#### Part 2. Ab initio Valence Bond

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navigation Main page CTTC 2014 Nha Trang VB tutorial center Recent changes Random page	page       discussion       edit       history       delete       move       pro         VB tutorial	An Ab Initio Non-orthogon	al Valence Bond Program	Benoit talk preferences	watchlist contri	ibutions log o	X
= Help search Search Go Search	Four tutorials ("hands-on" labs) sessions are planned, where partic concepts on some chemical problems of progressive difficulty. Pa book exercises. Rooms will also be at the disposal of participants	cipants will learn to use the X art of the «hands-on» session during the «hands-on» time f	MVB P program, as well as the BLW s can as well be used by participants or free discussions.	P and HuLis P programs, for questions/answers co	and apply the r	methods and ain lectures o	i or
toolbox What links here Related changes Upload file Special pages Printable version Permanent link Page information	Contents [hide] 1 The XMVB program 2 Tutorials 2.1 Tutorial n°1 : Basics of VB theory and XMVB program 2.2 Tutorial n°2 : VB applications on PI systems 2.3 Tutorial n°2 : State correlation diagrams 2.4 Tutorial n°3 : State correlation diagrams 3 The Valence Bond mailing list 4 Literature 5 VB lectures in Roscoff						
	The XMVB program					[6	edit]
	XMVB is an <i>ab initio</i> non-orthogonal Valence Bond program. It is here .	available upon request 🖾. Its	manual could be downloaded here D,	and a detailed presentation	on of input/outp	out structure	
	Tutorials					[0	edit]
	You can dowload all tutorial exercises from this link						

#### https://wiki.lct.jussieu.fr/workshop/index.php/VB\_tutorial



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## Ab initio Valence Bond

- Electron correlation
- VB methods including electron correlation
- VB computation in practice
- Illustrative application
- Limits of VB theory

## Stop me at any time !

- Nature of electron correlation :
  - «Uncorrelated» ( $\Psi_{HF}$  or  $\Psi_{HL}$ ) :

Each electron feels the **mean field** created by all the other electrons



• Correlated :

Each electron **dynamically** affects the configuration of all the other electrons



• «Uncorrelated» methods :  $\hat{H} = \hat{T} + \hat{V}$ 

 $\Psi_{HF}$  too much ionic :



Methods including electron correlation :



E (kcal/mole)



Including electron
 correlation is a tough problem,
 as we are dealing with very
 tiny energy differences

- The VBSCF method
- The **BOVB** method
- The VBCI method

- The VBSCF\* method :
  - Basically a MCSCF method with nonorthogonal orbitals :

$$\Psi_{VBSCF} = \sum_{K} C_{K} \Phi_{K} \quad \text{with} : \begin{cases} \Phi_{K}(1,...,N) = \hat{A} \{ \prod_{i=1}^{N} \varphi_{i}(1) \Theta_{K} \} : \text{VB structures} \\ \{\varphi_{i}\} : \text{set of non-orthogonal localized orbitals} \\ \text{expanded onto a set of basis functions} \{\chi_{m}\} : \varphi_{i}(1) = \sum_{m} d_{m}^{i} \chi_{m}(1) \\ \Theta_{K} \text{ spin function} \end{cases}$$

All w.f. parameters : structure coef. { $C_{\rm K}$ } and orb. coef. { $d^{\rm i}_{\rm m}$ } are **optimized simultaneously** minimizing  $\langle \Psi_{VBSCF} | \hat{H} | \Psi_{VBSCF} \rangle$ 

\*van Lenthe; Balint-Kurti, J. Chem. Phys. 1983, 78, 5699

#### • The VBSCF method :



The VBSCF method ensures a correct balance beween covalent and ionic configurations («left-right» static correlation)

- The VBSCF method :
  - Comparison with CASSCF method :

$$\Psi_{\text{CASSCF}} = \lambda \left( \begin{array}{c} \varphi_{a} \\ \varphi_{a} \\ \varphi_{g} \end{array} \right) + \mu \left( \begin{array}{c} \varphi_{a} \\ \varphi_{a} \\ \varphi_{g} \end{array} \right) + \mu \left( \begin{array}{c} \varphi_{a} \\ \varphi_{b} \\ \varphi_{g} \end{array} \right) + \mu \left( \begin{array}{c} \varphi_{a} \\ \varphi_{a} \\ \varphi_{g} \end{array} \right) + \left( \begin{array}{c} \varphi_{a} \\ \varphi_{a} \\ \varphi_{b} \\ \varphi_{a} \\ \varphi_{b} \\ \varphi_{a} \\ \varphi_{a} \\ \varphi_{a} \\ \varphi_{a} \\ \varphi_{a} \\ \varphi_{a} \\ \varphi_{b} \\$$

- The VBSCF method :
  - Accuracy :  $\mathbf{F}-\mathbf{F} \xrightarrow{\mathbf{\Delta E}} \mathbf{F} \cdot \mathbf{$

Some important physical ingredient is missing...

- The VBSCF method :
  - What the VBSCF method does :



Same set of orbitals for all VB structures : optimized for a mean situation

• A better wave function would be :

Each structure has its own specific set of orbitals

#### • The **BOVB**\* method :

• BOVB uses :

same number of structures as VBSCF, but different orbitals for the different structures :

$$\Psi_{BOVB} = \sum_{K} C_{K} \Phi_{K}^{BOVB} \text{ with } : \begin{cases} \Phi_{K}^{BOVB}(1,...,N) = \hat{A} \{\prod_{i=1}^{N} \varphi_{i}^{K}(1)\Theta_{K}\} : \text{VB structures} \\ \{\varphi_{i}^{K}\} : \text{ set of non-orthogonal (del)localized orbitals for the structure K} \\ \text{expanded onto a set of basis functions} \{\chi_{m}\} : \varphi_{i}^{K}(1) = \sum_{m} d_{m}^{i,K} \chi_{m}(1) \\ \Theta_{K} \text{ spin function} \end{cases}$$

All w.f. parameters : structure coef. { $C_{\rm K}$ } and orb. coef. sets { $d_{\rm m}^{\rm i}$ }<sup>K</sup> are optimized simultaneously minimizing  $\langle \Psi_{BOVB} | \hat{H} | \Psi_{BOVB} \rangle$ 

\* Hiberty, P. C. ; Humbel, S. ; Byrman, C. P. ; van Lenthe J. H. J. Chem. Phys. 1994, 101, 5969

#### • The **BOVB** method :

• BOVB brings that part of dynamic correlation that varies during a reaction, i.e. the **differential electron correlation** :



- BOVB keeps the same **compacity** as the VBSCF wave function
- BOVB provides a visual image of the role of electron correlation

The **BOVB** method :

• L-BOVB : EFOF C All orbitals are localized, ionics are closed-shell Spectator orbitals are delocalized in all structures • **SD-BOVB** - Active orbitals are split in ionics - Spectator orbitals are

delocalized in all structures

The **BOVB** method :

• L-BOVB : • E-BOVB : • E-B All orbitals are localized, ionics are closed-shell Spectator orbitals are delocalized in all structures • **SD-BOVB** EF • F E 

- Active orbitals are split in ionics - Spectator orbitals are delocalized in all structures

- The **BOVB** method :
  - Accuracy :  $F-F \longrightarrow F \cdot + F \cdot$

	ΔΕ	
RHF	$\frac{(\text{kcal.mol}^{-1})}{-37}$	
VBSCF	15	
L-BOVB	28.2	
SD-BOVB	33.6	
Exact	+39	∫ inco

Basis set incompleteness

• The VBCI\* method :

$$\Psi_{VBSCF} = \sum_{K} C_{K} \Phi_{K}^{0} \qquad \qquad \Psi_{VBCI} = \sum_{K} C_{K} \Phi_{K}^{0} + \sum_{K} (\sum_{i} C_{K}^{i} \Phi_{K}^{i})$$

- Where  $\Phi^{0}_{K}$  are the fundamental structures (obtained from VBSCF)

- and  $\Phi^{i}_{K}$  are excited structures, built from  $\Phi^{0}_{K}$  by replacing an occupied by a virtual orbital, but having the same physical meaning as  $\Phi^{0}_{K}$ :

$$\Phi_K^{VBCI} = C_K \Phi_K^0 + \sum_i C_K^i \Phi_K^i$$

The VBCI wave function is (much) larger than the VBSCF/BOVB one, but the interpretability is kept

\* Wu W., Song L., Cao Z., Zhang Q., and Shaik S. J. Phys. Chem. A, 2002, 106 (11), 2721-2726

• The VBCI method :

$$\Psi_{VBSCF} = \sum_{K} C_{K} \Phi_{K}^{0} \qquad \qquad \Psi_{VBCI} = \sum_{K} C_{K} \Phi_{K}^{0} + \sum_{K} (\sum_{i} C_{K}^{i} \Phi_{K}^{i})$$

- Occupied orbitals are obtained from a VBSCF calculation :

 $\varphi_i^A(1) = \sum_m d_m^{i,A} \chi_m^A(1)$  where A is a molecular fragment

 Localized virtual orbitals are generated from a Schmidt orthogonalization procedure to occupied orbitals on their own fragments :

$$\chi'_{m}^{A} = (\chi_{m} - S_{mn}T_{v}^{m}S^{\mu\nu}\varphi_{\mu})|_{m,n\in A}$$

 $\Rightarrow \Phi^{i}{}_{K}$  keep the same physical meaning as  $\Phi^{0}{}_{K}$ 

• The VBCI method :

$$\Psi_{VBSCF} = \sum_{K} C_{K} \Phi_{K}^{0} \qquad \qquad \Psi_{VBCI} = \sum_{K} C_{K} \Phi_{K}^{0} + \sum_{K} (\sum_{i} C_{K}^{i} \Phi_{K}^{i})$$

- The energy is obtained by :

$$E^{VBCI} = \frac{\left\langle \Psi^{VBCI} \middle| H \middle| \Psi^{VBCI} \right\rangle}{\left\langle \Psi^{VBCI} \middle| \Psi^{VBCI} \right\rangle} = \frac{\sum_{K,L} \sum_{i,j} C^{i}_{K} C^{j}_{L} \left\langle \Phi^{i}_{K} \middle| H \middle| \Phi^{j}_{L} \right\rangle}{\sum_{K,L} \sum_{i,j} C^{i}_{K} C^{j}_{L} \left\langle \Phi^{i}_{K} \middle| \Phi^{j}_{L} \right\rangle}$$

- A Davidson type of correction to the energy also exists :

$$\Delta E_{Q} = (1 - \sum_{K} W_{K}) \Delta E_{D} \text{ with } : W_{K} = \sum_{L} \sum_{i,j} C_{K}^{i} C_{L}^{j} \left\langle \Phi_{K}^{i} \middle| \Phi_{L}^{j} \right\rangle$$

(estimate the contribution of quadruple excitations that are product of double excitations)

• **Bond energies** (kcal.mol<sup>-1</sup>) with various methods :

molecule	$D^{ m HF}$	D <sup>B3LYP</sup>	$D^{\mathrm{CCSD}}$	DVBSCF	$D^{ ext{L-BOVB}}$	$D^{\mathrm{VBCIS}}$	$D^{\mathrm{VBCISD}}$
H <sub>2</sub>	84.6	111.7	105.9	95.8	96.0	96.0(11)	105.9(55)
LiH	32.5	57.2	49.5	42.4	43.0	42.8(27)	49.6(118)
HF	94.9	132.4	127.2	105.1	115.9	25.0(40)	126.0(274)
HC1	77.6	103.1	99.1	85.8	89.9	92.0(40)	98.0(274)
F <sub>2</sub>	-33.1	41.4	28.3	10.9	31.5	40.4(81)	33.9(1089)
Cl <sub>2</sub>	14.5	48.4	41.6	26.2	35.6	38.9(81)	42.1(1089)

		$D_{\rm e}$ (kcal/mol)				
bond	basis set	BOVB	VBCISD <sup>a</sup>	CCSD(T)	exptl	
F-F	6-31G*	36.2	32.3	32.8		
	cc-pVTZ	37.9	36.1	34.8	38.3	
Cl-Cl	6-31G*	40.0	41.6	40.5		
	cc-pVTZ	50.0	56.1	52.1	58.0	
Br-Br	6-31G*	41.3	44.1	41.2		
	cc-pVTZ	44.0	50.0	48.0	45.9	
F-Cl	6-31G*	47.9	49.3	50.2		
	cc-pVTZ	53.6	58.8	55.0	60.2	
H-H	6-31G**	105.4	105.4	105.9	109.6	
Li-Li	6-31G*	20.9	21.2	21.1	24.4	
$H_3C-H$	6-31G**	105.7	113.6	109.9	112.3	
$H_3C-CH_3$	6-31G*	94.7	90.0	95.6	96.7	
HO-OH	6-31G*	50.8	49.8	48.1	53.9	
$H_2N-NH_2$	6-31G*	68.5	70.5	66.5	$75.4 \pm 3$	
H <sub>3</sub> Si-H	6-31G**	93.6	90.2	91.8	97.6±3●	
H <sub>3</sub> Si-F	6-31G*	$140.4^{b}$	151.1	142.6	$160 \pm 7$	
H <sub>3</sub> Si-Cl	6-31G*	102.1	101.2	98.1	$113.7 \pm 4$	

<sup>*a*</sup> With Davidson correction. <sup>116 *b*</sup> Two-structure calculations ( $H_3Si^-F^+$  is omitted).

#### • The VBCI method :

If a complete basis of structures is included, if all VBCI excitations (S,D,T,Q,...) are included, if basis set extended to infinity...
⇒ converges to the «exact» (non-relativistic, 0K) solution of the Schrödinger equation

When pushed to their highest level, MO-based and
 VB-based methods ultimately converge to the same limit

- The VBCI method :
  - Very accurate (at the VBCISD+Davidson level)
  - Extremely simple to use
  - Full interpretability of the VB wave function is kept
  - but : it is very expensive...

Benchmark VB method

• New *ab initio* VB methods in development :

**- (ic)VBPT2 :** very cheap, the interpretability of the VB w.f. is lost at the moment (delocalized virtual orbitals)...

- **VBDFT** : cost «negligible» (VBSCF step), pb of double-counting of correlation effects, usual limitations of DFT

- **VB-QMC** : very expensive but massively parallel algorithms (thousands of processors)

- The VBSCF method
- The **BOVB** method
- The VBCI method





# What do you still have to know ?

#### Basis of structures :

The VB w.f. is a superposition of structures (covalent, ionic,...) :

$$\Psi_{VB} = \sum_{K} C_{K} \Phi_{K}$$

 $C_K$  : coefficients  $\Phi_K$  : VB structures

➡ Which structures to choose ?



 $\pi$  system of benzene (**6e**<sup>-</sup> in **6 orbitals**)

Three methods to generate a complete basis of structures :

- Rumer's Rules
- Weyl tableaux
- Young tableaux and operators

- Basis of structures / graphical Rumer's method :
  - Generation of a complete basis of covalent structures :
  - Put the orbitals around an imaginary circle
     Generate all possible couplings not displaying crossing bonds



 $\pi$  system of benzene (**6e**<sup>-</sup> in **6 orbitals**)

- Basis of structures / graphical Rumer's method :
  - Generation of a complete basis of covalent structures :
  - Put the orbitals around an imaginary circle
     Generate all possible couplings not displaying crossing bonds



- Complete and non-redundant set of VB structures
- VB structures are «chemically meaningful»

Redundant
 «Chemically
 meaninless»

- Basis of structures / graphical Rumer's method :
  - Generation of a complete basis of covalent structures :
  - Put the orbitals around an imaginary circle
     Generate all possible couplings not displaying crossing bonds



- Basis of structures / graphical Rumer's method :
  - Generation of a complete basis of ionic structures :
  - Choose a distribution of charges (configuration)
     Apply Rumer's rules on the rest of the system



3) Repeat 1) and 2) for all other configurations :


- Basis of structures / size limit :
  - Total number of covalent structures :

Weyl's formula:  $f_S^N = \frac{(2S+1)N!}{(\frac{1}{2}N+S+1)!(\frac{1}{2}N-S)!}.$ 

N	4	6	8	10	12	•••	
$f_s^N$	2	5	14	42	132	•••	

- Basis of structures / size limit :
  - Total number of covalent structures :

Weyl's formula:  $f_S^N = \frac{(2S+1)N!}{(\frac{1}{2}N+S+1)!(\frac{1}{2}N-S)!}$ .



• Total number of covalent + ionic structures :

Exponential wall !

Weyl's formula : 
$$g_{S}^{N,m} = \frac{2S+1}{m+1} \begin{pmatrix} m+1 \\ \frac{N}{2}+S+1 \end{pmatrix} \begin{pmatrix} m+1 \\ \frac{N}{2}-S \end{pmatrix}$$
  

$$\frac{N=m}{f_{S}^{N}} = \frac{4 \quad 6 \quad 8 \quad 14 \quad 28}{20 \quad 175 \quad 1764 \quad 2.76 \times 10^{6} \quad 2.07 \times 10^{14}}$$

#### Choice of an active space :

- Not all electrons are treated at the VB level :
  - an **active space** of electrons/orbitals treated at the **VB level**
  - the rest (called inactive or «spectators») at the MO level

$$\Psi_{VB} = \sum_{K} C_{K} \Phi_{K} \text{ with: } \Phi_{K} = \left| \{ \text{inactives} \} \{ \text{actives} \} \right|$$

#### Choice of an active space :

- Not all electrons are treated at the VB level :
  - an **active space** of electrons/orbitals treated at the **VB level**
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$$\Psi_{VB} = \sum_{K} C_{K} \Phi_{K} \text{ with: } \Phi_{K} = \left| \{ \text{inactives} \} \{ \text{actives} \} \right|$$

• Active space chosen depending on the chemical problem :

Example : SN2 Transition state : a **4-e**/**3-orbital** VB system



#### Definition of the orbitals :

- Active (VB) orbitals must always be strictly localized
- Inactive (MO) orbitals are preferably delocalized

#### ➡ Exercise 3 :

1) We consider a Heitler-London wave-function for H<sub>2</sub> made of two orbitals  $\varphi_a$  and  $\varphi_b$  partly delocalized onto the other center :  $\varphi_a = a + \varepsilon b$ ,  $\varphi_b = b + \varepsilon a$ , with a and b (strictly localized) atomic orbitals *a* and *b*.

2) Expand this wave-function in a basis of the atomic orbitals.

3) Is it legitimate to also incorporate ionic structures to this w.f.?

#### ➡ Exercise 3 (answers) :

$$\Psi = \left| \varphi_a \overline{\varphi_b} \right| + \left| \varphi_b \overline{\varphi_a} \right| = \left| (a + \varepsilon b)(\overline{b + \varepsilon a}) \right| + \left| (b + \varepsilon a)(\overline{a + \varepsilon b}) \right|$$
$$= \dots = (1 + \varepsilon^2)(|\underline{ab}| + |\underline{ba}|) + 2\varepsilon(|\underline{aa}| + |\underline{bb}|)$$
<sub>covalent</sub>

Active orbitals delocalized ⇒ ionic structures *implicitely* included (GVB / Spin Coupled methods)

- Active orbitals delocalized + explicit ionic structures  $\Rightarrow$  redundancies in the wave functions (instabilities,...)

➡ Active (VB) orbitals <u>must always</u> be strictly localized

# What do you get out of the calculation ?

## (good) numbers...



# ...but also insight !



Chemical insight / structure weights :

From computed {C<sub>K</sub>} coefficients :  $\Psi_{VB}(1,...,N) = \sum_{K} C_{K} \Phi_{K}(1,...,N)$ we can extract {W<sub>K</sub>} structure weights which verify :  $\sum_{K} W_{K} = 1$ 

→ quantification of the statistical importance of the ≠ structures usually expressed in % (100W<sub>K</sub>)

Chemical insight / structure weights :

From computed {C<sub>K</sub>} coefficients :  $\Psi_{VB}(1,...,N) = \sum_{K} C_{K} \Phi_{K}(1,...,N)$ we can extract {W<sub>K</sub>} structure weights which verify :  $\sum_{K} W_{K} = 1$ 

- → quantification of the statistical importance of the ≠ structures usually expressed in % (100W<sub>K</sub>)
- Most used : Chirgwin-Coulson weights :

$$\langle \Psi_{VB} | \Psi_{VB} \rangle = \sum_{K} \sum_{I} C_{K} C_{I} S_{KI} = 1 \Longrightarrow W_{K} = C_{K}^{2} + \sum_{I \neq K} C_{K} C_{I} S_{KI}$$
 with:  $S_{KI} = \langle \Phi_{K} | \Phi_{I} \rangle$ 

• Other options : Hiberty, Lowdin, Gallup (inverse) weights

#### • Chemical insight / structure weights :

<u>Ex</u> : reactivity of ozone vs. trisulfur :



#### Chemical insight / «diabatic states» :

It is possible to compute a VB w.f. which does not correspond to a real quantum state : a single structure of a subset of structures

computation of Resonance Energies (R.E.):



**1)** Optimize  $\Psi_{(1\leftrightarrow 2)}$ 

**2)** Optimize  $\Psi_1$  separately

3) *R.E.* = 
$$E(\Psi_1) - E(\Psi_{(1\leftrightarrow 2)})$$

#### Chemical insight / «diabatic states» :

It is possible to compute a VB w.f. which does not correspond to a real quantum state : a single structure of a subset of structures

computation of Resonance Energies (R.E.):



**1)** Optimize  $\Psi_{(1\leftrightarrow 2)}$ 

**2)** Optimize  $\Psi_1$  separately

3)  $R.E. = E(\Psi_1) - E(\Psi_{(1\leftrightarrow 2)})$ 

 $R.E. = 37 \ kcal/mol$ 

 $\Rightarrow$  formamide rotation barrier is due to  $\pi$  resonance energy

#### Chemical insight / «diabatic states» :

It is possible to compute a VB w.f. which does not correspond to a real quantum state : a single structure of a subset of structures

➡ Valence Bond diagrams (Shaik and Pross) for reactivity :



Part 4. lecture

#### Application 1,3-dipolar cycloadditions

## 1-3 dipolar cycloadditions

• Method for heterocyclic compounds synthesis :



• 1-3 dipole :

- 1 or 2  $\pi$  systems, 4  $\pi$  electrons on 3 centers ;

- Globally neutrals but **polarized** ;



## Some families of dipoles

#### Azomethine betaines :



#### **Reactions studied**

• Dipolarophiles : ethylene et acetylene :



• Marcus : correlation between **barriers** and **enthalpies** :



the **more exothermic** the reaction, the **lower** the **barrier** 

• Allylic dipole ; addition on ethylene :



• Allylic dipole ; addition on acetylene :



• Allylic dipole ; addition on acetylene :



• Propargylic dipole



• Propargylic dipole



## FMO theory ?



## Hammond postulate ?

• Hammond : TS is more reactant-like for more exothermic reactions



Different enthalpies, different **TS geometries** 

## Hammond principle ?

• All dipoles **distort** when going to TS :



## Hammond principle ?

• All dipoles **distort** when going to TS :



Ess & Houk\*: barriers only depend on dipole distortion energies ... but why ?...

\* Ess, D. H.; Houk, K. N. J. Am. Chem. Soc. 2008, 130, 10187

#### • VB description :



- Combination of 3 resonant structures

- Diradical structure (neutral) may have a significant weight, and is reactive...

• Importance of the diradical structure (BOVB calculations) :



#### - Diradical character is important !

B. Braida, C. Walter, B. Engels, P. C. Hiberty, J. Am. Chem. Soc. 2010, 132, 7631

dimanche 28 juillet 13

• Importance of the diradical structure (BOVB calculations) :



#### - Diradical character is important !

- Diradical character increases from reactants  $\rightarrow$  TS !

B. Braida, C. Walter, B. Engels, P. C. Hiberty, J. Am. Chem. Soc. 2010, 132, 7631

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• Importance of the diradical structure (BOVB calculations) :

H N	Ι	Reactants :	] <b>TS</b> :
H <sub>2</sub> C Z	$Z = O$ $Z = NH$ $Z = CH_2$	33.7 38.0 41.3	38.6 43.2 46.6
HC N Z	$Z = O$ $Z = NH$ $Z = CH_2$	21.3 26.5 26.3	32.1 35.7 35.4
• • • • • Z	$Z = O$ $Z = NH$ $Z = CH_2$	21.6 25.1 27.7	31.6 34.4 36.4

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H <sub>2</sub> C Z	Z = O	33.7	38.6
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HC N Z	Z = O Z = NH Z = CH2	21.3 26.5 26.3	32.1 35.7 35.4
• N=NZ	$Z = O$ $Z = NH$ $Z = CH_2$	21.6 25.1 27.7	31.6 34.4 36.4

What if... dipole distorsion would serve to **increase the diradical character** ?...

$$\begin{array}{c} a \ X = Y - \ddot{Z} \\ \oplus \ \odot \end{array} \\ b \ \ddot{X} - Y = Z \\ \odot \ \oplus \end{array} \\ c \ \dot{X} - \ddot{Y} - \dot{Z} \quad ground \ state \end{array}$$



1) Dipole distortion  $\rightarrow$  reach a «critical» diradical character

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1) Dipole distortion  $\rightarrow$  reach a «critical» diradical character

2) Dienophile attack (almost barrierless) !

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#### 1) If our postulated mechanism is correct :



If a **«critical» diradical character** has to be reached for the reaction to proceed : the **larger** the diradical weight in reactant, the **easier** the reaction

=> *Inverse correlation* between diradical *weights* and reaction *barriers* 

#### 1) Correlation diradical weights / barriers



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#### 1) Correlation diradical weights / barriers



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#### 1) Correlation diradical weights / barriers





#### 2) If our postulated mechanism is correct :



#### **2)** Correlation $G / \Delta H^{\ddagger}$ :



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• G : reactants :



#### • *G* : reactants vs. TS :



#### • *G* : reactants vs. TS :



If dipoles bend to reach some particular «critical» diradical character → similar G for all dipoles in TS geometry !

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#### Conclusions

- 1,3-dipoles are special reactants (violate ordinary laws)
- A mechanism is proposed, consistent with accurate ab initio data
- The **diradical character** is a key factor and the correlating quantity
- Reaction barriers can be predicted from reactants' properties



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B. Braida\*, C. Walter, B. Engels, P. C. Hiberty, J. Am. Chem. Soc. 2010, 132, 7631

#### Conclusions

VB theory allowed us to consider the 1,3 dipolar cycloaddition reaction from a different perspective :





# So, why everybody has not been doing VB ?

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#### • Historical reasons :



• Moderate number of structures :

• Given an active space, the selection of meaningful structures should be moderate (VBSCF : ~100s' ; BOVB : up to ~20)



• Large number of structure : computational issue, but also interpretative issue : VB theory may not be **relevant** in such cases

→ VB is not a «universal theory» of electronic structure

- Nonorthogonality :
  - Orthogonal orbitals cannot be strictly localized :



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To get strictly localized orbitals we need to use nonorthogonal orb

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- Nonorthogonality :
  - Working with non-orthogonal orbitals complicates a <u>LOT</u> all formulas :

$$\hat{H} = \sum_{i} \hat{h}_{i} + \sum_{i,j} \frac{1}{r_{ij}}$$

$$\overline{\overline{H}}\overline{C} = E\overline{\overline{S}}\overline{C}$$

<u>Ex</u> : a 4e/4o pb :



- Nonorthogonality :
  - Working with non-orthogonal orbitals complicates a LOT all formulas :

$$\hat{H} = \sum_{i} \hat{h}_{i} + \sum_{i,j} \frac{1}{r_{ij}} \qquad \overline{\overline{H}}\overline{C} = E\overline{\overline{S}}\overline{C} \qquad |a\overline{a}b\overline{b}| \\ |a\overline{a}c\overline{c}| \qquad |a\overline{a}b\overline{b}| \qquad |a\overline{a}c\overline{c}| \qquad 1 \quad 0 \quad 0 \quad 0 \quad 0 \\ 0 \quad 1 \quad 0 \quad 0 \quad 0 \quad 0 \\ 0 \quad 0 \quad 1 \quad 0 \quad 0 \quad 0 \\ 0 \quad 0 \quad 0 \quad 1 \quad 0 \quad 0 \\ 0 \quad 0 \quad 0 \quad 1 \quad 0 \quad 0 \\ 0 \quad 0 \quad 0 \quad 0 \quad 1 \quad 0 \\ 0 \quad 0 \quad 0 \quad 0 \quad 1 \quad 0 \\ 0 \quad 0 \quad 0 \quad 0 \quad 1 \quad 0 \\ 0 \quad 0 \quad 0 \quad 0 \quad 1 \quad 0 \\ 0 \quad 0 \quad 0 \quad 0 \quad 0 \quad 1 \\ |a\overline{a}b\overline{b}| \qquad |c\overline{c}d\overline{d}| \qquad \cdots \qquad |a\overline{a}b\overline{b}| \qquad |c\overline{c}d\overline{d}|$$

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$$\hat{H} = \sum_{i} \hat{h}_{i} + \sum_{i,j} \frac{1}{r_{ij}} \qquad \overline{H}\overline{C} = E\overline{S}\overline{C} \qquad |a\overline{a}b\overline{b}| \\ \mathbf{VB theory :} \\ (nonorthogonal orbs.) \qquad \vdots \\ (nonorthogonal orbs.) \qquad \vdots \\ \underline{Ex :} a 4e/4o pb : \qquad a\overline{b} \\ |a\overline{a}b\overline{b}| \\ |c\overline{c}d\overline{d}| \\ |c\overline{c}d\overline{d}| \\ |c\overline{c}d\overline{d}| \\ (D|D') = S_{ac}^{2}S_{bd}^{2} - 2S_{ad}S_{ac}S_{bd}S_{bc} + S_{ad}^{2}S_{bc}^{2}$$

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$$\hat{H} = \sum_{i} \hat{h}_{i} + \sum_{i,j} \frac{1}{r_{ij}}$$

$$\overline{H}\overline{C} = E\overline{S}\overline{C}$$

$$|a\overline{a}b\overline{b}|$$

$$|a\overline{a}c\overline{c}|$$

$$(nonorthogonal orbs.)$$

- Nonorthogonality :
  - Working with non-orthogonal orbitals complicates a <u>LOT</u> all formulas :
  - <u>Ex 2 :</u> (ic)VBPT2 formulas :

$$\begin{split} & \left\langle \Psi_{x_{1}y_{1}}^{i} \left| \hat{H}_{0} \right| \Psi_{x_{2}y_{2}}^{i} \right\rangle = \left( E_{imacl}^{(0)} - \varepsilon_{i} - \varepsilon_{j} \right) \left\langle \Psi_{x_{1}y_{1}}^{i} \left| \Psi_{x_{2}y_{2}}^{i} \right\rangle \\ & + \left\{ \left( 4 s_{x_{2}x_{1}} s_{y_{2}y_{1}} - 2 s_{y_{2}x_{1}} s_{x_{2}y_{1}} \right) E_{acl}^{(0)} \right. \\ & + 4 f_{x_{2}x_{1}} s_{y_{2}y_{1}} + 4 f_{y_{2}y_{1}} s_{x_{2}x_{1}} - 2 f_{x_{2}y_{1}} s_{y_{2}y_{1}} - 2 f_{y_{2}x_{1}} s_{x_{2}y_{1}} \\ & + \left[ \left( f_{x_{1}} s_{x_{2}y_{1}} s_{y_{2}u} + f_{y_{1}} s_{y_{2}x_{1}} s_{x_{2}u} - 2 f_{x_{1}} s_{y_{2}y_{1}} s_{x_{2}u} - 2 f_{y_{2}y_{1}} s_{x_{2}x_{1}} s_{y_{1}u} \right) \\ & + \left[ f_{x_{2}} s_{y_{2}x_{1}} s_{y_{1}u} + f_{y_{2}} s_{x_{2}y_{1}} s_{x_{1}u} - 2 f_{x_{2}} s_{y_{2}y_{1}} s_{x_{1}u} - 2 f_{y_{2}y_{1}} s_{x_{2}x_{1}} s_{y_{1}u} \right) \\ & + \left( f_{y_{2}x_{1}} s_{y_{1}u} s_{x_{2}t} + f_{x_{2}y_{1}} s_{x_{1}u} s_{y_{2}t} - 2 f_{x_{2}x_{1}} s_{y_{1}u} s_{y_{2}t} - 2 f_{y_{2}y_{1}} s_{x_{1}u} s_{y_{2}t} \right] D^{uu} \\ & + \left[ f_{u} \left( s_{y_{2}x_{1}} s_{x_{2}v} s_{y_{1}w} + s_{x_{2}y_{1}} s_{y_{2}v} s_{x_{1}w} - 2 s_{y_{2}y_{1}} s_{x_{2}v} s_{x_{1}w} - 2 s_{x_{2}x_{1}} s_{y_{2}v} s_{x_{1}w} - 2 s_{x_{2}x_{1}} s_{y_{2}v} s_{y_{1}w} \right] \right| \Pi^{uv,uw} \\ & + f_{u} s_{y_{1}v_{1}} s_{x_{2}v} s_{y_{1}w} + s_{x_{2}y_{1}} s_{y_{2}v} s_{x_{1}w} - 2 s_{y_{2}y_{1}} s_{x_{2}v} s_{x_{1}w} - 2 s_{x_{2}x_{1}} s_{y_{2}v} s_{y_{1}w} + s_{y_{2}v} s_{y_{1}w} s_{x_{1}w} \right] \right| \Pi^{uv,uw} \\ & + f_{u} s_{y_{1}v_{1}} s_{y_{2}x_{1}} - 2 s_{x_{2}x_{1}} s_{y_{2}v} s_{x_{1}w} - 2 s_{x_{2}x_{1}} s_{y_{2}y_{1}} \\ & + \delta^{ii} \left\{ \left( 4 s_{x_{2}y_{1}} s_{y_{2}x_{1}} - 2 s_{x_{2}x_{1}} s_{y_{1}y_{1}} - 2 f_{x_{2}x_{1}} s_{y_{2}y_{1}} \right) \right. \\ & + \left( f_{y_{1}} s_{x_{2}x_{1}} s_{y_{2}u} + f_{y_{2}} s_{x_{2}x_{1}} s_{y_{1}u} - 2 f_{y_{2}} s_{x_{2}y_{1}} s_{x_{1}u} - 2 f_{x_{2}} s_{y_{2}x_{1}} s_{y_{1}u} \right) \\ & + \left( f_{y_{1}} s_{x_{2}x_{1}} s_{y_{1}u} + f_{y_{2}} s_{x_{2}x_{1}} s_{y_{1}u} - 2 f_{y_{2}} s_{y_{2}x_{1}} s_{y_{1}u} \right) \\ & + \left( f_{y_{1}} s_{x_{2}x_{1}} s_{y_{2}u} + f_{x_{1}} s_{y_{1}y_{1}} s_{y_{2}u} - 2 f_{y_{2}} s_{y_{1}} s_{y_{1}u} - 2 f_{y_{2}} s_{y_{2}x_{1}$$

- Nonorthogonality :
  - Working with non-orthogonal orbitals complicates a **LOT**
  - Slowed dramatically the extension of VB community

#### But :

- Now extremely efficient algorithms, based on advanced algebra:\*
  - 1) Second Quantization Technique for non-orthogonal orbitals ;
  - 2) Tensor Analysis ;
  - 3) Automatic Formula/Code Generator

\* Chen Z., Chen X, and Wu W. J. Chem. Phys. 138, 164119 (2013) \* Chen Z., Chen X, and Wu W. J. Chem. Phys. 138, 164120 (2013)

- Moderate overlap between active orbitals :
  - Large overlap between active orbitals : the meaning of the VB wave function (covalent vs. ionic) may be lost :



 $S \approx 0.5 \text{ or less} : ok$ 

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H_2/LiH: S \approx 0.8!
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basically the only two really pathologic cases

• Moderate overlap between active orbitals :



• If very large basis sets are used, any one-center AO can take the shape of a fully delocalized MO  $\Rightarrow$  R.E. abnormally small, technical issue (BOVB),...

*«The more accurate the calculations become, the more the concepts tend to vanish into thin air.» (R. S. Mulliken)* 

• Moderate overlap between active orbitals :



#### The good news is :

• Up to triple-zeta basis set without diffuse functions\* usually allows to get close enough to chemical accuracy while still avoiding this problem

\* *Except if the system is anionic* 

- Moderate number of structures
- Nonorthogonality (not any more)
- Moderate basis set size
- **Kow-how and expertise**, whereas VB is hardly taught any more



An Ab Initio Non-orthogonal Valence Bond Program

## It's up to you (tutorials) !

• Tutoriel 1 (Thurs. 15h-18h) : Basics of XMVB computations

https://wiki.lct.jussieu.fr/workshop/index.php/VB\_tutorial