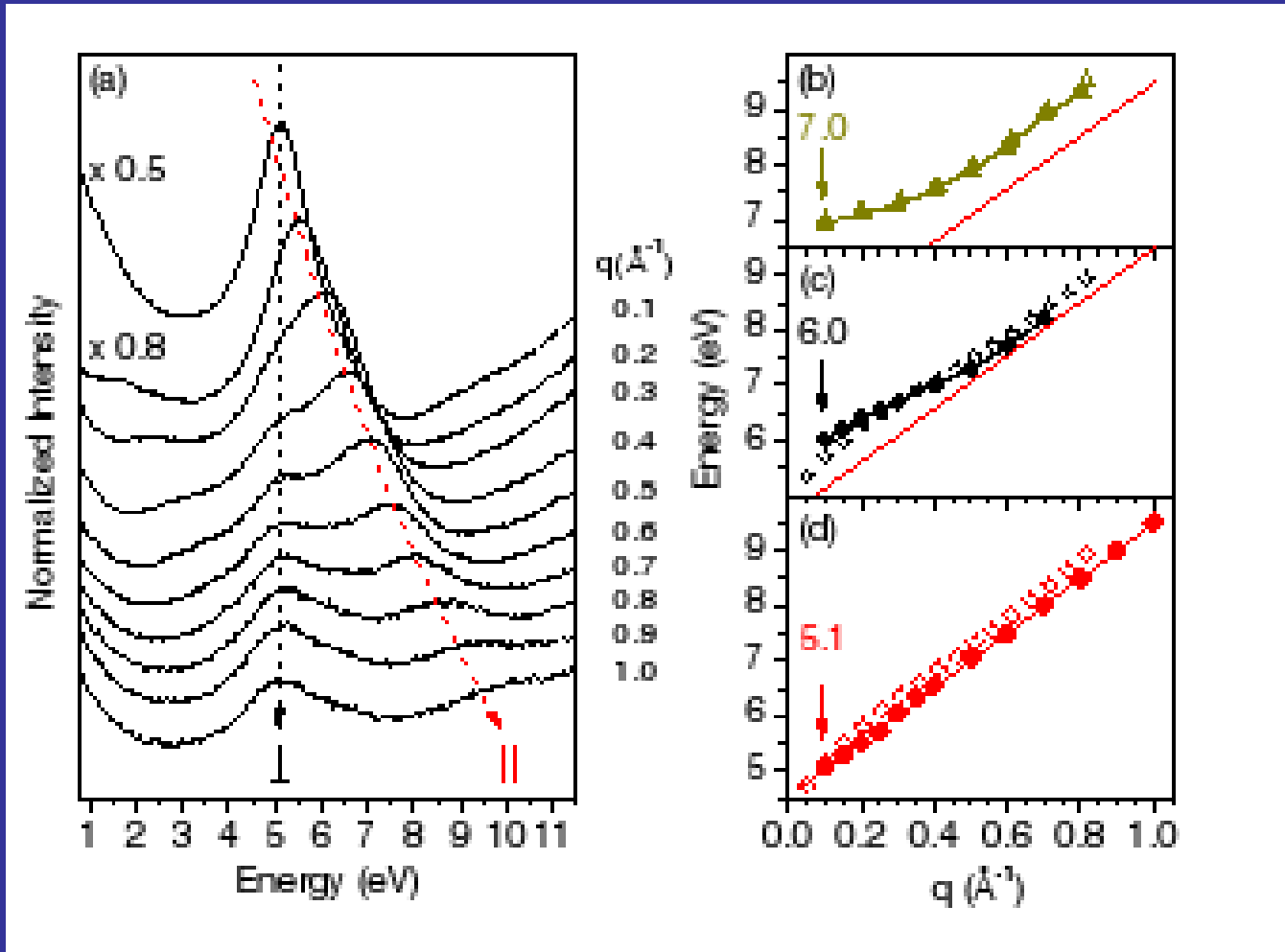
C. Kramberger et al., PRL 100, 196803 (2008)

Isolated carbon nanotube, π plasmon



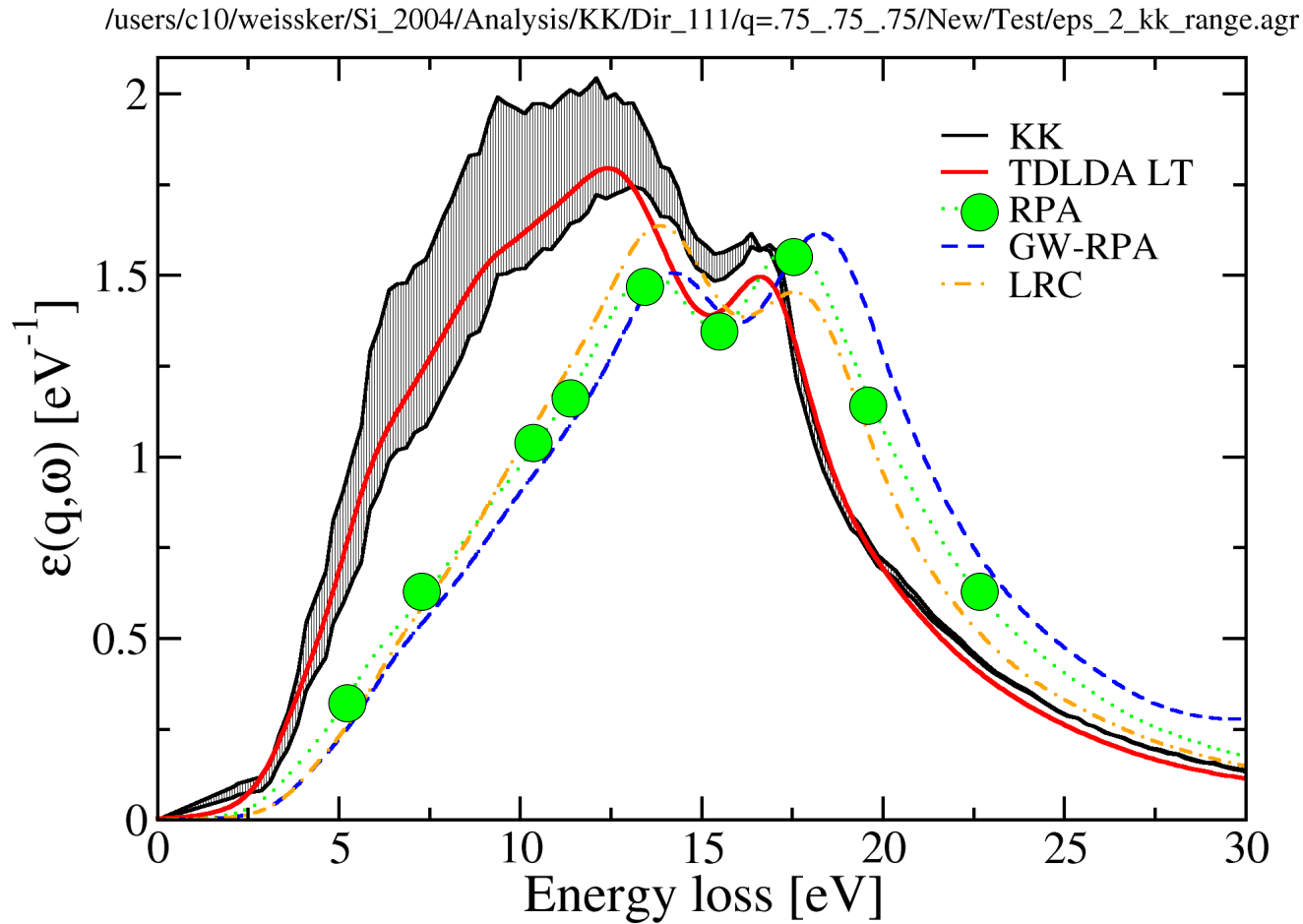
C. Kramberger et al., PRL 100, 196803 (2008)

Nanotubes and graphene, π plasmon dispersion

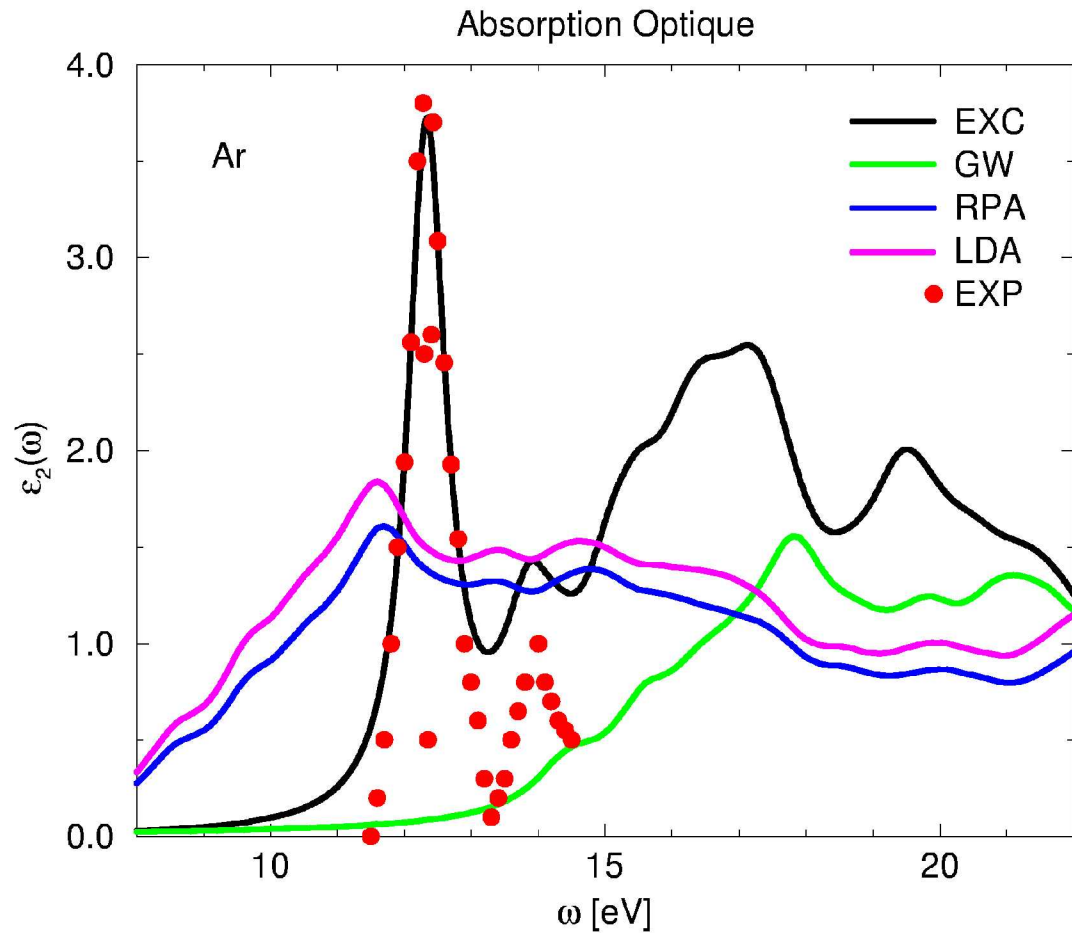


C. Kramberger et al., PRL 100, 196803 (2008)

More quantitatively?



H. Weissker et al., Phys. Rev. Lett. **97**, 237602 (2006)
Collaboration LSI – ESRF (ID16)



...or even qualitatively?

3. Look at a spectral function

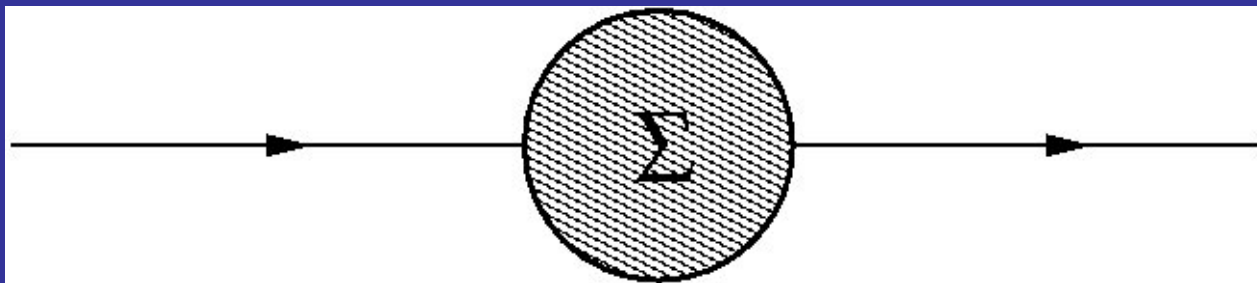
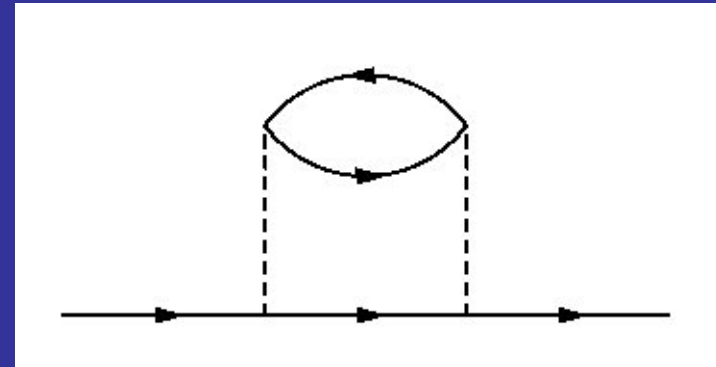
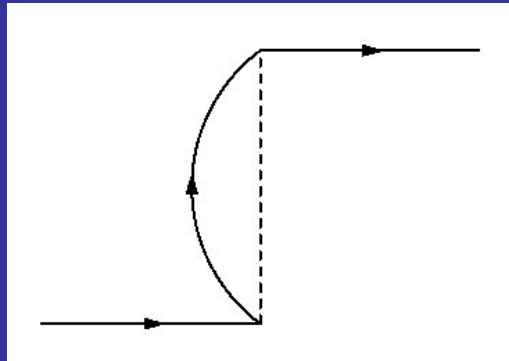
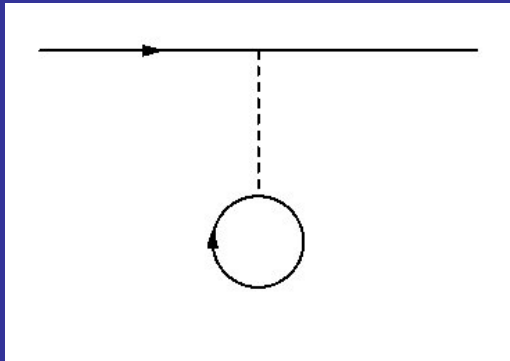
- a. Appearance of the (RPA) response function
- b. Effects of the (RPA) response function
- c. Failures

$G = \dots$



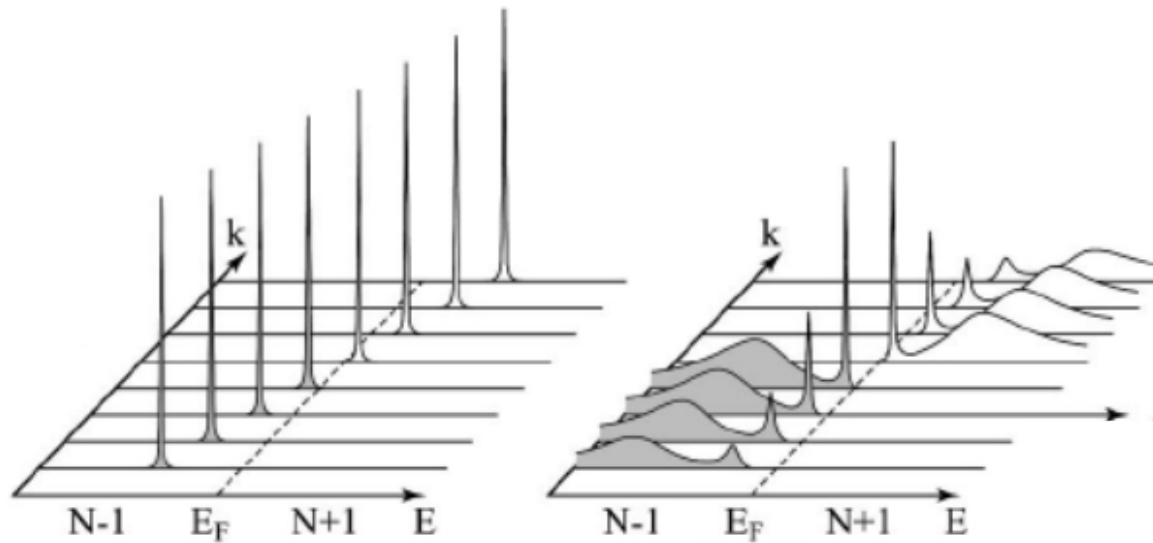
$$G(1,2) = -i \langle T[\psi(1)\psi^\dagger(2)] \rangle$$

$$1=(r_1, \sigma_1, t_1)$$



$$\text{Dyson equation: } G = G_0 + G_0 \Sigma G$$

$$A(\omega) \sim \text{Im}[G(\omega)]$$



From Damascelli et al., RMP 75, 473 (2003)

$$V_H(12) = -i\delta(12)G(33^+)v_c(31)$$

$$\Sigma_x(12) = iG(12^+)v_c(21)$$

$$\Sigma_{xc}(1, 2) = \Sigma_x + iG(1\bar{3})\Xi(\bar{3}, \bar{5}; 2, \bar{6})L(\bar{6}\bar{2}; \bar{5}\bar{2})v_c(\bar{2}, 1) \quad (1)$$

$$\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}') := -i\delta(\bar{3}\bar{3}')\delta(\bar{2}', \bar{2})v_c(\bar{3}\bar{2}) + \frac{\delta\Sigma_{xc}(\bar{3}, \bar{3}')}{\delta G(\bar{2}', \bar{2})} \quad (2)$$

Interaction = variation of “potential”

$$L(1, 2, 1', 2') = L_0(1, 2, 1', 2') + L_0(1, \bar{3}', 1', \bar{3})\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}')L(\bar{2}', 2, \bar{2}, 2') \quad (3)$$

What does the system do (create e-h pairs...)

$$L_0(1, 2, 1', 2') = G(1, 2')G(2, 1') \quad (4)$$

$$V_H(12) = -i\delta(12)G(33^+)v_c(31)$$

$$\Sigma_x(12) = iG(12^+)v_c(21)$$

$$\Sigma_{xc}(1, 2) = \Sigma_x + iG(1\bar{3})\Xi(\bar{3}, \bar{5}; 2, \bar{6})L(\bar{6}\bar{2}; \bar{5}\bar{2})v_c(\bar{2}, 1) \quad (1)$$

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$$L_0(1, 2, 1', 2') = G(1, 2')G(2, 1') \quad (4)$$

System does nothing: HF

Only classical part of Ξ : $L \rightarrow i\chi$, $\Sigma_{xc} \rightarrow iG(v_c + v_c \chi v_c) = iGW$

$$V_H(12) = -i\delta(12)G(33^+)v_c(31)$$

$$\Sigma_x(12) = iG(12^+)v_c(21)$$

$$\Sigma_{xc}(1, 2) = \Sigma_x + iG(1\bar{3})\Xi(\bar{3}, \bar{5}; 2, \bar{6})L(\bar{6}\bar{2}; \bar{5}\bar{2})v_c(\bar{2}, 1) \quad (1)$$

$$\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}') := -i\delta(\bar{3}\bar{3}')\delta(\bar{2}', \bar{2})v_c(\bar{3}\bar{2}) + \frac{\delta\Sigma_{xc}(\bar{3}, \bar{3}')}{\delta G(\bar{2}', \bar{2})} \quad (2)$$

Interaction = variation of “potential”

$$L(1, 2, 1', 2') = L_0(1, 2, 1', 2') + L_0(1, \bar{3}', 1', \bar{3})\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}')L(\bar{2}', 2, \bar{2}, 2') \quad (3)$$

What does the system do (create e-h pairs...)

$$L_0(1, 2, 1', 2') = G(1, 2^+)v_c(2, 1')$$

W consistently in RPA

System does nothing: HF

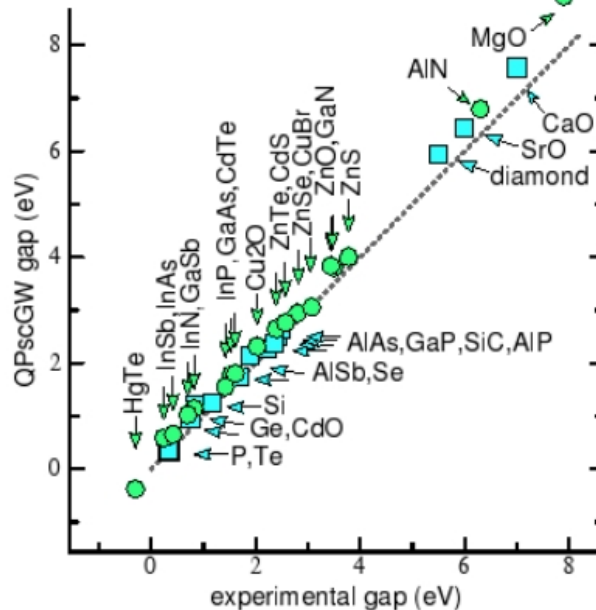
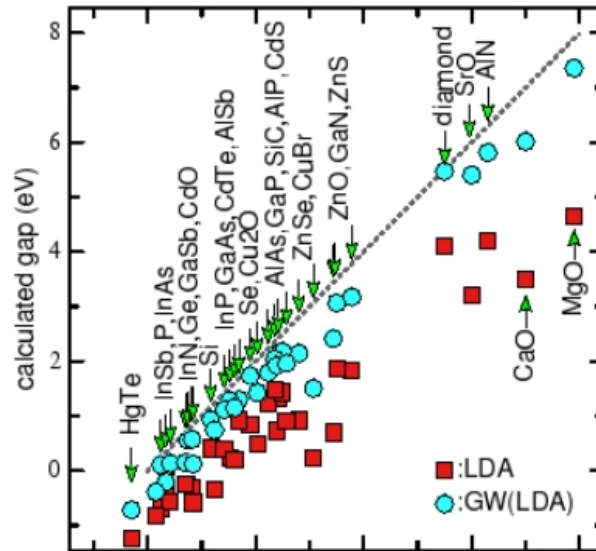
Only classical part of Ξ : $L \rightarrow i\chi$, $\Sigma_{xc} \rightarrow iG(v_c + v_c \chi v_c) = iGW$

RPA screening in GW?

→ self-consistent screening of Fock exchange

localization

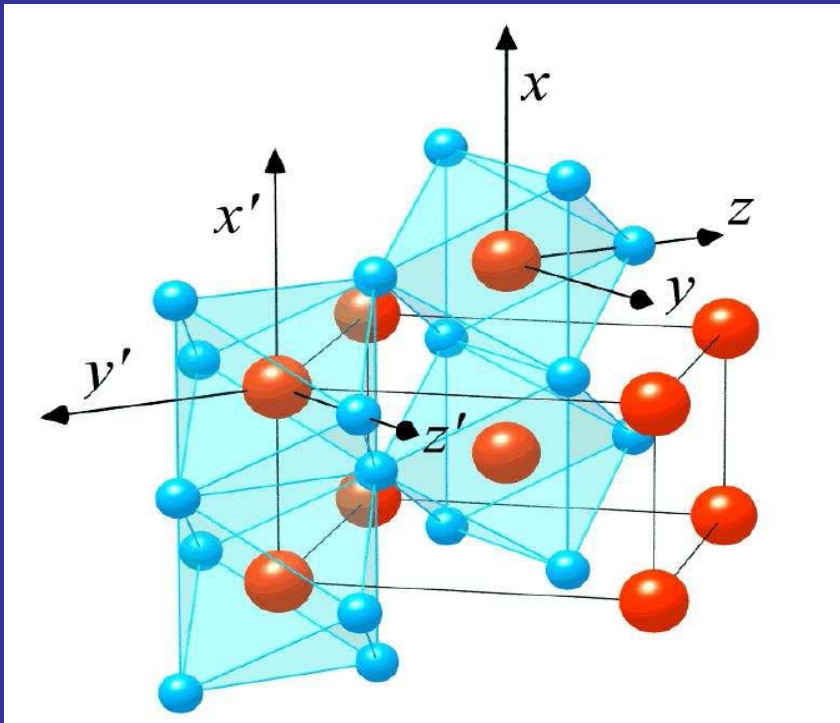
Improve QP energies and wavefunctions



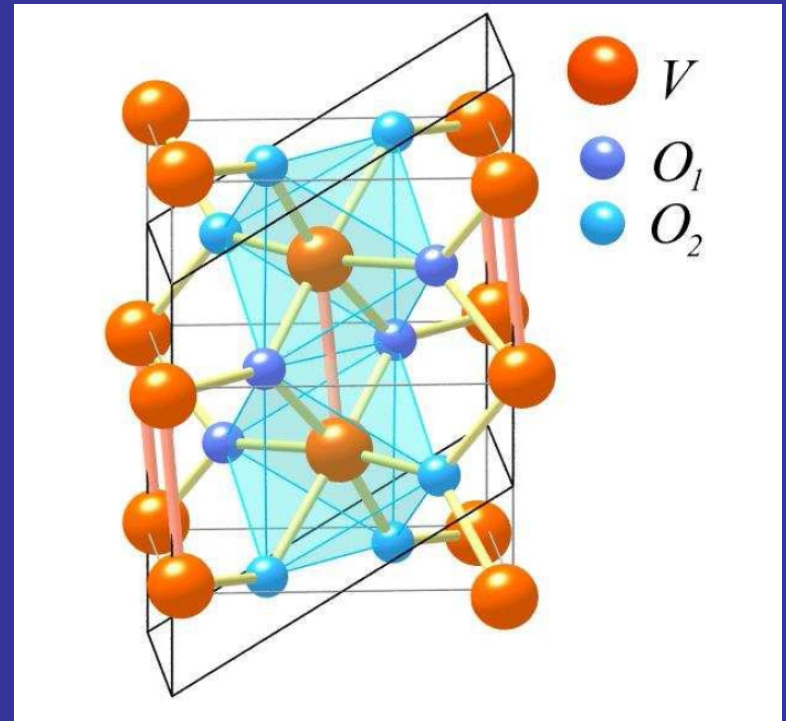
“RPA form” (bubble)

van Schilfgaarde, Kotani, Faleev,
Phys. Rev. Lett. 96, 226402 (2006)

Even qualitatively important - VO_2

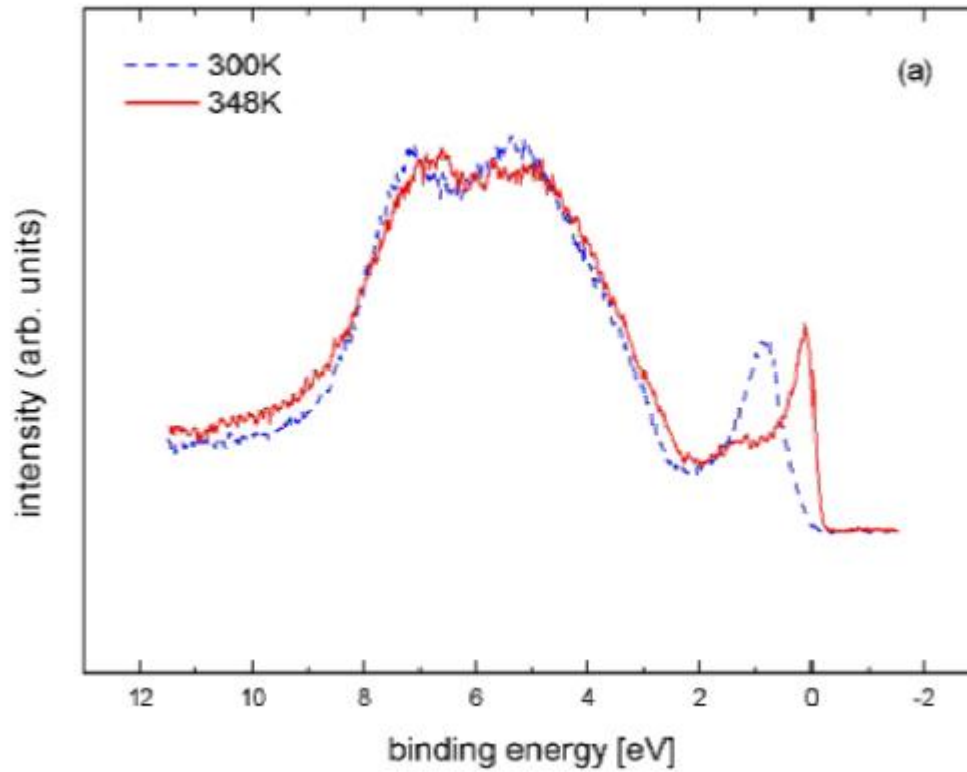


rutile

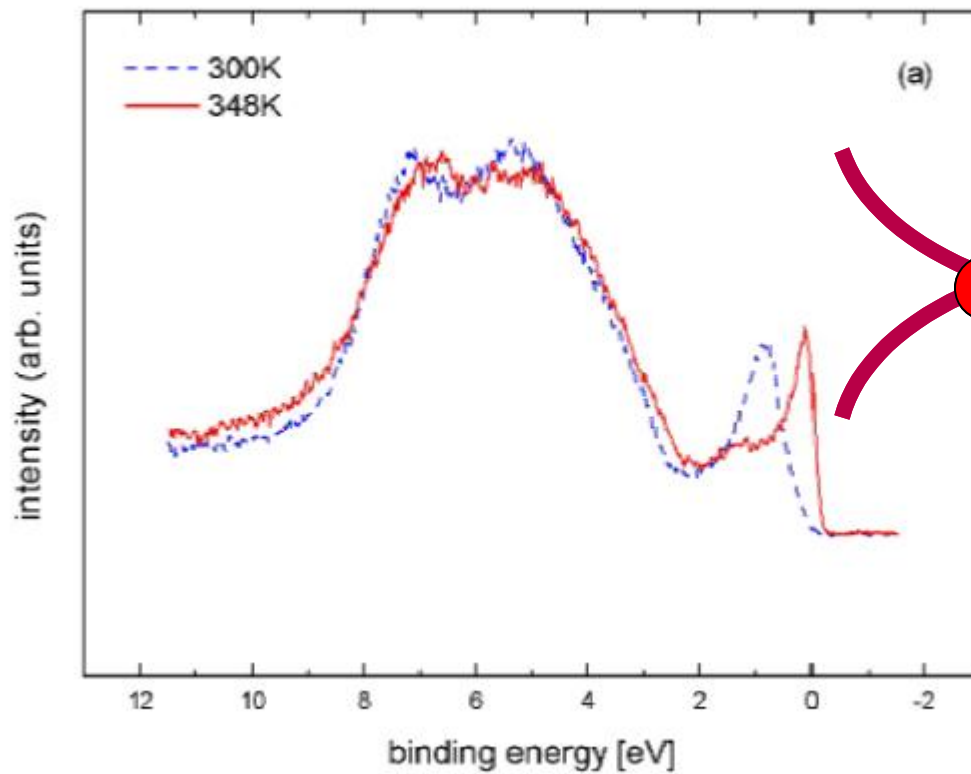


monoclinic

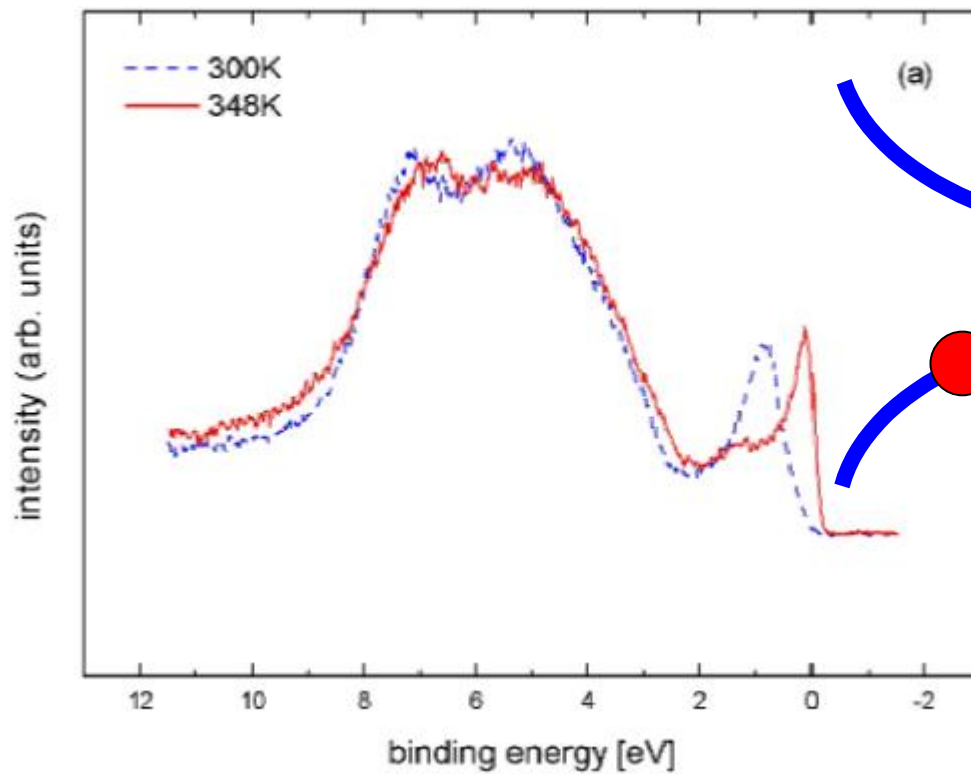
Matteo Gatti et al.



T. C. Koethe et al., Phys. Rev. Lett. 97, 116402 (2006).

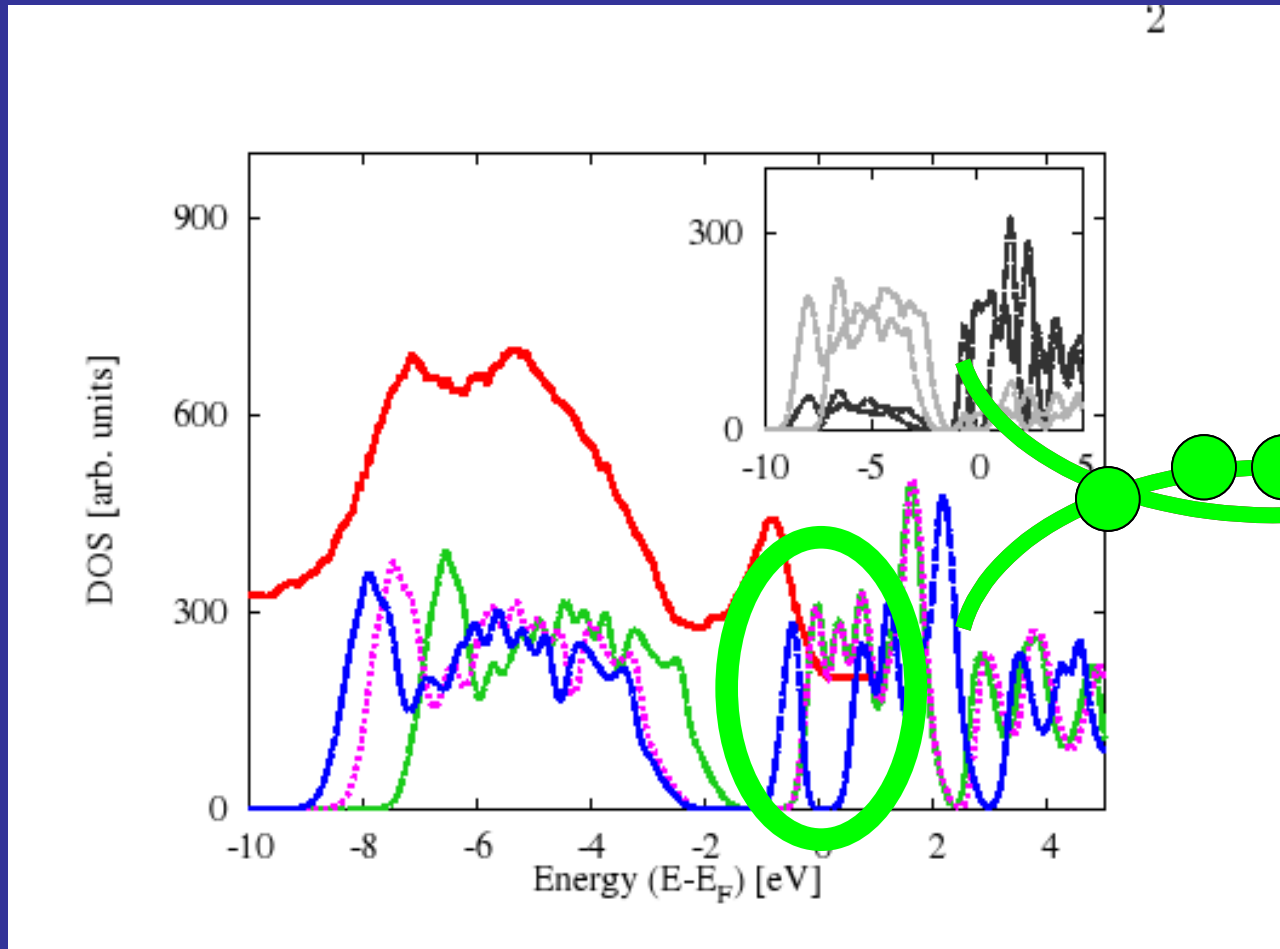


T. C. Koethe et al., Phys. Rev. Lett. 97, 116402 (2006).



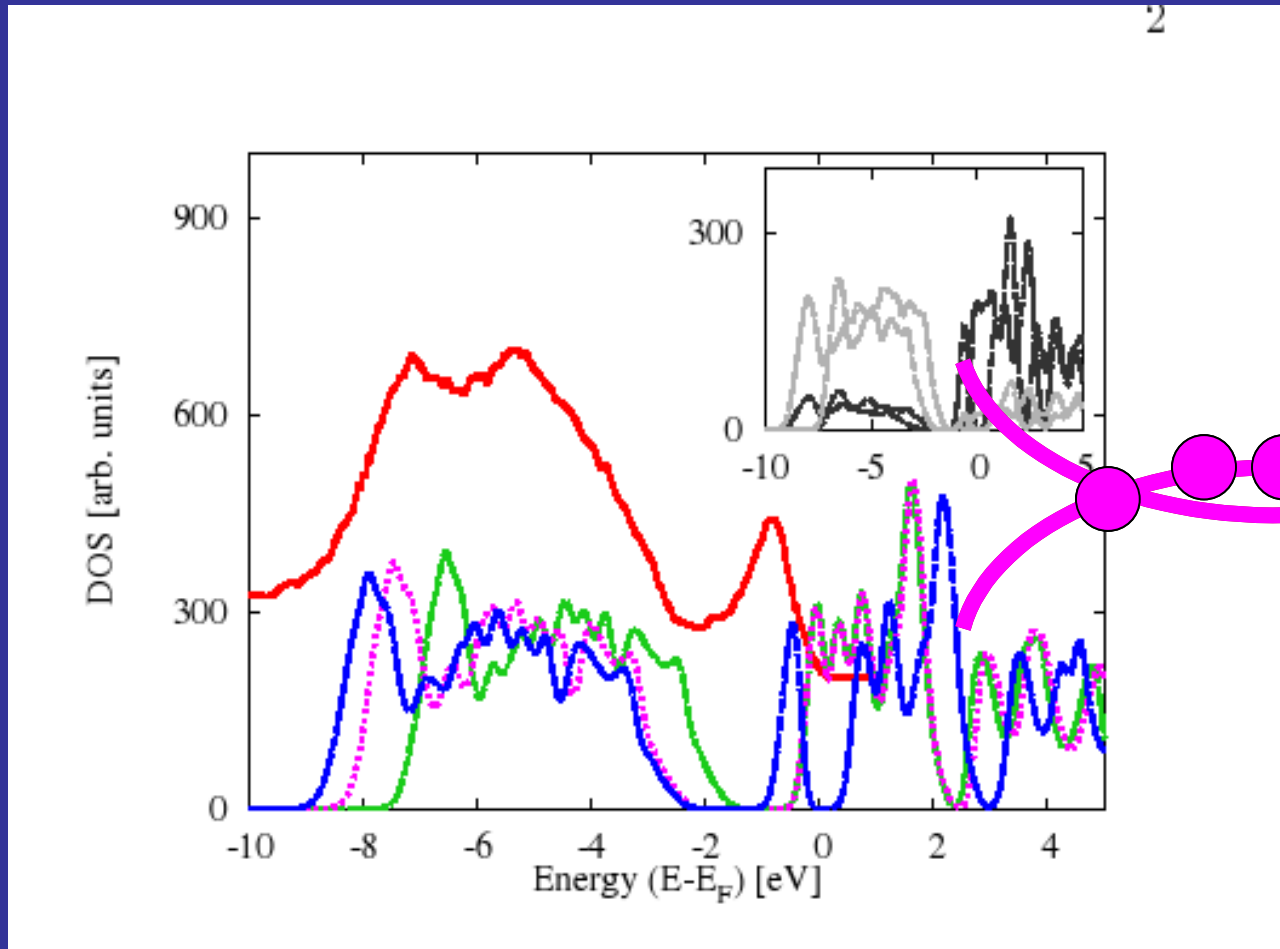
T. C. Koethe et al., Phys. Rev. Lett. 97, 116402 (2006).

LDA insulator \rightarrow metal ...



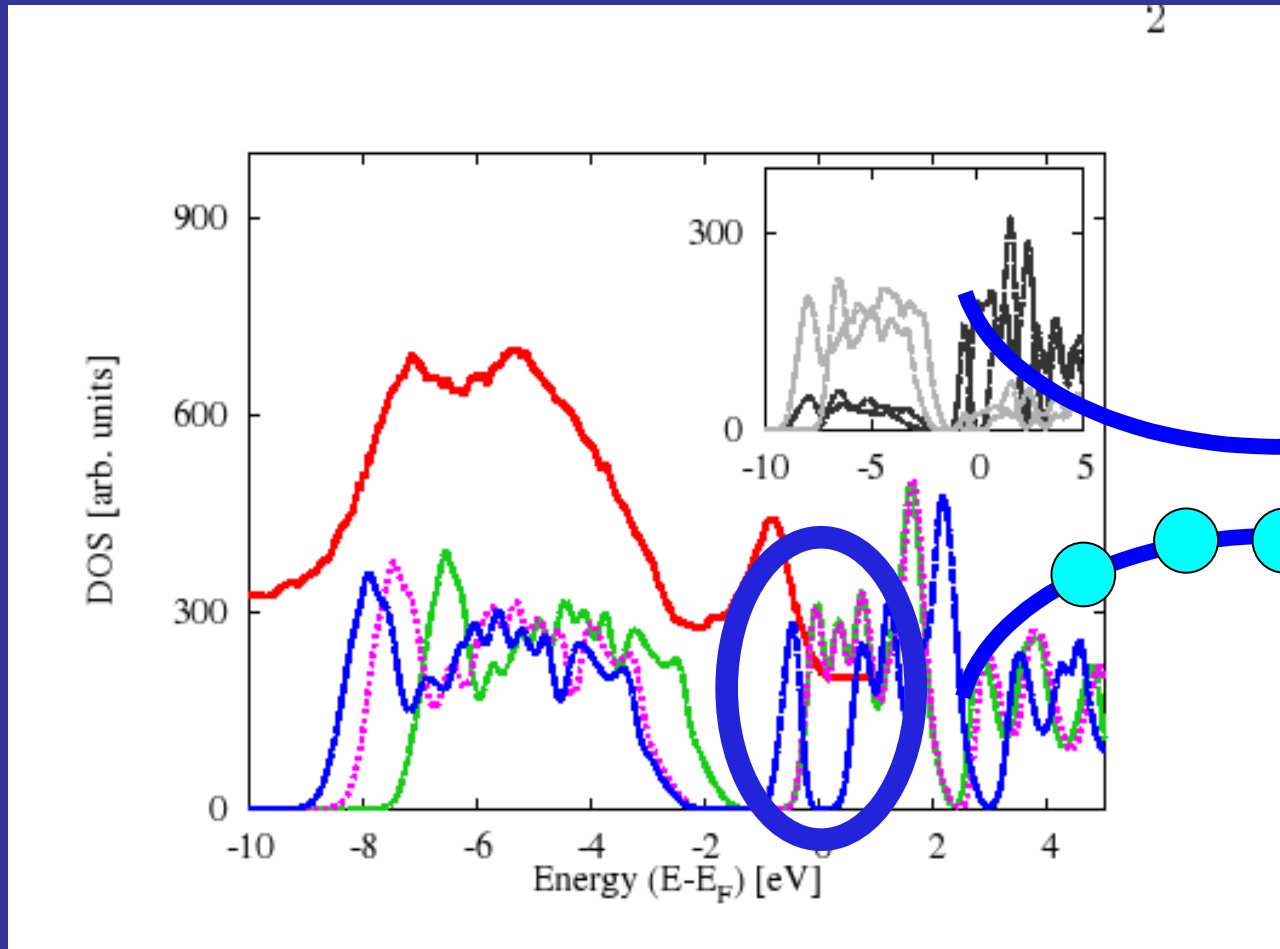
M. Gatti, F. Bruneval, V. Olevano and L. Reining,
Phys. Rev. Lett. **99**, 266402 (2007)

G_0W_0 insulator \rightarrow metal !!!!!



M. Gatti, F. Bruneval, V. Olevano and L. Reining,
Phys. Rev. Lett. **99**, 266402 (2007)

sc GW insulator = insulator !!!



M. Gatti, F. Bruneval, V. Olevano and L. Reining,
Phys. Rev. Lett. **99**, 266402 (2007)

Away from the LDA starting point!

- S. V. Faleev, M. van Schilfgaarde, and T. Kotani, PRL 93, 126406 (2004).
- Bruneval, Vast, and Reining, Phys. Rev. B 74, 045102 (2006). *COHSEX*
- T. Miyake et al., Phys. Rev. B 74, 245231 (2006). *LDA+U*
- F. Fuchs, et al., Phys. Rev. B 76, 115109 (2007). *Hybrids*
- Hong Jiang et al., , Phys. Rev. Lett. 102, 126403 (2009). *LDA+U*
- P. Rinke et al., New J. Phys. 7, 126 (2005). *KS-EXX*
-

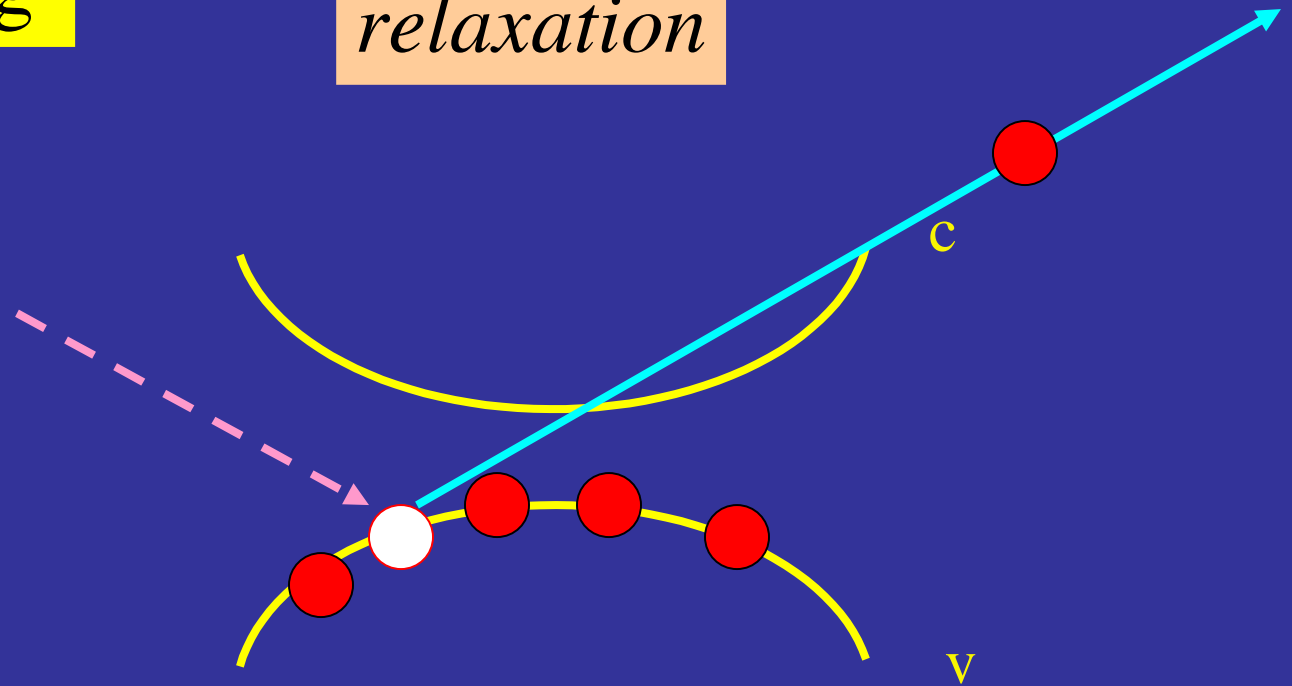
RPA: flexible screening

→ Fock exchange

localization

→ Screening

relaxation



...of practical importance: example photovoltaics

Beyond Standard GW

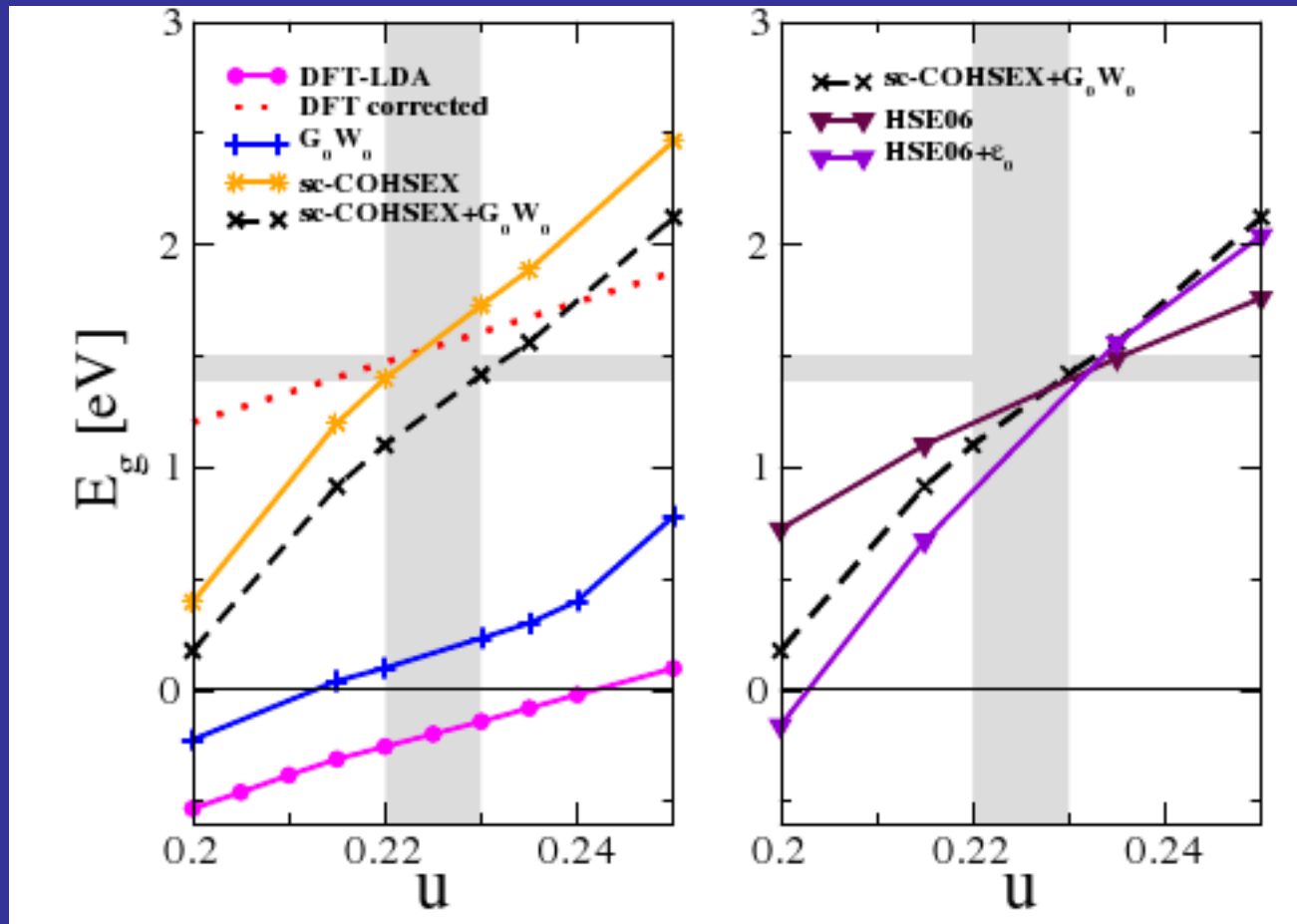
Quasiparticle energies within sc-GW for CIS

	DFT-LDA	G_0W_0	sc-GW	exp.
CuInS ₂				
E_g	-0.11	0.28	1.48	1.54
In-(S,Se)	6.5	6.9	7.0	6.9
(Se,S) s band	12.4	13.0	13.6	12.0
In 4 d band	14.6	16.4	18.2	18.2
CuInSe ₂				
	DFT-LDA	G_0W_0	sc-GW	exp.
E_g	-0.29	0.25	1.14	1.05 (+0.2)
In-(S,Se)	5.8	6.15	6.64	6.5
(Se,S) s band	12.6	12.9	13.6	13.0
In 4 d band	14.7	16.2	17.8	18.0

J. Vidal et al., PRL in press (2010)



Bandgaps in function of structure (Cu-S)



Hybrids ~ approximate GW with almost fixed screening

RPA: dynamical screening

→ Fock exchange

→ Screening

→ Screening (dynamical)

$W(\omega)$ leads to imaginary part

$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_i \frac{f^i(\mathbf{r}) f^{i*}(\mathbf{r}')}{\omega - \varepsilon_i} \quad (1)$$

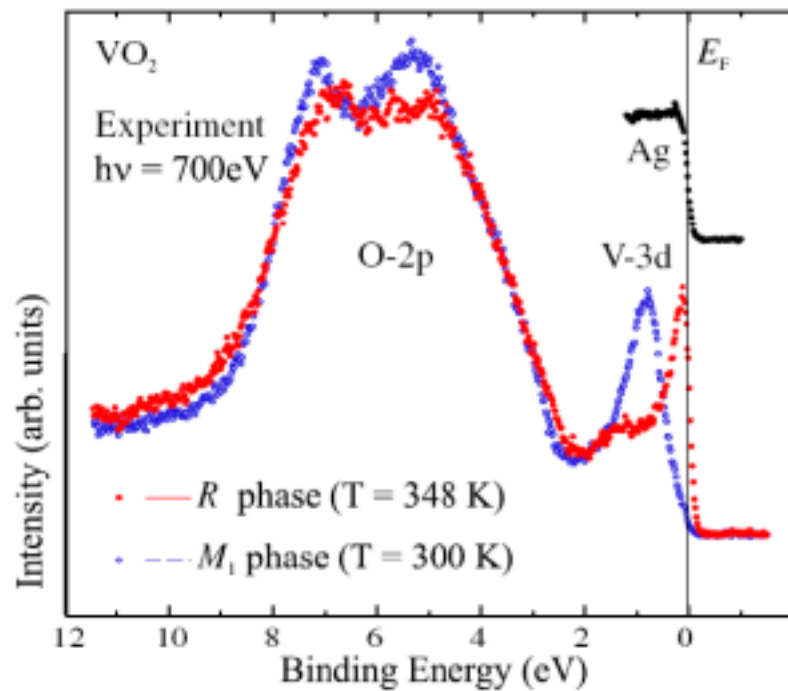
$$W(\mathbf{r}, \mathbf{r}', \omega) = \frac{1}{|\mathbf{r} - \mathbf{r}'|} + \sum_s \frac{2\omega_s V^s(\mathbf{r}) V^s(\mathbf{r}')}{\omega^2 - \omega_s^2}, \quad (2)$$

$$\langle k | \Sigma_c | k \rangle = \sum_{i,s \neq 0} \frac{|V_{ki}^s|^2}{\omega - \omega_s \text{sgn}(\mu - \varepsilon_i) - \varepsilon_i} \quad (3)$$

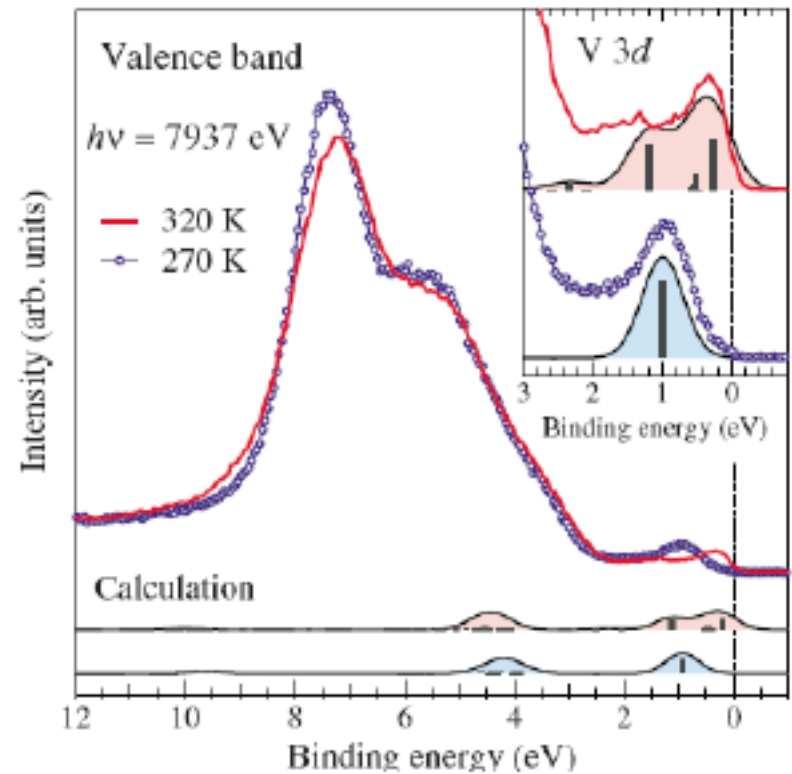
$$A_k(\omega) = \frac{1}{\pi} |Im G_{kk}(\omega)| = \frac{1}{\pi} \frac{|Im \Sigma_{kk}(\omega)|}{(\omega - \varepsilon_k - Re \Sigma_{kk}(\omega))^2 + (Im \Sigma_{kk}(\omega))^2} \quad (4)$$

Broadening (lifetime), satellites

Photoemission spectra



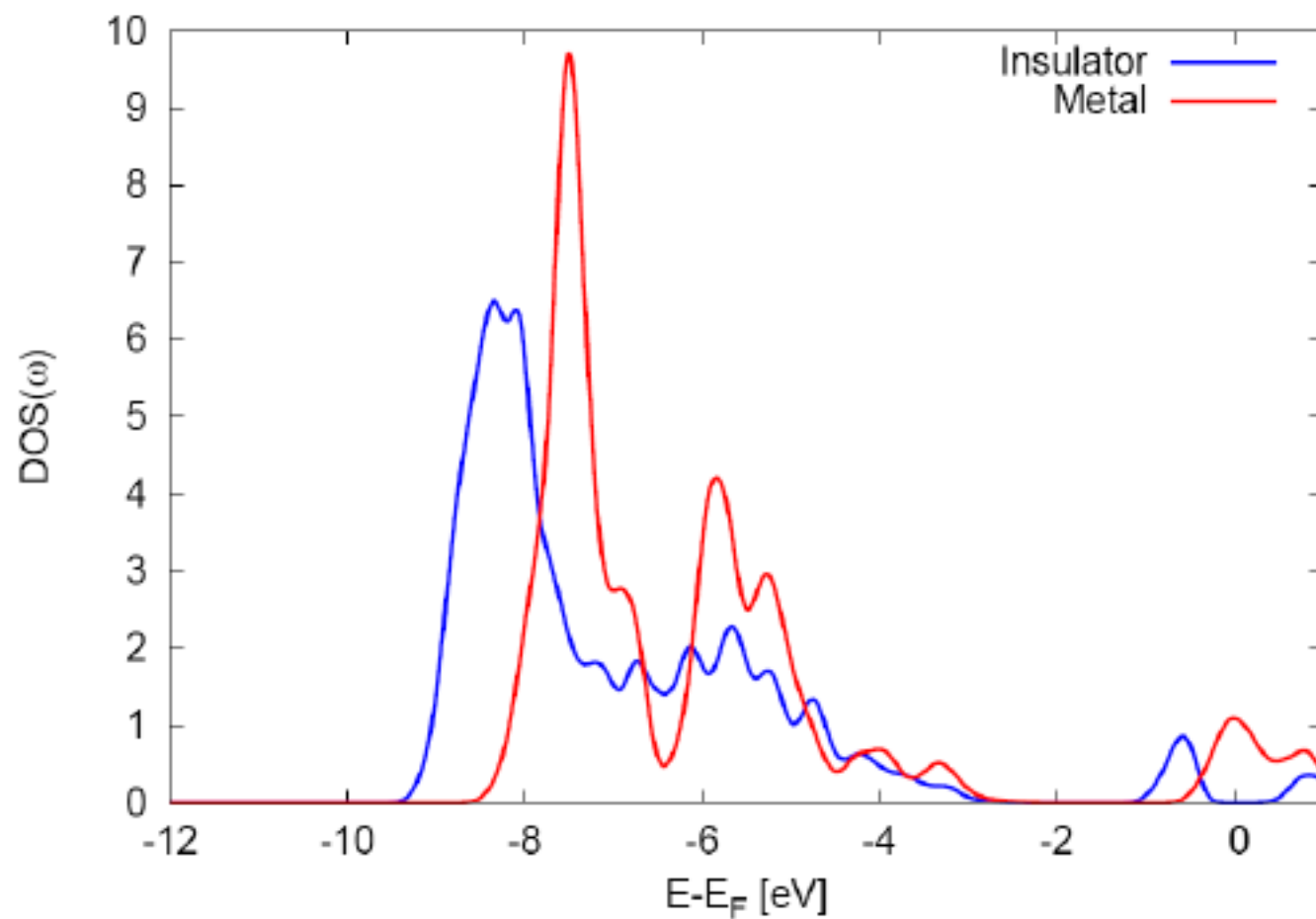
From: Koethe *et al.* PRL 97 (2006)



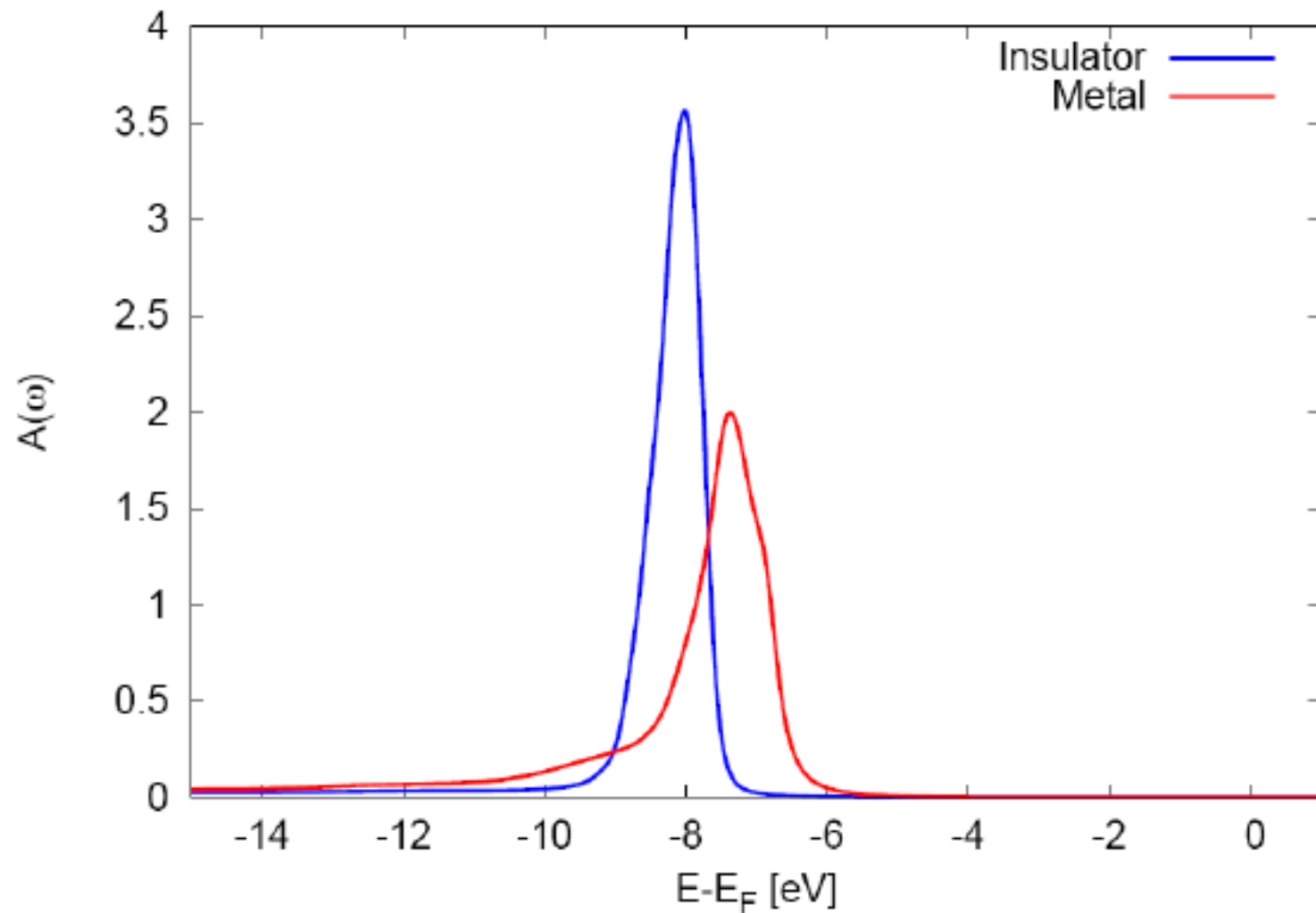
From: Eguchi *et al.* PRB 78 (2008)

Identification of the peak: partial DOS

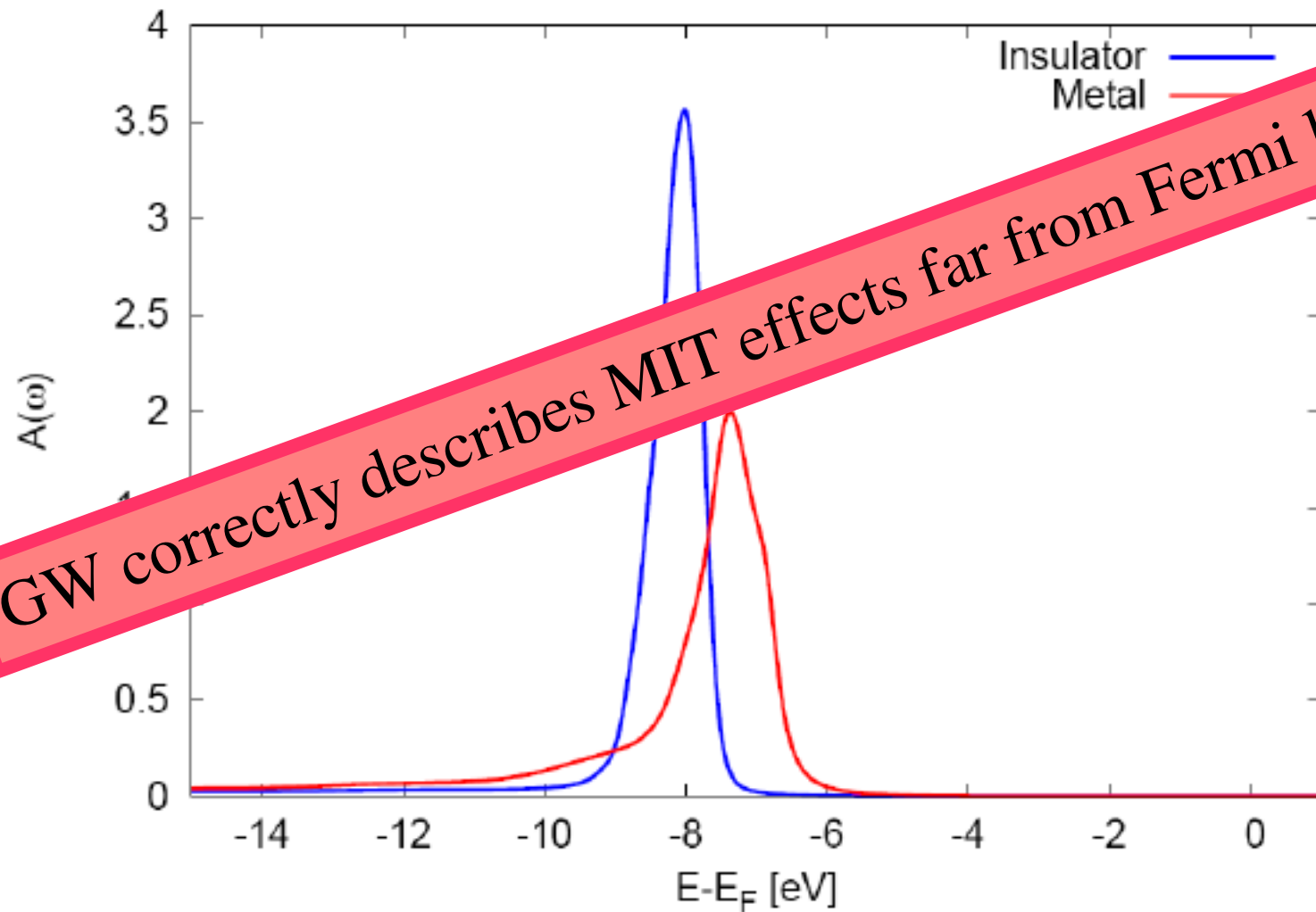
Partial DOS: V s component



GW spectral function



GW spectral function



GW correctly describes MIT effects far from Fermi level

Role of Surface Plasmons in the Decay of Image-Potential States on Silver SurfacesA. García-Lekue,¹ J. M. Pitarke,^{1,2} E. V. Chulkov,^{2,3} A. Liebsch,⁴ and P. M. Echenique^{2,3}PHYSICAL REVIEW B **76**, 195116 (2007)**Many-pole model of inelastic losses in x-ray absorption spectra**J. J. Kas,¹ A. P. Sorini,¹ M. P. Prange,¹ L. W. Cambell,² J. A. Soininen,³ and J. J. Rehr¹**Anomalous Quasiparticle Lifetime in Graphite: Band Structure Effects**Catalin D. Spataru,^{1,2} Miguel A. Cazalilla,³ Angel Rubio,^{4,5} Lorin X. Benedict,⁶

$W(\omega)$ leads to structured imaginary part

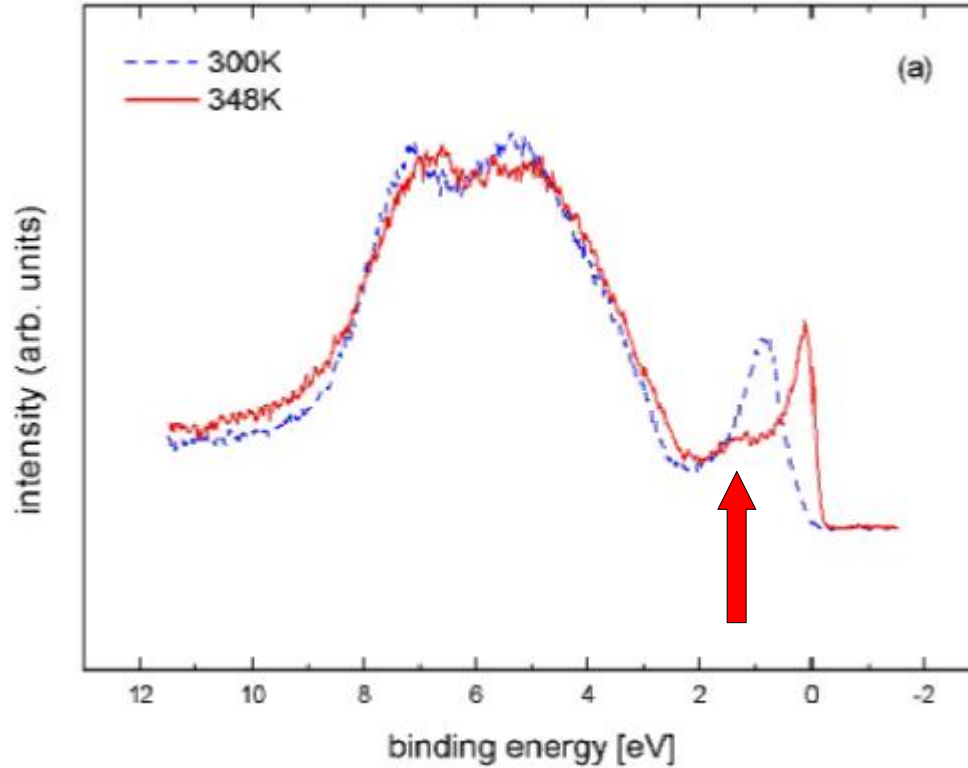
$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_i \frac{f^i(\mathbf{r}) f^{i*}(\mathbf{r}')}{\omega - \varepsilon_i} \quad (1)$$

$$W(\mathbf{r}, \mathbf{r}', \omega) = \frac{1}{|\mathbf{r} - \mathbf{r}'|} + \sum_s \frac{2\omega_s V^s(\mathbf{r}) V^s(\mathbf{r}')}{\omega^2 - \omega_s^2}, \quad (2)$$

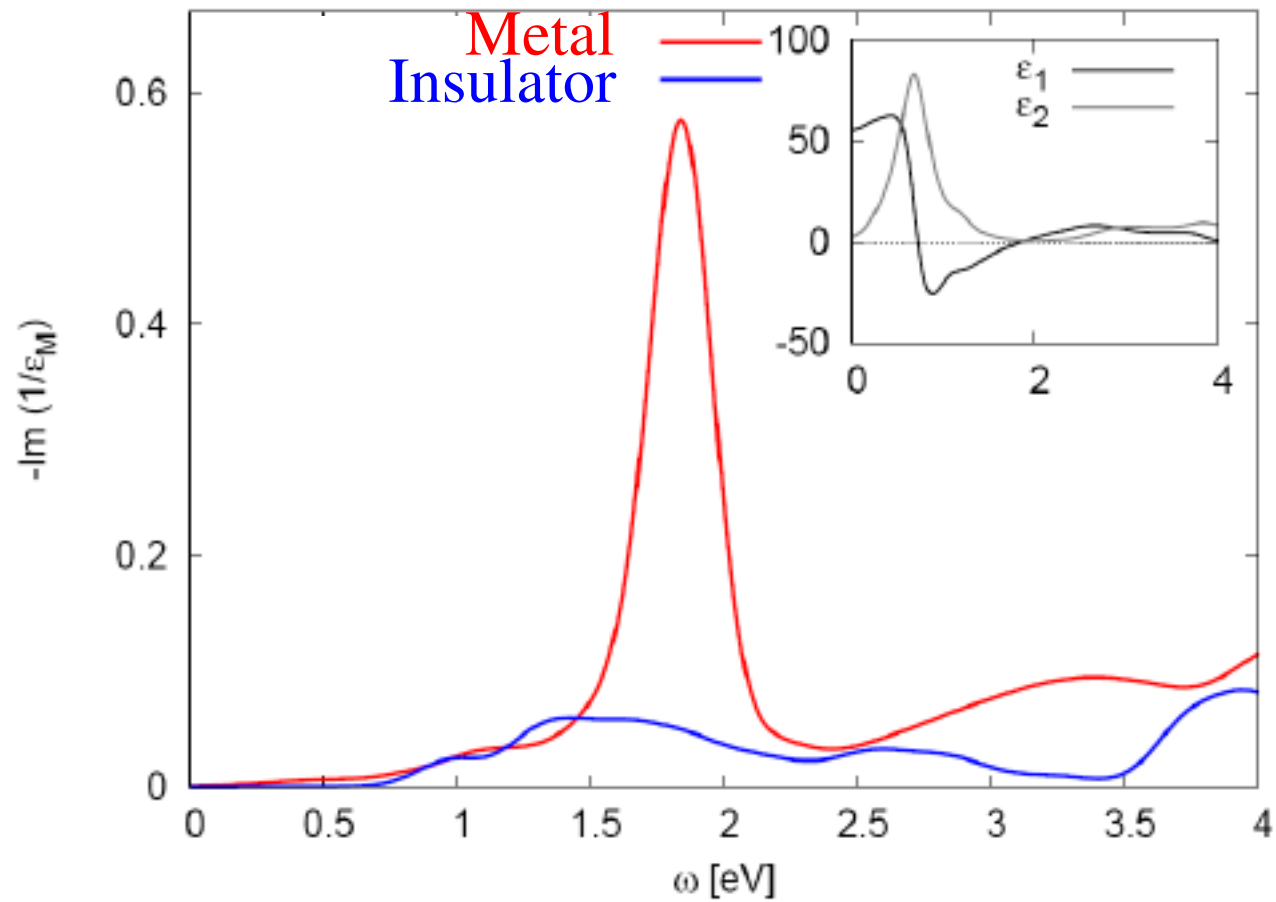
$$\langle k | \Sigma_c | k \rangle = \sum_{i,s \neq 0} \frac{|V_{ki}^s|^2}{\omega - \omega_s \text{sgn}(\mu - \varepsilon_i) - \varepsilon_i} \quad (3)$$

$$A_k(\omega) = \frac{1}{\pi} |Im G_{kk}(\omega)| = \frac{1}{\pi} \frac{|Im \Sigma_{kk}(\omega)|}{(\omega - \varepsilon_k - Re \Sigma_{kk}(\omega))^2 + (Im \Sigma_{kk}(\omega))^2} \quad (4)$$

Broadening (lifetime), satellites

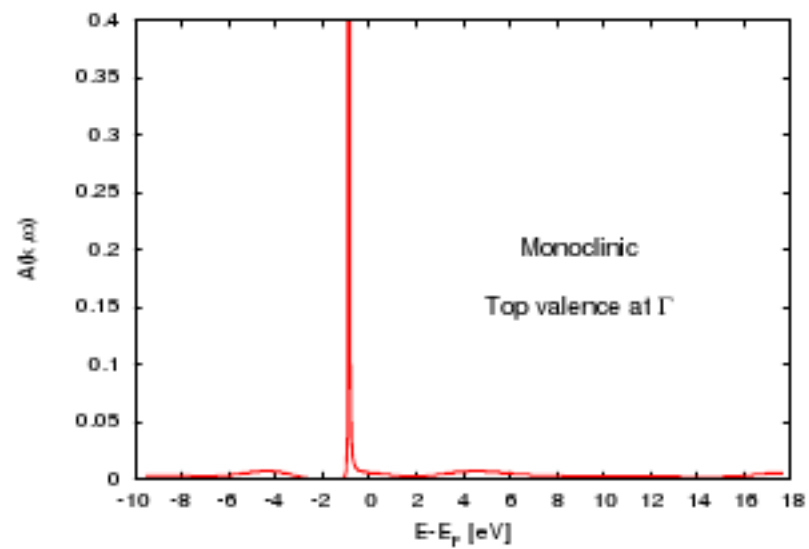
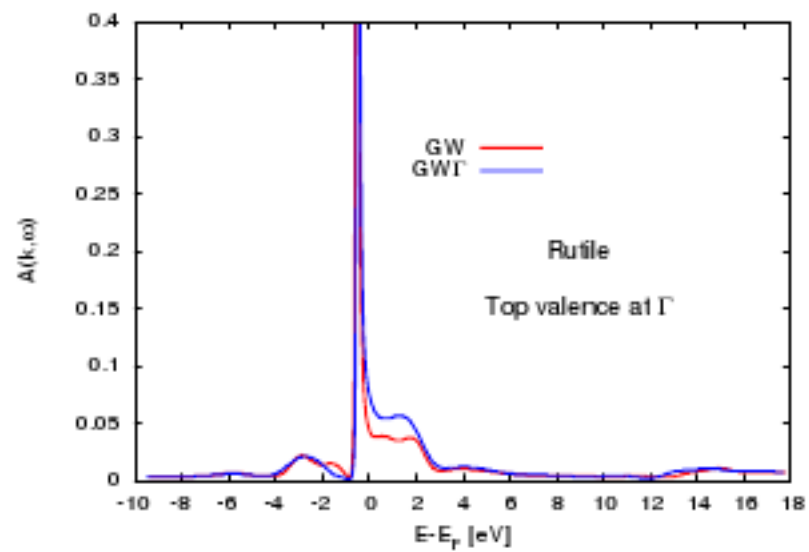


T. C. Koethe et al., Phys. Rev. Lett. 97, 116402 (2006).

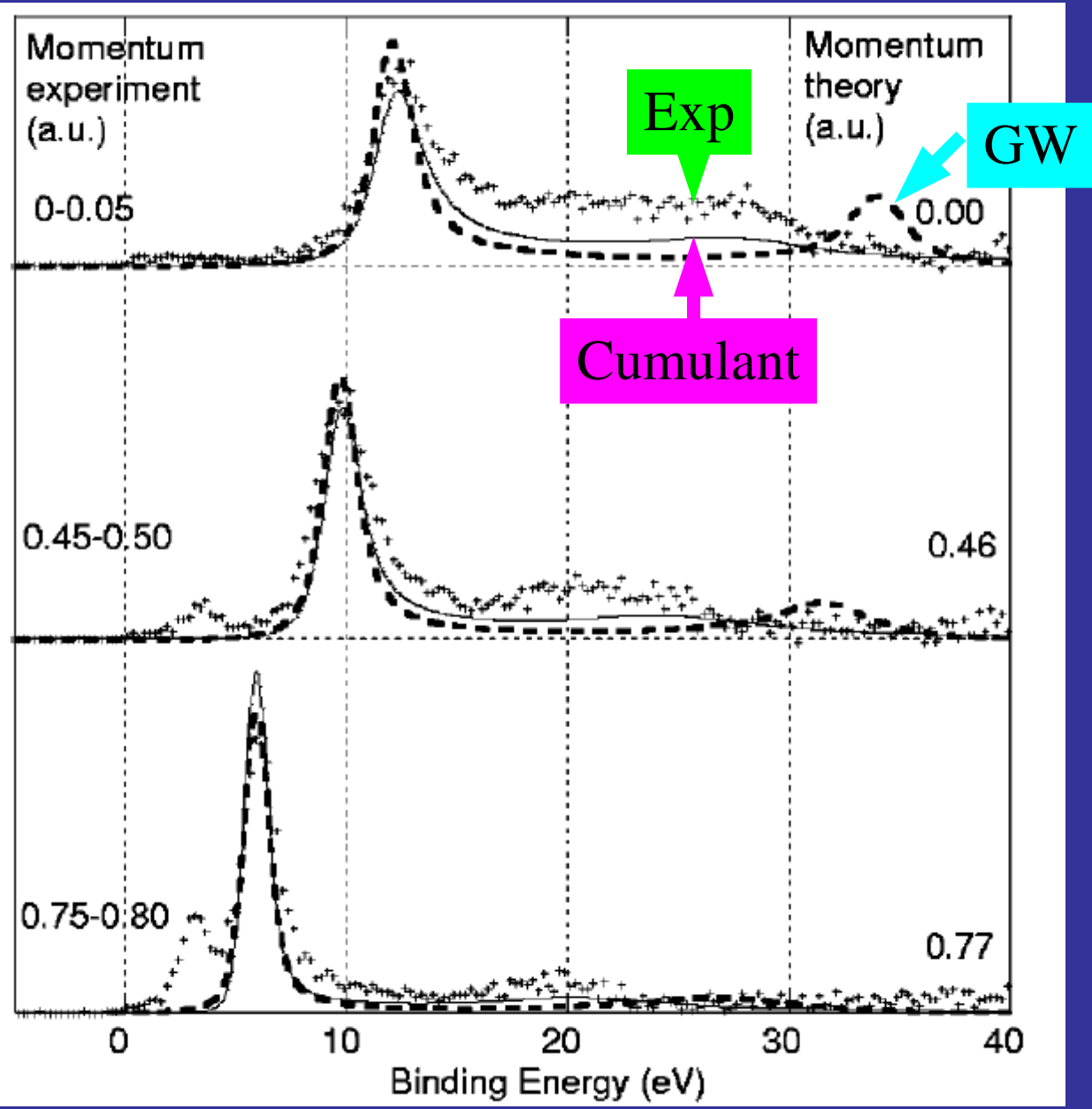


(small finite q)

see also Exp.: Abe *et al.* Jpn. J. Appl. Phys (1997)



Silicon



Satellites qualitative.

But coupling to excitations overall important

→ renormalizes energies

→ explains lifetimes

→ gives rise to van der Waals

e.g. Garcia Gonzalez + Godby, PRL 88, 056406 (2002)

→ distorts bands

e.g. Trevisanutto et al., PRL 101, 226405 (2008);

Park et al., PRL 102, 076803 (2009)

→ can be generalized (e.g. phonons)

e.g. A. Marini, PRL 101, 106405 (2008)

e.g. Eiguren + Ambrosch Draxl, PRL 101, 036402 (2008)

Is there anything **BAD** about the RPA W in GW?

→ Fock exchange

→ Screening

→ Screening (dynamical)

→ Widely valid and efficient

→ Room for speedup

$$\Sigma_{xc}(1, 2) = \Sigma_x + iG(1\bar{3})\Xi(\bar{3}, \bar{5}; 2, \bar{6})L(\bar{6}\bar{2}; \bar{5}\bar{2})v_c(\bar{2}, 1) \quad (1)$$

$$\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}') := -i\delta(\bar{3}\bar{3}')\delta(\bar{2}', \bar{2})v_c(\bar{3}\bar{2}) + \frac{\delta\Sigma_{xc}(\bar{3}, \bar{3}')}{\delta G(\bar{2}', \bar{2})} \quad (2)$$

Interaction = variation of “potential”

$$L(1, 2, 1', 2') = L_0(1, 2, 1', 2') + L_0(1, \bar{3}', 1', \bar{3})\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}')L(\bar{2}', 2, \bar{2}, 2') \quad (3)$$

What does the system do (create e-h pairs...)

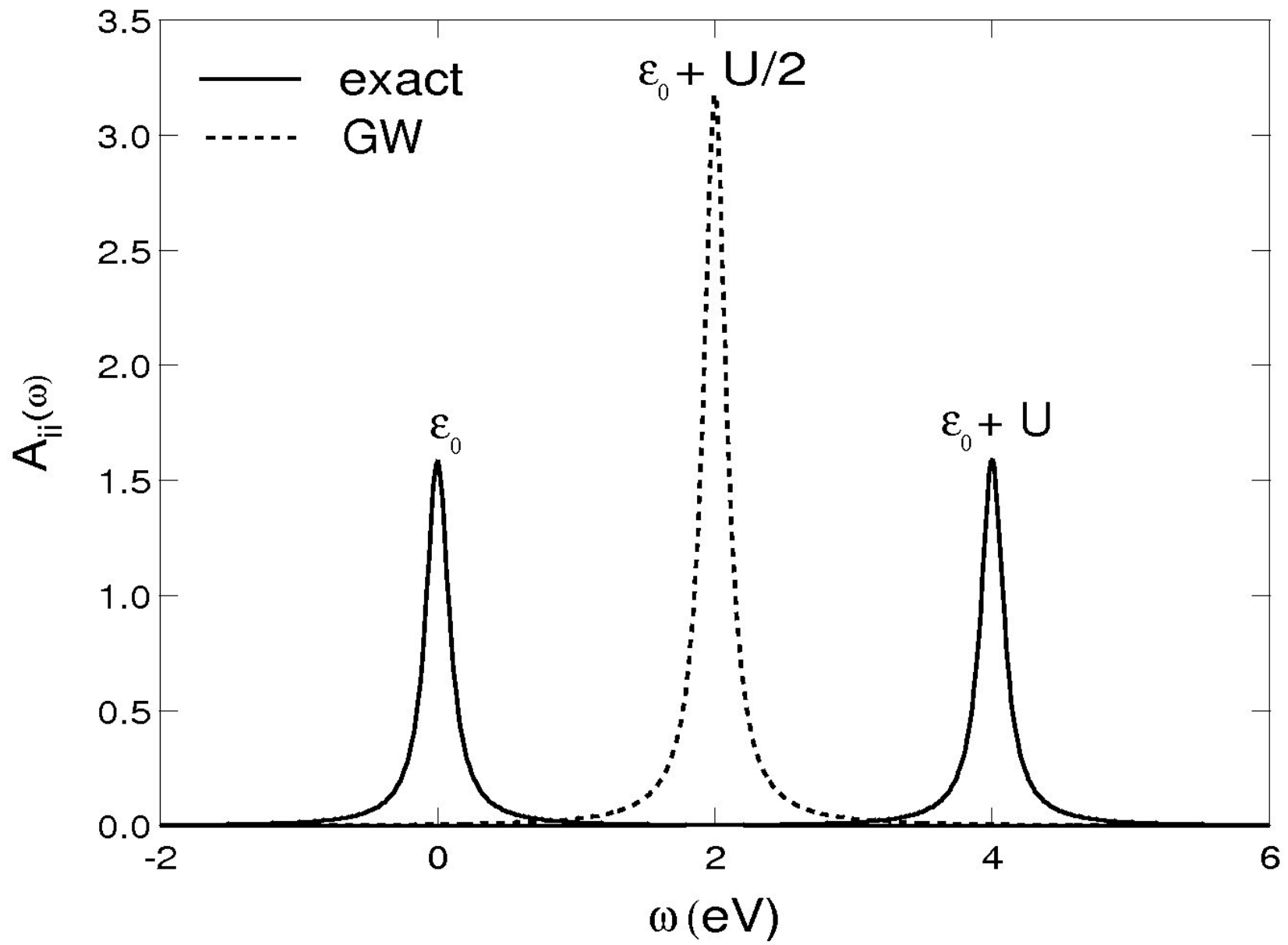
$$L_0(1, 2, 1', 2') = G(1, 2')G(2, 1') \quad (4)$$

System does nothing: HF

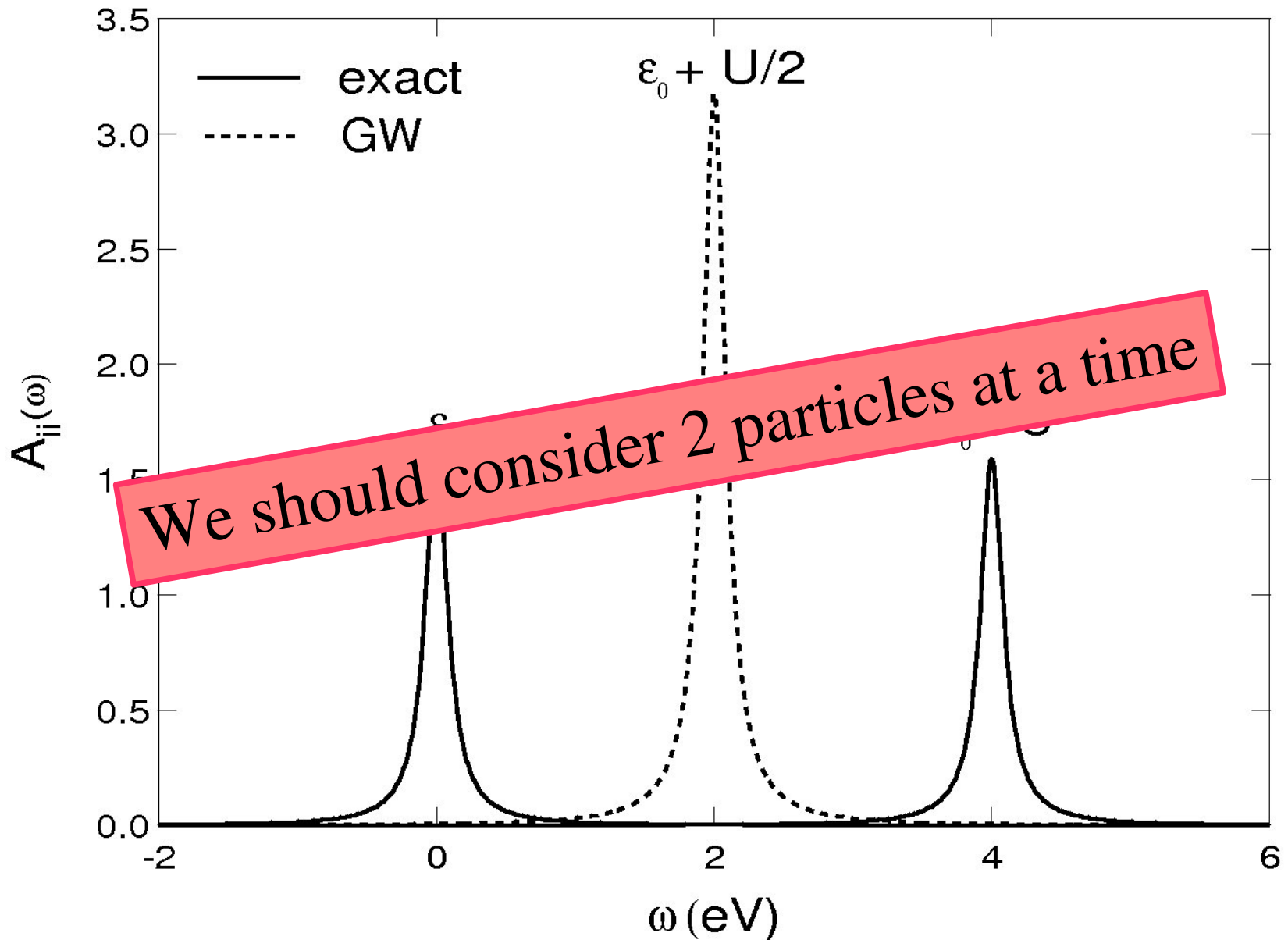
Only classical part of Ξ : $L \rightarrow i\chi$, $\Sigma_{xc} \rightarrow iGW$

Additional particle interacts in a classical way

2 site 1 electron Hubbard model; P. Romaniello et al. JPC 2009



2 site 1 electron Hubbard model; P. Romaniello et al. JPC 2009



$$\Sigma_{xc}(1, 2) = \Sigma_x + iG(1\bar{3})\Xi(\bar{3}, \bar{5}; 2, \bar{6})L(\bar{6}\bar{2}; \bar{5}\bar{2})v_c(\bar{2}, 1) \quad (1)$$

$$\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}') := -i\delta(\bar{3}\bar{3}')\delta(\bar{2}', \bar{2})v_c(\bar{3}\bar{2}) + \frac{\delta\Sigma_{xc}(\bar{3}, \bar{3}')}{\delta G(\bar{2}', \bar{2})} \quad (2)$$

Interaction = variation of “potential”

$$L(1, 2, 1', 2') = L_0(1, 2, 1', 2') + L_0(1, \bar{3}', 1', \bar{3})\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}')L(\bar{2}', 2, \bar{2}, 2') \quad (3)$$

What does the system do (create e-h pairs....)

$$L_0(1, 2, 1', 2') = G(1, 2')G(2, 1') \quad (4)$$

System does nothing: HF

Only classical part of Ξ : $L \rightarrow i\chi$, $\Sigma_{xc} \rightarrow iGW$

Better Ξ : e.g. T-matrix approach

We are used to solve 2-particle equations: BSE

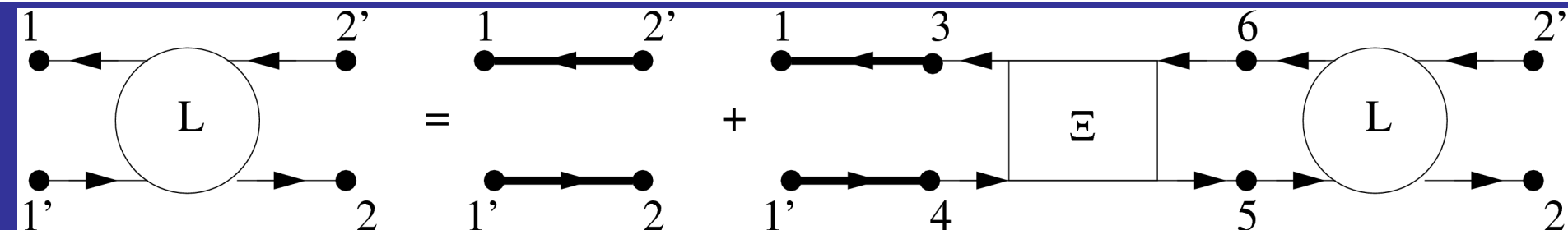
$$\Sigma_{xc}(1, 2) = \Sigma_x + iG(1\bar{3})\Xi(\bar{3}, \bar{5}; 2, \bar{6})L(\bar{6}\bar{2}; \bar{5}\bar{2})v_c(1, \bar{2}) \quad (1)$$

$$\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}') := -i\delta(\bar{3}\bar{3}')\delta(\bar{2}', \bar{2})v_c(\bar{3}\bar{2}) + \frac{\delta\Sigma_{xc}(\bar{3}, \bar{3}')}{\delta G(\bar{2}', \bar{2})} \quad (2)$$

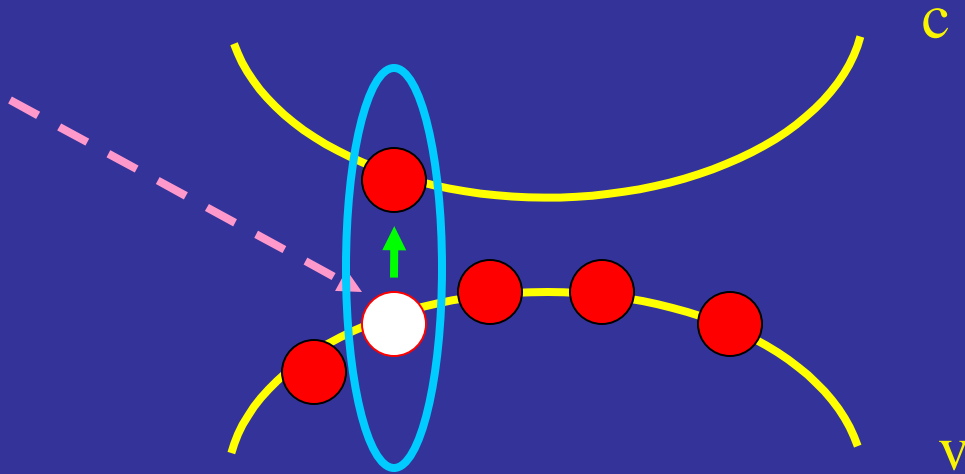
$$L(1, 2, 1', 2') = L_0(1, 2, 1', 2') + L_0(1, \bar{3}', 1', \bar{3})\Xi(\bar{3}, \bar{2}, \bar{3}', \bar{2}')L(\bar{2}', 2, \bar{2}, 2') \quad (3)$$

Dyson-like equation: BSE

$$L_0(1, 2, 1', 2') = G(1, 2')G(2, 1') \quad (4)$$



Absorption ?



Electron-hole interaction

Excitonic effects

Bethe-Salpeter Equation

Like in GW:

→ Fock exchange

Makes the electron-hole attraction

→ Screening

Calculated consistently

→ Screening (dynamical)

Makes more excitations

→ Widely valid and efficient

Many people

→ Room for speedup

*Wannier functions,
propagation, mixed TDDFT,...*

4. Look at the total energy

Adiabatic connection fluctuation dissipation theorem
(ACFDT)

$$E_X^{\text{EXX}}[n] = -2 \times \frac{1}{2} \sum_{nm}^{\text{occ}} \int d^3\mathbf{r} d^3\mathbf{r}' \frac{c_{nm}(\mathbf{r}) c_{mn}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

$$c_{nm}(\mathbf{r}) = \phi_n^*(\mathbf{r}) \phi_m(\mathbf{r})$$

$$E_C[n] = - \int_0^{\infty} \frac{du}{2\pi} \int_0^1 d\lambda \int d^3\mathbf{r} d^3\mathbf{r}' \frac{1}{|\mathbf{r} - \mathbf{r}'|} \times \\ [\chi_{\lambda}(\mathbf{r}, \mathbf{r}'; iu) - \chi_0(\mathbf{r}, \mathbf{r}'; iu)]$$

.....in DFT framework

Total energy in GW framework

$$\begin{aligned}\Phi_{GW} &= \frac{1}{2}Tr [GV_H] + \frac{1}{2}Tr [G\Sigma_x] - Tr \left[\frac{1}{4}(v_cGG)^2 + \frac{1}{6}(v_cGG)^3 + \dots \right] \\ &= \frac{1}{2}Tr [GV_H] + \frac{1}{2}Tr [G\Sigma_x] + \frac{1}{2}Tr [v_cGG + \ln(1 - v_cGG)].\end{aligned}$$

$$\Omega_{Klein}[G_s] = \Phi[G_s] - Tr(G_0^{-1}G_s - 1) - Tr \ln(-G_s^{-1})$$

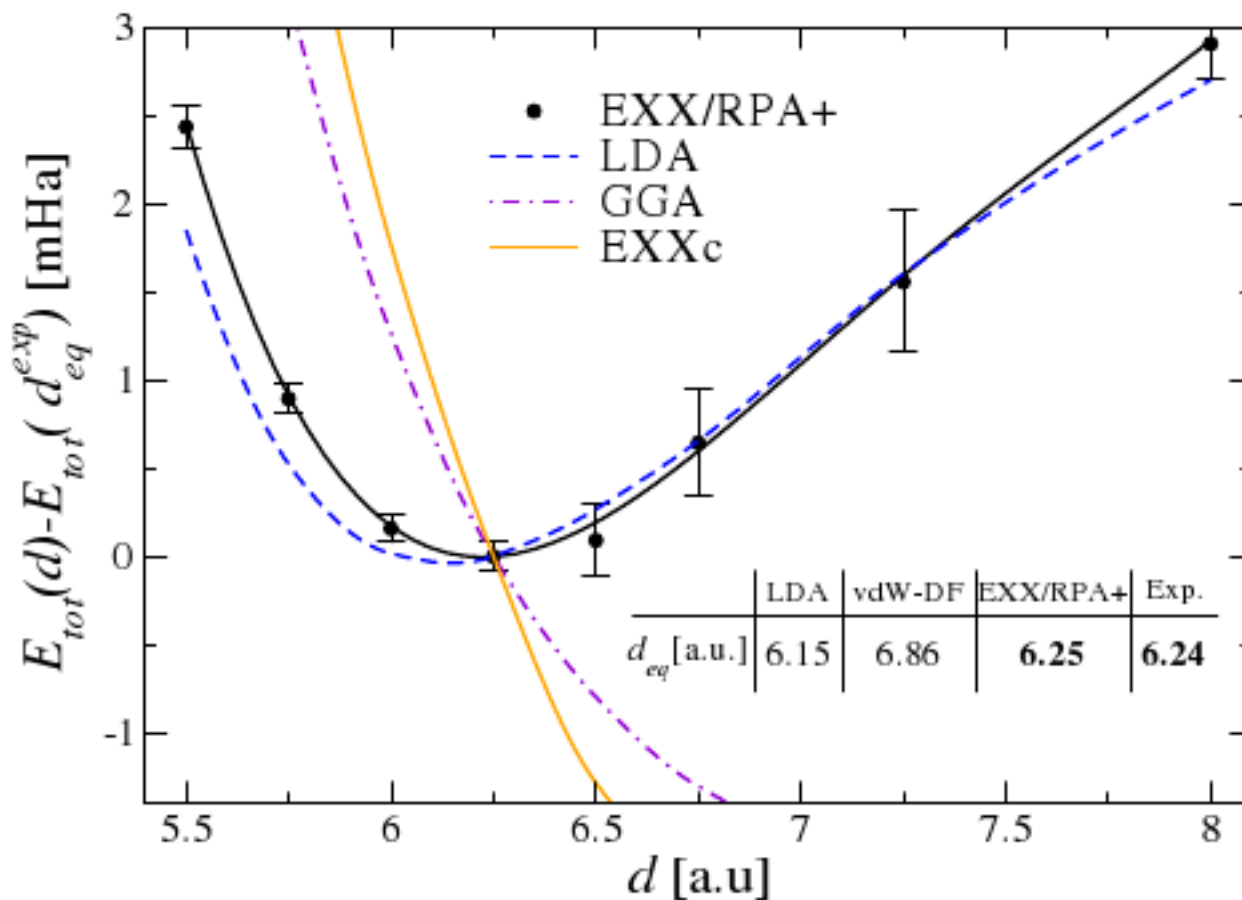
$$E^{GWA,s} = \Omega_{Klein}[G_s] + \mu N =$$

$$= \langle \Psi_s | \hat{H} | \Psi_s \rangle + \frac{1}{2}Tr [v_cG_sG_s + \ln(1 - v_cG_sG_s)]$$

.....evaluated at $G=G_s$, equivalent to RPA ACFDT

First-Principles Description of Correlation Effects in Layered Materials

Andrea Marini,¹ P. García-González,² and Angel Rubio^{3,4}



RPA+ : because RPA has deficiencies

Uniform gas correlation energy too negative.

Can be corrected by LDA like term.

S. Kurth and J. P. Perdew, Phys. Rev. B 59, 10 461
(1999); M. Lein, E. K. U. Gross, and J. P. Perdew,
Phys. Rev. B 61, 13 431 (2000).

Conclusions

RPA = simplest way to include response
=== a lot of physics!!!

Spectra: plasmons, crystal LFE, but e.g. no excitons
(Hartree only!!!)

Many-body effects:

renormalization of energies,
lifetimes
plasmon satellites

but no “strong correlation”

Total energies: long range contributions, improved correlation energies

Conclusions

RPA = simplest way to include response
=== a lot of physics!!!

Spectra: plasmons, crystal LFE, but e.g. no excitons

RPA = excellent starting point for doing better

Many-body effects:

- renormalization of energies,
- lifetimes
- plasmon satellites

but no “strong correlation”

Total energies: long range contributions, improved correlation energies

Palaiseau Theoretical Spectroscopy Group



Thank you!!!