



Extended RPA with 3-body ground-state correlations

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- Formulation
- Application to Lipkin model
- Summary

Formulation

TDDM equations

$$\rho(11', t) = \sum_{\alpha\alpha'} n_{\alpha\alpha'}(t) \psi_{\alpha}(1, t) \psi_{\alpha'}^*(1', t)$$

$$C_2'(121'2', t) = \rho_2 - A(\rho\rho)$$

$$= \sum_{\alpha\beta\alpha'\beta'} C_{\alpha\beta\alpha'\beta'}(t) \psi_{\alpha}(1, t) \psi_{\beta}(1, t) \psi_{\alpha'}^*(1', t) \psi_{\beta'}^*(1', t)$$

$$i\hbar \frac{\partial}{\partial t} \psi_{\alpha} = h(\rho) \psi_{\alpha}$$

$$i\hbar \dot{n}_{\alpha\alpha'} = \sum_{\lambda_1\lambda_2\lambda_3} \left(\langle \alpha\lambda_3 | v | \lambda_1\lambda_2 \rangle C_{\lambda_1\lambda_2\alpha'\lambda_3} - C_{\alpha\lambda_3\lambda_1\lambda_2} \langle \lambda_1\lambda_2 | v | \alpha'\lambda_3 \rangle \right)$$

$$i\hbar \dot{C}_{\alpha\beta\alpha'\beta'} = B_{\alpha\beta\alpha'\beta'} + P_{\alpha\beta\alpha'\beta'} + H_{\alpha\beta\alpha'\beta'}$$

Extended RPA(ERPA)

Excitation operator

$$Q_{\mu}^{+} = \sum (x_{\lambda\lambda'}^{\mu} : a_{\lambda}^{+} a_{\lambda'} : + X_{\lambda_1\lambda_2\lambda_1'\lambda_2'}^{\mu} : a_{\lambda_1}^{+} a_{\lambda_2}^{+} a_{\lambda_2'} a_{\lambda_1'} :)$$

$$Q_{\mu}^{+} |\Psi_0\rangle = |\Psi_{\mu}\rangle$$

$$Q_{\mu} |\Psi_0\rangle = 0$$

Equation-of-motion approach

$$\langle \Psi_0 | [[: a_{\alpha}^{+} a_{\alpha} :, H], Q_{\mu}^{+}] | \Psi_0 \rangle = \omega_{\mu} \langle \Psi_0 | [[: a_{\alpha}^{+} a_{\alpha} :, Q_{\mu}^{+}] | \Psi_0 \rangle$$

$$\langle \Psi_0 | [[: a_{\alpha}^{+} a_{\beta}^{+} a_{\beta} a_{\alpha} :, H], Q_{\mu}^{+}] | \Psi_0 \rangle = \omega_{\mu} \langle \Psi_0 | [[: a_{\alpha}^{+} a_{\beta}^{+} a_{\beta} a_{\alpha} :, Q_{\mu}^{+}] | \Psi_0 \rangle$$

ERPA equation

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} x^\mu \\ X^\mu \end{pmatrix} = \omega_\mu \begin{pmatrix} S_{11} & T_{12} \\ T_{21} & S_{22} \end{pmatrix} \begin{pmatrix} x^\mu \\ X^\mu \end{pmatrix}$$

$$A(\alpha\alpha' : \lambda\lambda') = \langle 0 | [[: a_\alpha^+ a_\alpha : , H], : a_\lambda^+ a_\lambda :] | 0 \rangle$$

$$B(\alpha\alpha' : \lambda_1\lambda_2\lambda_1'\lambda_2') = \langle 0 | [[: a_\alpha^+ a_\alpha : , H], : a_{\lambda_1}^+ a_{\lambda_2}^+ a_{\lambda_2'} a_{\lambda_1'} :] | 0 \rangle$$

$$C(\alpha\beta\alpha' \beta' : \lambda\lambda') = \langle 0 | [[: a_\alpha^+ a_\beta^+ a_\beta a_\alpha : , H], : a_\lambda^+ a_\lambda :] | 0 \rangle$$

$$D(\alpha\beta\alpha' \beta' : \lambda_1\lambda_2\lambda_1'\lambda_2') = \langle 0 | [[: a_\alpha^+ a_\beta^+ a_\beta a_\alpha : , H], : a_{\lambda_1}^+ a_{\lambda_2}^+ a_{\lambda_2'} a_{\lambda_1'} :] | 0 \rangle$$

$$S_{11}(\alpha\alpha' : \lambda\lambda') = \langle 0 | [[: a_\alpha^+ a_\alpha : , : a_\lambda^+ a_\lambda :] | 0 \rangle$$

$$T_{12}(\alpha\alpha' : \lambda_1\lambda_2\lambda_1'\lambda_2') = \langle 0 | [[: a_\alpha^+ a_\alpha : , : a_{\lambda_1}^+ a_{\lambda_2}^+ a_{\lambda_2'} a_{\lambda_1'} :] | 0 \rangle$$

$$T_{21}(\alpha\beta\alpha' \beta' : \lambda\lambda') = \langle 0 | [[: a_\alpha^+ a_\beta^+ a_\beta a_\alpha : , : a_\lambda^+ a_\lambda :] | 0 \rangle$$

$$S_{22}(\alpha\beta\alpha' \beta' : \lambda_1\lambda_2\lambda_1'\lambda_2') = \langle 0 | [[: a_\alpha^+ a_\beta^+ a_\beta a_\alpha : , : a_{\lambda_1}^+ a_{\lambda_2}^+ a_{\lambda_2'} a_{\lambda_1'} :] | 0 \rangle$$

Jacobi 's identity must be fulfilled
for ERPA to be Hermitian

$$\langle 0 | [[B, H], A] | 0 \rangle - \langle 0 | [[A, H], B] | 0 \rangle = \langle 0 | [H, [A, B]] | 0 \rangle = 0$$

$$A, B = a_{\alpha}^{+} a_{\alpha}, \text{ or } a_{\alpha}^{+} a_{\beta}^{+} a_{\beta} a_{\alpha}$$

This implies

$$\langle 0 | [a_{\alpha}^{+}, a_{\alpha}, H] | 0 \rangle = 0$$

$$\langle 0 | [a_{\alpha}^{+}, a_{\beta}^{+}, a_{\beta} a_{\alpha}, H] | 0 \rangle = 0$$

$$\langle 0 | [a_{\alpha}^{+}, a_{\beta}^{+}, a_{\gamma}^{+}, a_{\gamma} a_{\beta} a_{\alpha}, H] | 0 \rangle = 0$$



One-body part of ERPA

$$Ax^\mu = \omega_\mu S_{11} x^\mu$$

is the same as self-consistent RPA (SCRPA) equation

Ground state $|0\rangle$ is given by

$n_{\alpha\alpha'}$: Occupation matrix

$C_{2\alpha\beta\alpha'\beta'}$: 2-body correlation matrix

$C_{3\alpha\beta\gamma\alpha'\beta'\gamma'}$: 3-body correlation matrix

Stationary conditions

$$\frac{d}{dt}n_{\alpha\alpha'} = F_1 = \langle 0|[a_{\alpha'}^+ a_{\alpha}, H]|0\rangle = 0$$

$$\frac{d}{dt}C_{2\alpha\beta\alpha'\beta'} = F_2 = \langle 0|[a_{\alpha'}^+ a_{\beta'}^+ a_{\beta} a_{\alpha}, H]|0\rangle = 0$$

$$\frac{d}{dt}C_{3\alpha\beta\gamma\alpha'\beta'\gamma'} = F_3 = \langle 0|[a_{\alpha'}^+ a_{\beta'}^+ a_{\gamma'}^+ a_{\gamma} a_{\beta} a_{\alpha}, H]|0\rangle = 0$$



$n_{\alpha\alpha'}$

$C_{2\alpha\beta\alpha'\beta'}$

$C_{3\alpha\beta\gamma\alpha'\beta'\gamma'}$

Gradient method can be used

$$\begin{pmatrix} \frac{\delta F_1}{\delta n} & \frac{\delta F_1}{\delta C_2} & 0 \\ \frac{\delta F_2}{\delta n} & \frac{\delta F_2}{\delta C_2} & \frac{\delta F_2}{\delta C_3} \\ \frac{\delta F_3}{\delta n} & \frac{\delta F_3}{\delta C_2} & \frac{\delta F_3}{\delta C_3} \end{pmatrix} \begin{pmatrix} \Delta n \\ \Delta C_2 \\ \Delta C_3 \end{pmatrix} = \begin{pmatrix} \Delta F_1 \\ \Delta F_2 \\ \Delta F_3 \end{pmatrix} = \begin{pmatrix} 0 - F_1 \\ 0 - F_2 \\ 0 - F_3 \end{pmatrix}$$



$$\begin{pmatrix} n(N+1) \\ C_2(N+1) \\ C_3(N+1) \end{pmatrix} = \begin{pmatrix} n(N) \\ C_2(N) \\ C_3(N) \end{pmatrix} - \alpha \begin{pmatrix} \frac{\delta F_1}{\delta n} & \frac{\delta F_1}{\delta C_2} & 0 \\ \frac{\delta F_2}{\delta n} & \frac{\delta F_2}{\delta C_2} & \frac{\delta F_2}{\delta C_3} \\ \frac{\delta F_3}{\delta n} & \frac{\delta F_3}{\delta C_2} & \frac{\delta F_3}{\delta C_3} \end{pmatrix}^{-1} \begin{pmatrix} F_1(N) \\ F_2(N) \\ F_3(N) \end{pmatrix}$$

N : Iteration step

Relation to the variational principle

$$|\Phi'\rangle = e^{iF} |0\rangle$$

$$F = F^\dagger = \sum (f_{\alpha\alpha'} a_\alpha^\dagger a_{\alpha'} + f_{\alpha\beta\alpha'\beta'} a_\alpha^\dagger a_\beta^\dagger a_{\beta'} a_{\alpha'})$$

$$E = \langle 0 | e^{-iF} H e^{iF} | 0 \rangle \approx E_0 + i \sum \left(\langle 0 | [H, a_\alpha^\dagger a_{\alpha'}] | 0 \rangle, \langle 0 | [H, a_\alpha^\dagger a_\beta^\dagger a_{\beta'} a_{\alpha'}] | 0 \rangle \right) \begin{pmatrix} f_{\alpha\alpha'} \\ f_{\alpha\beta\alpha'\beta'} \end{pmatrix}$$

$$\frac{\delta E}{\delta f_{\alpha\alpha'}} = i \langle 0 | [H, a_\alpha^\dagger a_{\alpha'}] | 0 \rangle = 0$$

$$\frac{\delta E}{\delta f_{\alpha\beta\alpha'\beta'}} = i \langle 0 | [H, a_\alpha^\dagger a_\beta^\dagger a_{\beta'} a_{\alpha'}] | 0 \rangle = 0$$

Expressions for F_1 , F_2 and F_3

$$F_{1\alpha\alpha'}(n, C_2) = (\varepsilon_\alpha - \varepsilon_{\alpha'})n_{\alpha\alpha'} + \sum_{\lambda_1\lambda_2\lambda_3} \left\{ \langle \alpha\lambda_3 | v | \lambda_1\lambda_2 \rangle C_{2\lambda_1\lambda_2\alpha'\lambda_3} - C_{2\alpha\lambda_3\lambda_1\lambda_2} \langle \lambda_1\lambda_2 | v | \alpha'\lambda_3 \rangle \right\}$$

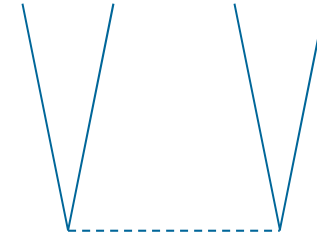
$$F_{2\alpha\beta\alpha'\beta'}(n, C_2, C_3) = (\varepsilon_\alpha + \varepsilon_\beta - \varepsilon_{\alpha'} - \varepsilon_{\beta'})C_{2\alpha\beta\alpha'\beta'} + B_{\alpha\beta\alpha'\beta'} + P_{\alpha\beta\alpha'\beta'} + H_{\alpha\beta\alpha'\beta'} + T_{\alpha\beta\alpha'\beta'}$$

$$F_{3\alpha\beta\gamma\alpha'\beta'\gamma'}(n, C_2, C_3) = (\varepsilon_\alpha + \varepsilon_\beta + \varepsilon_\gamma - \varepsilon_{\alpha'} - \varepsilon_{\beta'} - \varepsilon_{\gamma'})C_{3\alpha\beta\gamma\alpha'\beta'\gamma'} + \dots$$

B, P, H terms

$$B_{\alpha\beta\alpha'\beta'} = \sum_{\lambda_1\lambda_2\lambda_3\lambda_4} \langle \lambda_1\lambda_2 | v | \lambda_3\lambda_4 \rangle_A$$

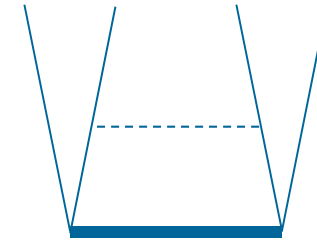
$$\times [(\delta_{\alpha\lambda_1} - n_{\alpha\lambda_1})(\delta_{\beta\lambda_2} - n_{\beta\lambda_2})n_{\lambda_3\alpha'}n_{\lambda_4\beta'} - (\alpha, \beta \Leftrightarrow \alpha', \beta')]$$



2p-2h excitations

$$P_{\alpha\beta\alpha'\beta'} = \sum_{\lambda_1\lambda_2\lambda_3\lambda_4} \langle \lambda_1\lambda_2 | v | \lambda_3\lambda_4 \rangle$$

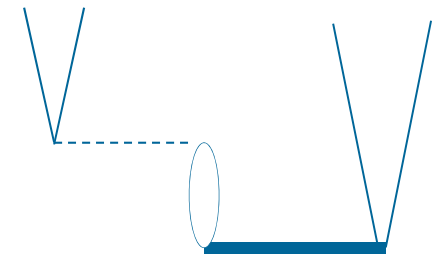
$$\times [(\delta_{\alpha\lambda_1}\delta_{\beta\lambda_2} - \delta_{\alpha\lambda_1}n_{\beta\lambda_2} - \delta_{\beta\lambda_2}n_{\alpha\lambda_1})C_{2\lambda_3\lambda_4\alpha'\beta'} - (\alpha, \beta \Leftrightarrow \alpha', \beta')]$$



p-p and h-h correlations

$$H_{\alpha\beta\alpha'\beta'} = \sum_{\lambda_1\lambda_2\lambda_3\lambda_4} \langle \lambda_1\lambda_2 | v | \lambda_3\lambda_4 \rangle_A$$

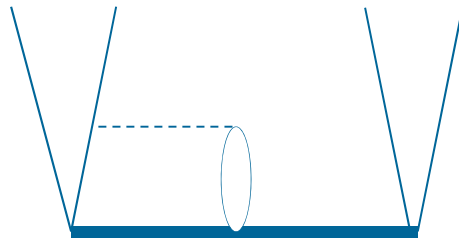
$$\times [(\delta_{\alpha\lambda_1}n_{\lambda_3\alpha'} - \delta_{\lambda_3\alpha'}n_{\alpha\lambda_1})C_{2\lambda_4\beta\lambda_2\beta'} + (\alpha, \beta \Leftrightarrow \alpha', \beta') - (\delta_{\beta\lambda_1}n_{\lambda_3\alpha'} - \delta_{\lambda_3\alpha'}n_{\beta\lambda_1})C_{2\lambda_4\alpha\lambda_2\beta'} - (\alpha, \beta \Leftrightarrow \alpha', \beta')]$$



P-h correlations

T term

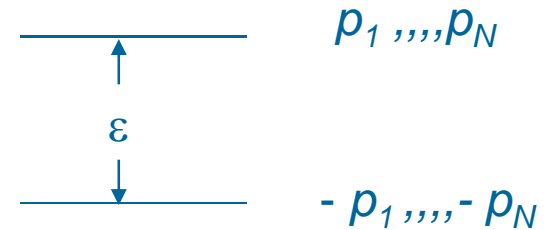
$$T_{\alpha\beta\alpha'\beta'} = \sum_{\lambda_1\lambda_2\lambda_3\lambda_4} \langle \lambda_1\lambda_2 | v | \lambda_3\lambda_4 \rangle$$
$$\times [\delta_{\alpha\lambda_1} C_{3\lambda_3\lambda_4\beta\alpha'\lambda_2\beta'} + \delta_{\beta\lambda_2} C_{3\lambda_4\lambda_3\alpha\beta'\lambda_1\alpha'}$$
$$- \delta_{\alpha'\lambda_3} C_{3\alpha\lambda_4\beta\lambda_1\lambda_2\beta'} - \delta_{\beta'\lambda_4} C_{3\alpha\lambda_3\beta\lambda_1\lambda_2\alpha'}]$$



interferes with p-p and p-h correlations

Application to Lipkin model

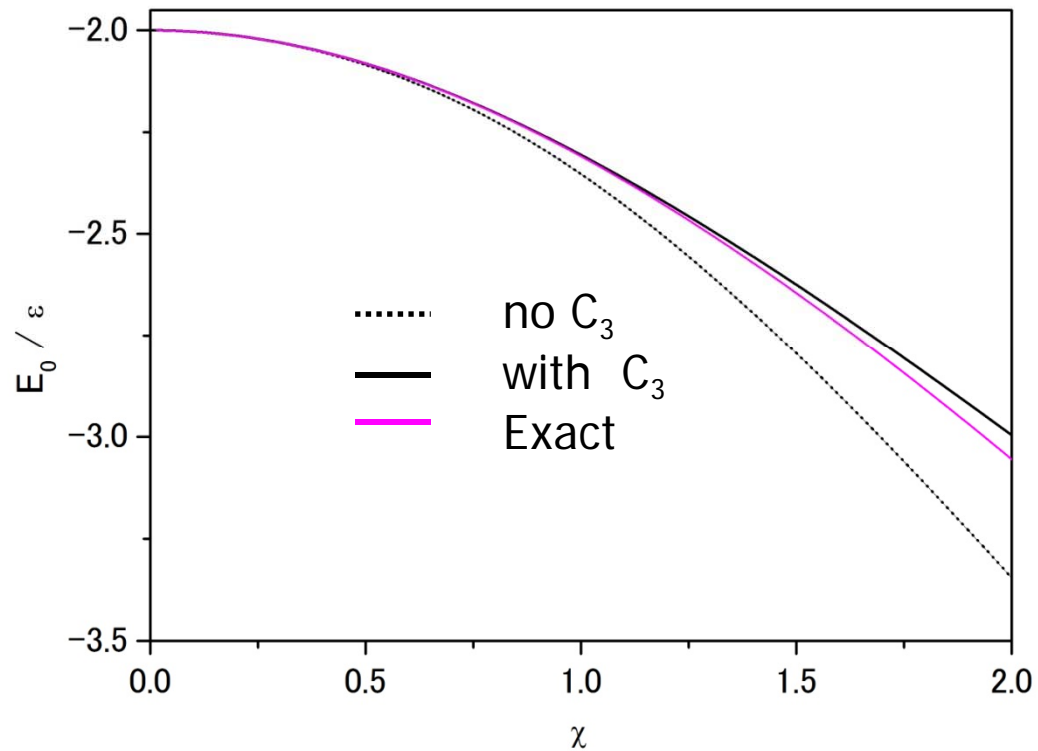
$$H = \varepsilon J_z + \frac{1}{2} V (J_+^2 + J_-^2)$$



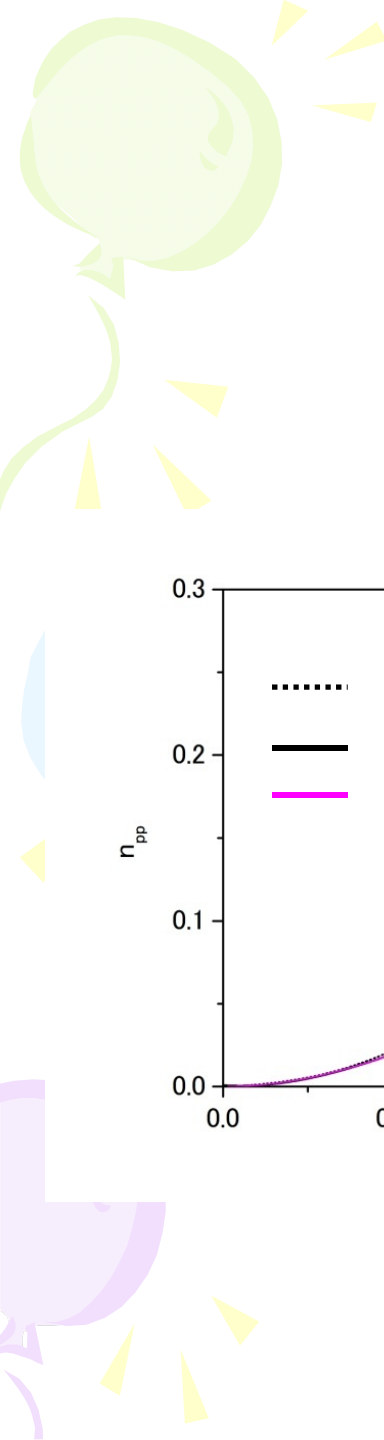
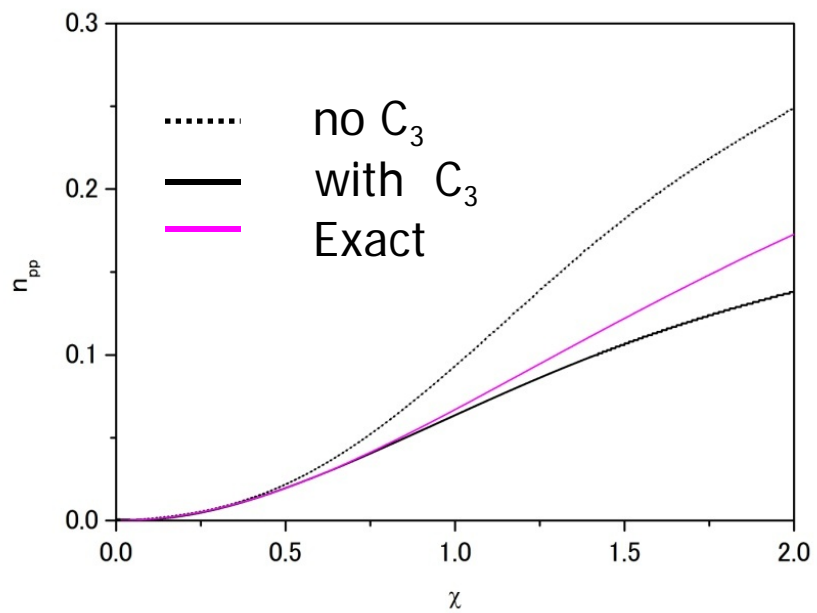
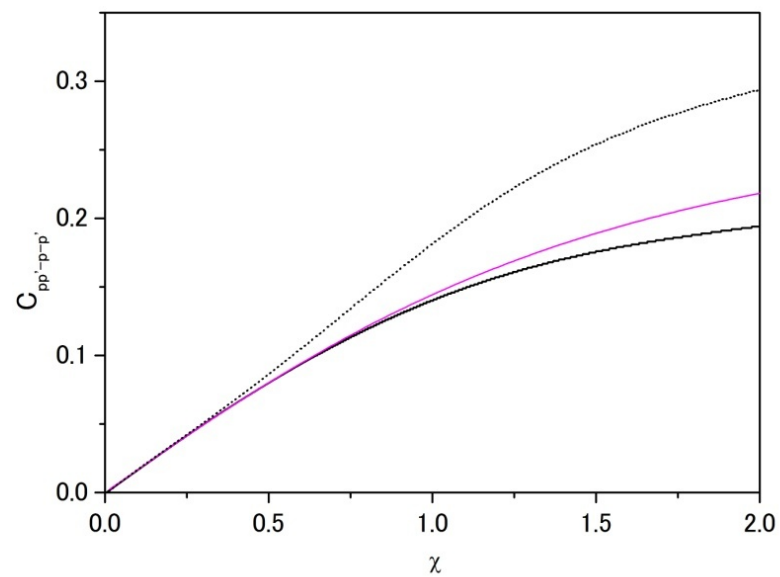
$$J_z = \frac{1}{2} \sum_p (a_p^+ a_p - a_{-p}^+ a_{-p}), \quad J_+ = J_-^+ = \sum_p a_p^+ a_{-p}$$

$$[J_+, J_-] = 2J_z, \quad [J_z, J_{\pm}] = \pm J_{\pm}$$

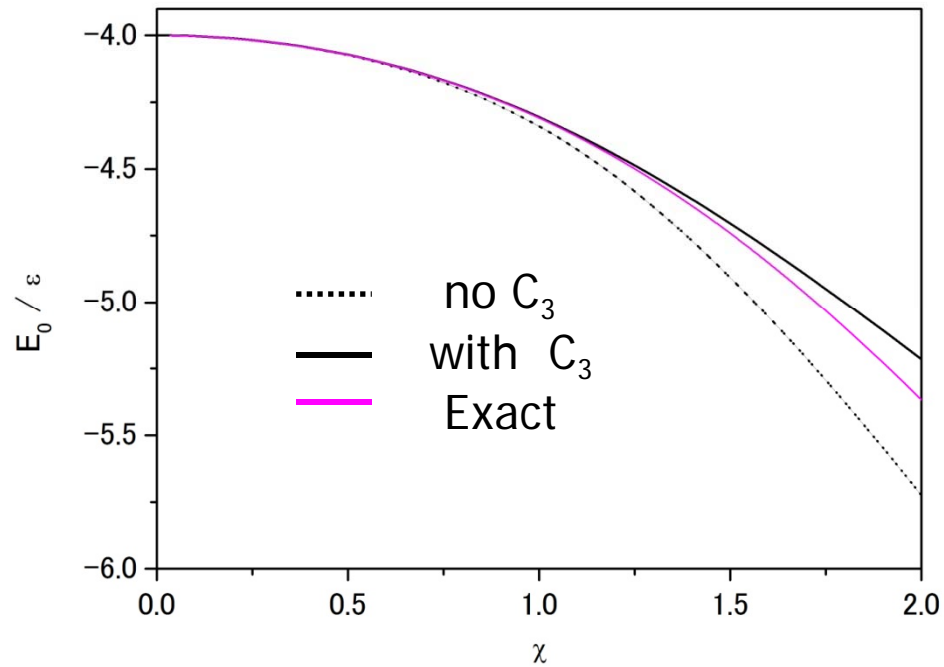
Ground-state energy (N=4)



$$\chi = (N-1)V/\epsilon$$

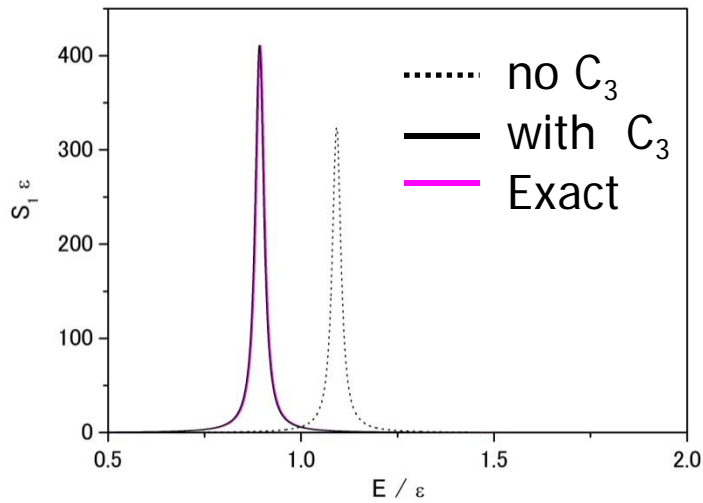

$$n_{pp}$$

$$C_{pp'-p-p'}$$


Ground-state energy (N=8)

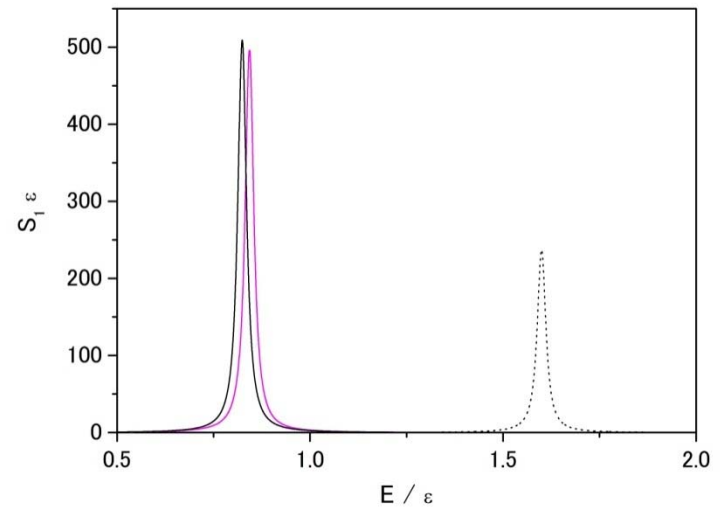


One-phonon state (N=4)

$$Q_{\mu}^{+} = J_{+} + J_{-}$$



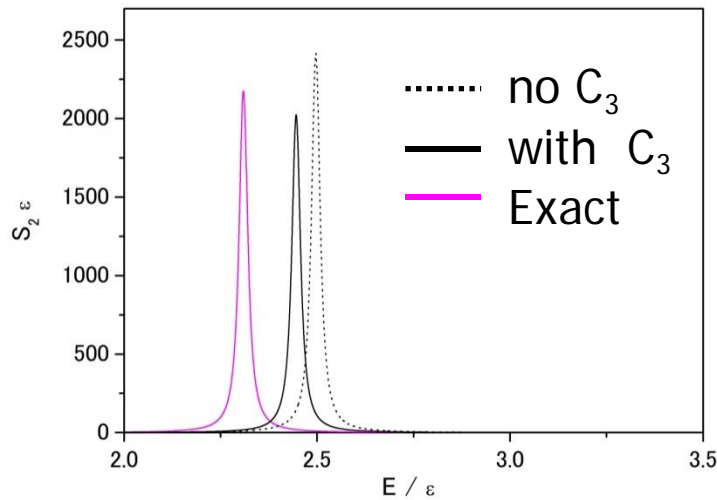
$$\chi = (N-1)V/\epsilon = 1$$



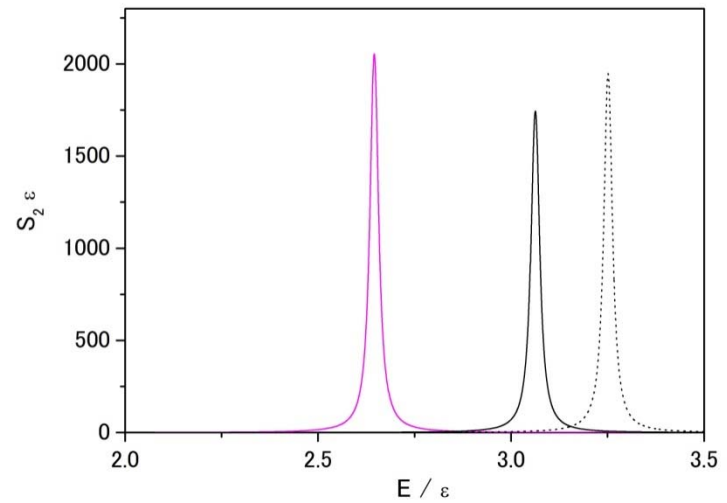
$$\chi = 1.5$$

2 phonon state (N=4)

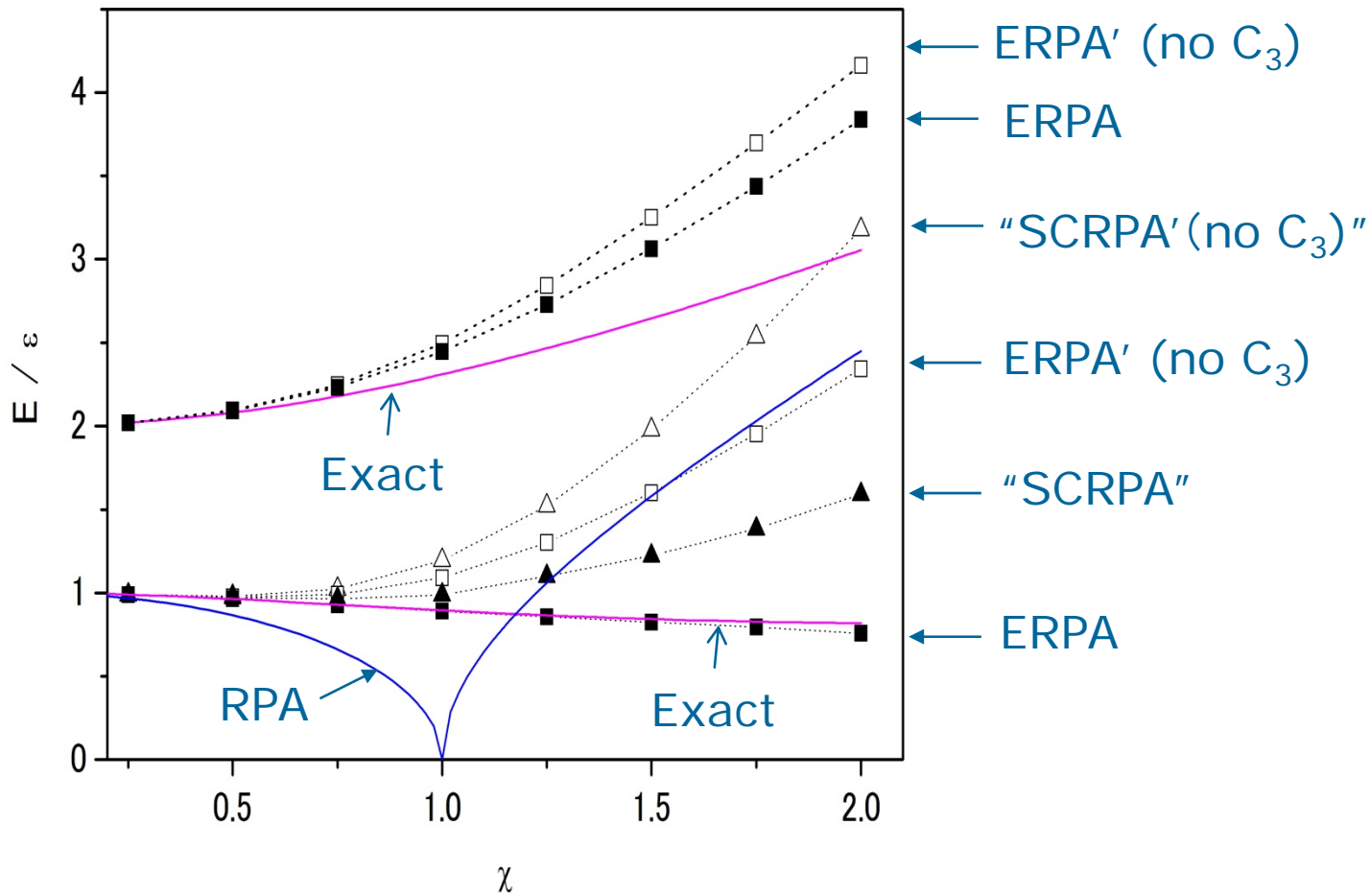
$$Q_{\mu}^{+} = (J_{+} + J_{-})^2$$



$\chi=1$



$\chi=1.5$





Summary

- 3-body g.s. correlations guarantee Hermiticity
- 3-body g.s. correlations drastically improve the results for the Lipkin model