Algorithms for the Optimization of Non-orthogonal Orbitals Using Biothogonal Expansions A Perspective Talk

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Orbital Optimization

March 27, 2017 1 / 17

The purpose is to synthetically and didactically present the state-of-the-art of a specific approach

- pinpointing questions
- opening current challenges to solve
- and providing food for thought for discussion.
- (Lecturers time: 20-25 minutes)

Interpretation and chemical insight

- Realizes that non-orthogonal problems leads to inherent computational problems
- Instead: Gives a better way of understanding bonding
- (From my point of view: far too defensive)

An efficient computational tool

- Methods based on non-othogonal orbitals will not replace standard methods in general
- But, there is a significant group of molecules containing atoms with partial filled shells and limited overlap, that are best described using non-orthogonal methods

The major remaining challenge in quantum chemistry: Static correlation

Status

- XASSCF, X=C,R,G leads to expansions that will not be able to treat say 30 electrons in 30 orbitals for the next few decades
- A number of newer methods are emerging: DMRG, Tensor-methods, Monte Carlo Cl ... - impressive progress

But

- I would like to do the expansions directly in terms of quantities that are chemically meaningfull.
- Thus, whereas most are moving forward (with impressive speed), I will take the opposite approach: Look back (a prerogative of the older generation) and re-examine the valence bond method

When do we perform standard XASSCF(X=C,R,G) calculations?

Molecules with partly occupied shells with limited overlap

- Includes molecules of main-group elements at stretched geometries and transition, lanthanide, and actinice compounds at equilibrium.
- Small overlap \rightarrow limited splitting between bonding and anti-bonding orbitals.
- A RHF calculation produces a wave function with many ionic terms not present in the full wave function.
- The major correction of the MCSCF calculation is to introduce left-right correlation to clean up the mess of RHF.
- (MCSCF is also needed when there are atomic configurations with nearly identical energies)

The standard procedure

- Use atomic rather than delocalized orbitals
- The number of electrons on each atom/ set of atomic orbitals is specified several combinations may be needed
- When the partly filled atomic orbitals on different atoms have a small overlap, the covalent configuration is strongly dominating
- Molecules, where the standard MO-RHF-MCSCF approach is most requiring, have (often) simple non-orthogonal expansions.
- lonic terms are introduced by orbital-optimization allowing limited delocalization of orbitals.

Describes electronic wave functions as products of creation operators a_{P}^{\dagger}

- Example, the ground state of H₂: $a^{\dagger}_{1\sigma\alpha}a^{\dagger}_{1\sigma\beta}|\text{Vac}\rangle$
- The conjugate operators, a_P removes (if possible an electron in spin-orbital P.

• The simplicity of the standard orthogonal theories arise from the simple anticommutator $[a_P^\dagger,a_Q]_+=\delta_{QP}$

Non-orthogonal orbitals

- For non-orthogonal orbitals, all the problems arises from anti-commutation relation $[a_P^{\dagger},a_Q]_+=S_{QP}$
- One may introduce another set of orbitals, the bi-orthonormal orbitals, $\tilde{a}_Q=\sum_{Q'}S_{Q'Q}^{-1}a_{Q'}$

• Their anticommutator with the original basis is $[a_P^{\dagger}, \tilde{a}_Q]_+ = \delta_{QP}$

Two types two be discussed

- A single set of non-orthogonal orbitals
- Several sets of orbitals
 - Each set is orthogonal and uses all SDs of the N-electron space, say FCI
 - 2 Each set is non-orthogonal and uses part of the N-electron space

Algorithms for optimizing a single set of orbitals

Target form of wave function

- A set of orthonormal double occupied inactive orbitals
- A set of variably occupied non-orthogonal orbitals
- A set of orthonormal unoccupied orbitals
- The various spaces are kept orthonormal to each other

Design goals

- Complete and non-redundant parameterization
- Must at all points be faster and simpler that the corresponding CASSCF expansion (requires that we bypass the three- and four-particle density matrices of standard non-orthogonal orbital optimization).
- Should be easy to extent to response theory for calculating non-linar and excited state propeties including magnetic properties

Algorithms for optimizing a single set of orbitals

Parameterization of orbital rotations

Use exponential operators

•
$$|\tilde{0}\rangle = \exp(-\hat{\kappa}_A) \exp(-\hat{\kappa}_S)|0\rangle / \sqrt{\langle 0| \exp(-2\hat{\kappa}_S)|0\rangle}$$

• $\kappa_A = \sum_{p>q} \kappa^A_{pq} (E_{pq} - E_{qp})$
• $\kappa_S = \sum_{x>y} \kappa^S_{xy} (E_{xy} + E_{yx})$ (sum only over active orbitals)
• $E_{pq} = \sum_{\sigma=\alpha,\beta} a^{\dagger}_{p\sigma} a_{q\sigma}$

Note

- Operators written in terms of the original operators not bi-orthogonality here
- The operator exp(-κ̂_A) conserves the metric. The operator exp(-κ̂_S) changes the metric, but only between the active orbitals.
- Internal rotations in the inactive and secondary orbitals are redundant and therefore not in use.

The energy

• $E(\kappa) = \frac{\langle \tilde{0} | \hat{H} | \tilde{0} \rangle}{\langle \tilde{0} | \tilde{0} \rangle}$

•
$$| ilde{0}
angle = \exp(-\hat{\kappa}_{\mathcal{A}})\exp(-\hat{\kappa}_{\mathcal{S}})|0
angle/\sqrt{\langle 0|\exp(-2\hat{\kappa}_{\mathcal{S}})|0
angle}$$

•
$$E = \frac{\langle 0|\exp(-\hat{\kappa}_S)\exp(\hat{\kappa}_A)\hat{H}\exp(-\hat{\kappa}_A)\exp(-\hat{\kappa}_S)|0\rangle}{\langle 0|\exp(-2\hat{\kappa}_S)|0\rangle}$$

- The κ_A terms may be expanded as usual using the BCH expansion
- No BCH expansion for the κ_S terms
- Bioorthogonal reexpressions are used, either the right or left hand side is expressed in the bioorthogonl expansion

Algorithms for optimizing a single set of orbitals The gradient

(The energy:
$$E = \frac{\langle 0| \exp(-\hat{\kappa}_S) \exp(\hat{\kappa}_A) \hat{H} \exp(-\hat{\kappa}_A) \exp(-\hat{\kappa}_S) | 0 \rangle}{\langle 0| \exp(-2\hat{\kappa}_S) | 0 \rangle}$$
)

The gradient (assuming $\langle 0|0 angle = 1$

2
$$E_{S,xy}^{[1]} = \frac{\partial E}{\partial \kappa_{xy}^{S}} = -2\langle 0|(E_{xy} + E_{yx})(H - E)|0\rangle$$

The gradient wrt symmetric operators

- Is the term that in standard algorithms requires the 3-electron density
- Note that the excitation is in the active space
- Proceed therefore instead as
 - Calculate sigma-vector, $| ilde{S}
 angle = (H-E)|0
 angle$ in biorthogonal basis
 - Calculate one-electron transition density $\langle 0|(E_{xy}+E_{yx})| ilde{S}
 angle$

A novel approach to orbital optimization of non-orthogonal expansions

Example of convergence: Cr₂ in ano-4 4321, 3 Ångstöm

Iteration	$E - E_{conv}$	Step-size
1	0.0431467207	$0.339 imes10^{00}$
2	0.0018950153	$0.162 imes10^{00}$
3	0.0000270094	$0.508 imes10^{-1}$
4	0.0000001751	$0.414 imes10^{-2}$
5	0.000000003	$0.206 imes 10^{-3}$

Comments

- Distance is where 3d-3d bonding occurs
- The accuracy of initial guess (atomic orbitals) is noteworthy
- Optimization takes about 50 seconds on my (now retired) Macbook for one geometry

Convergence towards CASSCF of the hierarchy for the chromium trimer



- Figure includes the results using the initial and optimised orbitals
- Energies are relative to those at 3.0 Å.
- Orbital optimisation is perhaps not essential?
- Very accurate results are obtained by including up to double excitations
- Dimensions: N(18,18,0) = 4 862, N(18,18,2) = 1 986 842, N(18,18,9) = 112 318 492.

- up to 18-20 unpaired electrons can be used. Selected of the crucial parts of the code has been extended to treat 30 unpaired electrons.
- Standard CI is still going in several parts of the code, but these will be removed in the next year.
- Full second order (including orbital-Cl coupling is also under development, as are response theory in the frequency domain.
- Non-orthogonal expansions containing hundred of millions of SDs have been realized (yes, in the Mac-book)

Algorithms for optimizing wave functions containing several sets of orbitals

Why several sets of orbitals, consider F_2

- The correlation energy increases (in absolute magnitude) when the two F-atoms approach each other
- Explanations:
 - The standard answer: When the two F-atoms approaches each other, the electron density increases, and therefore also the correlation energy
 - The VB person: Yes, but there is an effect more: When the two atoms appoaches each other, there is also a deviation from RHF by the tendency for an electron to jump from on F-atom to the other, forming the F+ F- ionic term.
- The 2p-orbitals for F⁻ are pretty different from the orbitals for the neutral F and therefore to describe the above effect in an efficient manner, one needs several sets of orbitals.
- The simultaneous optimization of several sets of orbitals is also needed to desribe for example mixed valence states.

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Orbital Optimization

Algorithms for optimizing wave functions containing several sets of orbitals

Two sets of codes

- One where, for example numerous HF or CAS states, each with different orthonormal orbitals are added and optimized
- One where, say a CAspace is divided into several parts, and each part is given a set of non-orthonormal orbitals
- Not simple coding

Central aspects

- The gradient can again be described only in terms of one- and two-electron densities using a number of tricks and rewrites
- The explicit Hessian requires now a three-body term
- A direct approach, where the Hessian times a vector is calculated, may again be formulated in terms of two-electron terms.