# Algorithms for the Optimization of Non-orthogonal Orbitals Using Biothogonal Expansions A Perspective Talk 

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## Perspective talks

The purpose is to synthetically and didactically present the state-of-the-art of a specific approach
(1) pinpointing questions
(2) opening current challenges to solve
(3) and providing food for thought for discussion.
(9) (Lecturers time: 20-25 minutes)

## What is the role of Nonorthogonal methods?

## Interpretation and chemical insight

- Realizes that non-orthogonal problems leads to inherent computational problems
- Instead: Gives a better way of understanding bonding
- (From my point of view: far too defensive)


## An efficient computational tool

- Methods based on non-othogonal orbitals will not replace standard methods in general
- But, there is a signficant group of molecules containing atoms with partial filled shells and limited overlap, that are best described using non-orthogonal methods


## The major remaining challenge in quantum chemistry:

 Static correlation
## Status

- XASSCF, $X=C, R, G$ leads to expansions that will not be able to treat say 30 electrons in 30 orbitals for the next few decades
- A number of newer methods are emerging: DMRG, Tensor-methods, Monte Carlo $\mathrm{Cl} . .$. - impressive progress


## But

- I would like to do the expansions directly in terms of quantities that are chemically meaningfull.
- Thus, whereas most are moving forward (with impressive speed), I will take the opposite approach: Look back (a prerogative of the older generation) and re-examine the valence bond method


## When do we perform standard $\operatorname{XASSCF}(X=C, R, G)$ calculations?

Molecules with partly occupied shells with limited overlap

- Includes molecules of main-group elements at stretched geometries and transition, lanthanide, and actinice compounds at equilibrium.
- Small overlap $\rightarrow$ limited splitting between bonding and anti-bonding orbitals.
- A RHF calculation produces a wave function with many ionic terms not present in the full wave function.
- The major correction of the MCSCF calculation is to introduce left-right correlation to clean up the mess of RHF.
- (MCSCF is also needed when there are atomic configurations with nearly identical energies)


## The non-orthogonal/valence bond approach

## The standard procedure

- Use atomic rather than delocalized orbitals
- The number of electrons on each atom/set of atomic orbitals is specified - several combinations may be needed
- When the partly filled atomic orbitals on different atoms have a small overlap, the covalent configuration is strongly dominating
- Molecules, where the standard MO-RHF-MCSCF approach is most requiring, have (often) simple non-orthogonal expansions.
- lonic terms are introduced by orbital-optimization allowing limited delocalization of orbitals.


## Second Quantization

Describes electronic wave functions as products of creation operators $a_{P}^{\dagger}$

- Example, the ground state of $\mathrm{H}_{2}: a_{1 \sigma \alpha}^{\dagger} a_{1 \sigma \beta}^{\dagger}|\mathrm{Vac}\rangle$
- The conjugate operators, $a_{P}$ removes (if possible an electron in spin-orbital $P$.
- The simplicity of the standard orthogonal theories arise from the simple anticommutator $\left[a_{P}^{\dagger}, a_{Q}\right]_{+}=\delta_{Q P}$


## Non-orthogonal orbitals

- For non-orthogonal orbitals, all the problems arises from anti-commutation relation $\left[a_{P}^{\dagger}, a_{Q}\right]_{+}=S_{Q P}$
- One may introduce another set of orbitals, the bi-orthonormal orbitals, $\tilde{a}_{Q}=\sum_{Q^{\prime}} S_{Q^{\prime} Q^{\prime}}^{-1} Q_{Q^{\prime}}$
- Their anticommutator with the original basis is $\left[a_{P}^{\dagger}, \tilde{a}_{Q}\right]_{+}=\delta_{Q P}$


## Algorithms for orbital optimization

Two types two be discussed

- A single set of non-orthogonal orbitals
- Several sets of orbitals
(1) Each set is orthogonal and uses all SDs of the N -electron space, say FCI
(2) Each set is non-orthogonal and uses part of the N -electron space


## Algorithms for optimizing a single set of orbitals

## Target form of wave function

- A set of orthonormal double occupied inactive orbitals
- A set of variably occupied non-orthogonal orbitals
- A set of orthonormal unoccupied orbitals
- The various spaces are kept orthonormal to each other


## Design goals

- Complete and non-redundant parameterization
- Must at all points be faster and simpler that the corresponding CASSCF expansion - (requires that we bypass the three- and four-particle density matrices of standard non-orthogonal orbital optimization).
- Should be easy to extent to response theory for calculating non-linar and excited state propeties including magnetic properties ....


## Algorithms for optimizing a single set of orbitals

Parameterization of orbital rotations

## Use exponential operators

- $|\tilde{0}\rangle=\exp \left(-\hat{\kappa}_{A}\right) \exp \left(-\hat{\kappa}_{S}\right)|0\rangle / \sqrt{\langle 0| \exp \left(-2 \hat{\kappa}_{S}\right)|0\rangle}$
(1) $\kappa_{A}=\sum_{p>q} \kappa_{p q}^{A}\left(E_{p q}-E_{q p}\right)$
(2) $\kappa_{S}=\sum_{x>y} \kappa_{x y}^{S}\left(E_{x y}+E_{y x}\right)$ (sum only over active orbitals)
- $E_{p q}=\sum_{\sigma=\alpha, \beta} a_{p \sigma}^{\dagger} a_{q \sigma}$


## Note

- Operators written in terms of the original operators - not bi-orthogonality here
- The operator $\exp \left(-\hat{\kappa}_{A}\right)$ conserves the metric. The operator $\exp \left(-\hat{\kappa}_{S}\right)$ changes the metric, but only between the active orbitals.
- Internal rotations in the inactive and secondary orbitals are redundant and therefore not in use.


## Algorithms for optimizing a single set of orbitals

## Energy

## The energy

- $E(\kappa)=\frac{\langle\tilde{0}| \hat{H}|\overline{0}\rangle}{\langle\tilde{0} \mid \hat{0}\rangle}$
- $|\tilde{0}\rangle=\exp \left(-\hat{\kappa}_{A}\right) \exp \left(-\hat{\kappa}_{S}\right)|0\rangle / \sqrt{\langle 0| \exp \left(-2 \hat{\kappa}_{S}\right)|0\rangle}$
- $E=\frac{\langle 0| \exp \left(-\hat{\kappa}_{S}\right) \exp \left(\hat{\kappa}_{A}\right) \hat{H} \exp \left(-\hat{\kappa}_{A}\right) \exp \left(-\hat{\kappa}_{S}\right)|0\rangle}{\langle 0| \exp \left(-2 \hat{\kappa}_{S}\right)|0\rangle}$
- The $\kappa_{A}$ terms may be expanded as usual using the BCH expansion
- No BCH expansion for the $\kappa_{S}$ terms
- Bioorthogonal reexpressions are used, either the right or left hand side is expressed in the bioorthogonl expansion


## Algorithms for optimizing a single set of orbitals

## The gradient

(The energy: $\left.E=\frac{\langle 0| \exp \left(-\hat{\kappa}_{S}\right) \exp \left(\hat{\kappa}_{A}\right) \hat{H} \exp \left(-\hat{\kappa}_{A}\right) \exp \left(-\hat{\kappa}_{S}\right)|0\rangle}{\langle 0| \exp \left(-2 \hat{\kappa}_{S}\right)|0\rangle}\right)$
The gradient (assuming $\langle 0 \mid 0\rangle=1$
(1) $E_{A, p q}^{[1]}=\frac{\partial E}{\partial \kappa_{p q}^{A}}=\langle 0|\left[E_{p q}-E_{q p}, H\right]|0\rangle$
(2) $E_{S, x y}^{[1]}=\frac{\partial E}{\partial \kappa_{x y}^{S}}=-2\langle 0|\left(E_{x y}+E_{y x}\right)(H-E)|0\rangle$

The gradient wrt symmetric operators

- Is the term that in standard algorithms requires the 3-electron density
- Note that the excitation is in the active space
- Proceed therefore instead as
- Calculate sigma-vector, $|\tilde{S}\rangle=(H-E)|0\rangle$ in biorthogonal basis
- Calculate one-electron transition density $\langle 0|\left(E_{x y}+E_{y x}\right)|\tilde{S}\rangle$


## A novel approach to orbital optimization of non-orthogonal expansions

## Example of convergence: $\mathrm{Cr}_{2}$ in ano-4 4321, 3 Ångstöm

| Iteration | $E-E_{\text {conv }}$ | Step-size |
| :---: | :---: | :---: |
| 1 | 0.0431467207 | $0.339 \times 10^{00}$ |
| 2 | 0.0018950153 | $0.162 \times 10^{00}$ |
| 3 | 0.0000270094 | $0.508 \times 10^{-1}$ |
| 4 | 0.0000001751 | $0.414 \times 10^{-2}$ |
| 5 | 0.0000000003 | $0.206 \times 10^{-3}$ |

## Comments

- Distance is where 3d-3d bonding occurs
- The accuracy of initial guess (atomic orbitals) is noteworthy
- Optimization takes about 50 seconds on my (now retired) Macbook for one geometry


## Convergence towards CASSCF of the hierarchy for the chromium trimer



- Figure includes the results using the initial and optimised orbitals
- Energies are relative to those at 3.0 Å.
- Orbital optimisation is perhaps not essential?
- Very accurate results are obtained by including up to double excitations
- Dimensions: $\mathrm{N}(18,18,0)=4$ 862, $\mathrm{N}(18,18,2)=1986$ 842, $N(18,18,9)=112318492$.


## Current status and the future plans/dreams

- up to 18-20 unpaired electrons can be used. Selected of the crucial parts of the code has been extended to treat 30 unpaired electrons.
- Standard Cl is still going in several parts of the code, but these will be removed in the next year.
- Full second order (including orbital-Cl coupling is also under development, as are response theory in the frequency domain.
- Non-orthogonal expansions containing hundred of millions of SDs have been realized (yes, in the Mac-book)


## Algorithms for optimizing wave functions containing several sets of orbitals

## Why several sets of orbitals, consider $\mathrm{F}_{2}$

- The correlation energy increases (in absolute magnitude) when the two F-atoms approach each other
- Explanations:
(1) The standard answer: When the two F-atoms approaches each other, the electron density increases, and therefore also the correlation energy
(2) The VB person: Yes, but there is an effect more: When the two atoms appoaches each other, there is also a deviation from RHF by the tendency for an electron to jump from on F-atom to the other, forming the $\mathrm{F}+\mathrm{F}$ - ionic term.
- The 2 p -orbitals for $\mathrm{F}^{-}$are pretty different from the orbitals for the neutral $F$ and therefore to describe the above effect in an efficient manner, one needs several sets of orbitals.
- The simultaneous optimization of several sets of orbitals is also needed to desribe for example mixed valence states.


## Algorithms for optimizing wave functions containing several sets of orbitals

## Two sets of codes

- One where, for example numerous HF or CAS states, each with different orthonormal orbitals are added and optimized
- One where, say a CAspace is divided into several parts, and each part is given a set of non-orthonormal orbitals
- Not simple coding


## Central aspects

- The gradient can again be described only in terms of one- and two-electron densities using a number of tricks and rewrites
- The explicit Hessian requires now a three-body term
- A direct approach, where the Hessian times a vector is calculated, may again be formulated in terms of two-electron terms.

