

Dynamic correlation for strongly orthogonal geminals

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Uncorrelated strongly orthogonal electron pairs approximation

- Many-electron wavefunction is given as an antisymmetrized product of two-electron functions (strongly orthogonal geminals)

$$\Psi(x_1, \dots, x_N) = \hat{A} \prod_{P=1}^{N/2} \phi_P(x_{2P-1}, x_{2P})$$

- Each geminal is of the form (APSG)

$$\phi_P(\mathbf{x}_1, \mathbf{x}_2) = 2^{-1/2} \sum_{p \in P} c_p \psi_p(\mathbf{r}_1) \psi_p^*(\mathbf{r}_2) [\alpha(s_1) \beta(s_2) - \alpha(s_2) \beta(s_1)]$$

where

$$\forall_q \quad (c_q)^2 = n_q \in (0, 1)$$

- Geminals are strongly orthogonal

$$\forall_{Q \neq P} \quad \forall_{\mathbf{x}_1, \mathbf{x}'_1} \int \phi_P(\mathbf{x}_1, \mathbf{x}_2) \phi_Q(\mathbf{x}'_1, \mathbf{x}_2) d\mathbf{x}_2 = 0$$

PP-GVB (perfect-pairing generalized valence bond)

- Each geminal is of the form

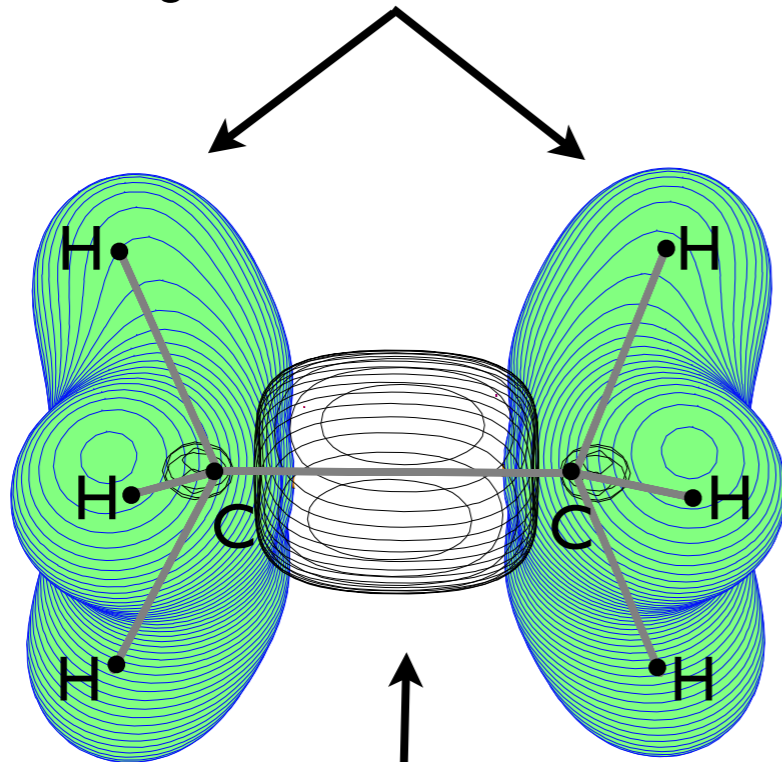
$$\forall_P \quad \phi_P^{GVB}(\mathbf{r}_1, \mathbf{r}_2) = c_{1P} \psi_{1P}(\mathbf{r}_1) \psi_{1P}(\mathbf{r}_2) + c_{2P} \psi_{2P}(\mathbf{r}_1) \psi_{2P}(\mathbf{r}_2)$$

where

$$\forall_q \quad (c_q)^2 = n_q \in (0, 1)$$

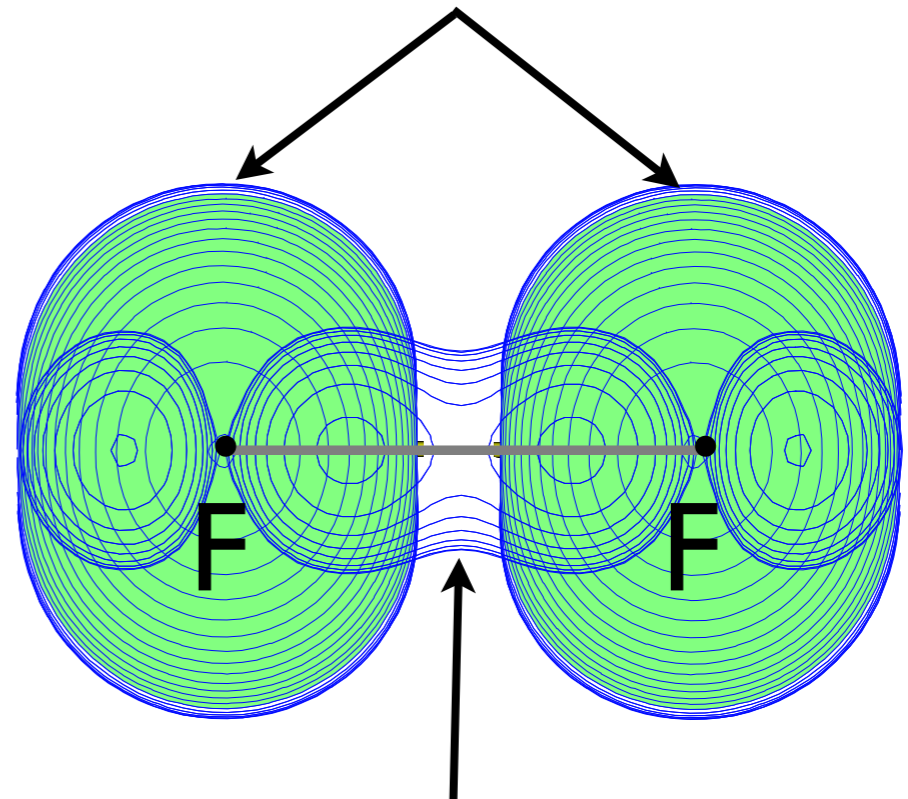
Localization on bonds and lone pairs of geminal densities

densities of geminals localized on C-H bonds



a geminal localized on bond

densities of geminals localized on atoms



a geminal localized on bond

Ground state energy functional in the APSG model

- The spin-summed expression for the electron energy takes a simple form

$$E^{APSG} = 2 \sum_p n_p h_{pp} + \sum_I \sum_{\substack{pq \\ I_p=I_q=I}} c_p c_q \langle pp|qq \rangle + \sum_{I \neq J} \sum_{\substack{pq \\ I_p=I, I_q=J}} n_p n_q [2 \langle pq|pq \rangle - \langle pq|qp \rangle]$$

intrageminal electron interaction

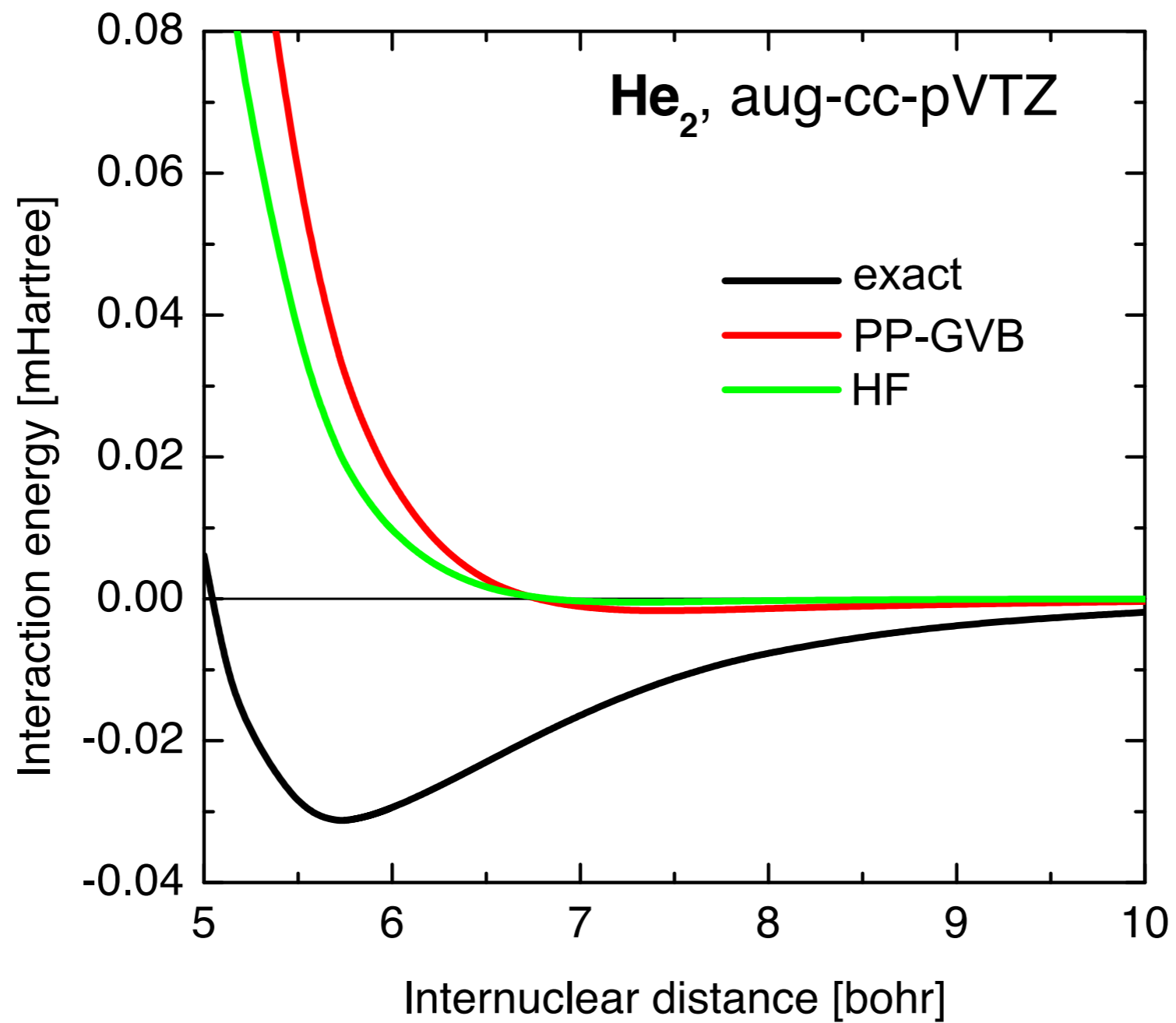
Coulomb and exchange intergeminal interaction

where I_p stands for the index of a geminal which the p th orbital belongs to.

- The ground state energy is obtained upon minimization with respect to (1) the orbitals, (2) the coefficients c_p under normalization constraint, and (3) Arai subspaces.

$$\forall_P \quad \sum_{p \in P} c_p^2 = 1$$

Lack of pair-pair correlation results in failure in capturing dispersion interaction



Inter-pair correlation energy

- Inter-domain correlation between I and J domains results from coupling density-density fluctuations

$$\rho_{IJ}^{(2)}(\mathbf{x}_1, \mathbf{x}_2) = \frac{1}{2} [\langle \delta \hat{\rho}_I(\mathbf{x}_1) \delta \hat{\rho}_J(\mathbf{x}_2) \rangle + \langle \delta \hat{\rho}_J(\mathbf{x}_1) \delta \hat{\rho}_I(\mathbf{x}_2) \rangle]$$

$$\delta \hat{\rho}_I(x) = \hat{\rho}_I(x) - \rho_I(x)$$

$$E_{corr}^{IJ} = \frac{1}{2} \sum_{\sigma_1, \sigma_2} \int \int \frac{\rho_{IJ}^{(2)}(\mathbf{x}_1, \mathbf{x}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|} d\mathbf{r}_1 d\mathbf{r}_2$$

K. Pernal, *J. Chem. Theory Comput.* **10**, 4332 (2014).

K. Pernal, *Phys. Chem. Chem. Phys.*, **18**, 21111 (2016).

Intergeminal correlation from the fluctuation-dissipation theorem for geminals

- It is more convenient to write the Inter-Geminal (IG) correlation in terms of the transition density matrix elements

$$(T_\nu)_{pq} = \langle 0 | \hat{a}_q^\dagger \hat{a}_p | \nu \rangle$$

$$E_{IGcorr} = \frac{1}{2} \sum_{\substack{I,J \\ I>J}} \sum_{pq \in I} \sum_{rs \in J} \sum_{\nu \neq 0} (T_\nu)_{qp} (T_\nu)_{rs}^* \langle pr | qs \rangle$$

Equation of motion

- By considering the Rowe's equation of motion and the **ERPA (extended random phase approximation)** excitation operator

$$\hat{O}_{ERPA}^\dagger = \sum_{p>q} (X_\nu)_{pq} (\hat{a}_{p\alpha}^\dagger \hat{a}_{q\alpha} + \hat{a}_{p\beta}^\dagger \hat{a}_{q\beta}) + \sum_{p>q} (Y_\nu)_{pq} (\hat{a}_{q\alpha}^\dagger \hat{a}_{p\alpha} + \hat{a}_{q\beta}^\dagger \hat{a}_{p\beta})$$

where p, q pertain to the natural spinorbitals, we have obtained equations for excitation energies

$$\begin{aligned} (\mathbf{A} + \mathbf{B}) (\mathbf{X} + \mathbf{Y}) &= \omega \mathcal{N} (\mathbf{Y} - \mathbf{X}) \\ (\mathbf{A} - \mathbf{B}) (\mathbf{Y} - \mathbf{X}) &= \omega \mathcal{N} (\mathbf{X} + \mathbf{Y}) \end{aligned}$$

The eigenvectors are related to the transition density matrix elements

$$\forall_{p>q} (n_p - n_q) [(Y_\nu)_{pq} - (X_\nu)_{pq}] = (T_\nu)_{pq} + (T_\nu)_{qp}$$

$$(T_\nu)_{pq} = \langle 0 | \hat{a}_q^\dagger \hat{a}_p | \nu \rangle$$

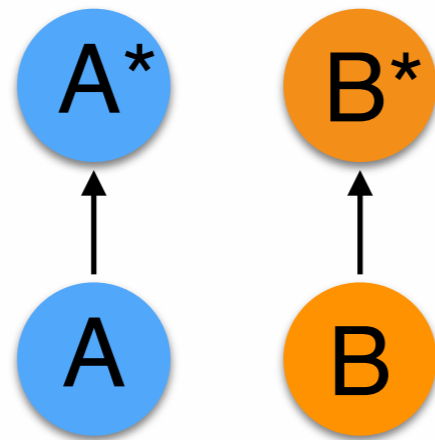
D. J. Rowe, *Rev. Mod. Phys.* **40**, 153 (1968).

K. Chatterjee and K. Pernal, *J. Chem. Phys.* **137**, 204109 (2012).

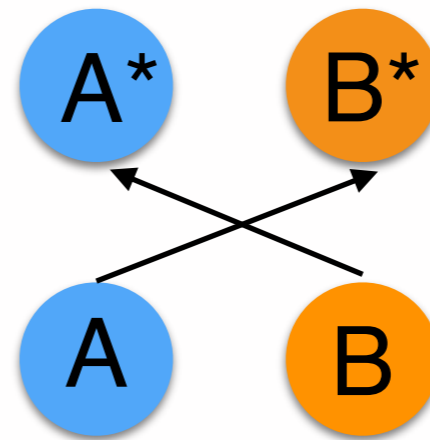
K. Pernal, K. Chatterjee, and P. H. Kowalski, *J. Chem. Phys.* **140**, 014101 (2014).

Obtained expression for the inter-pair correlation involves pairwise terms representing interactions of transition densities from A and B electron pairs

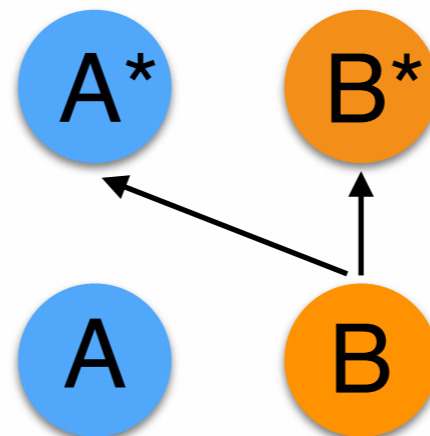
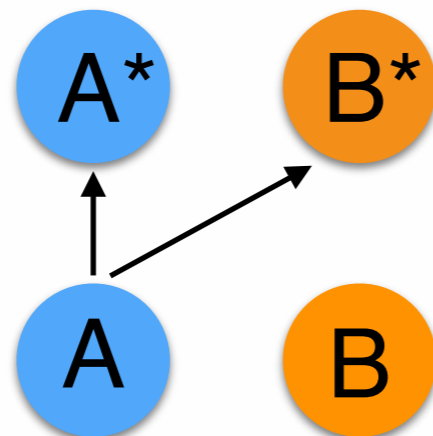
Dispersion



Exchange-Dispersion



Charge transfer



Many-body correlation terms in ERPA-GVB

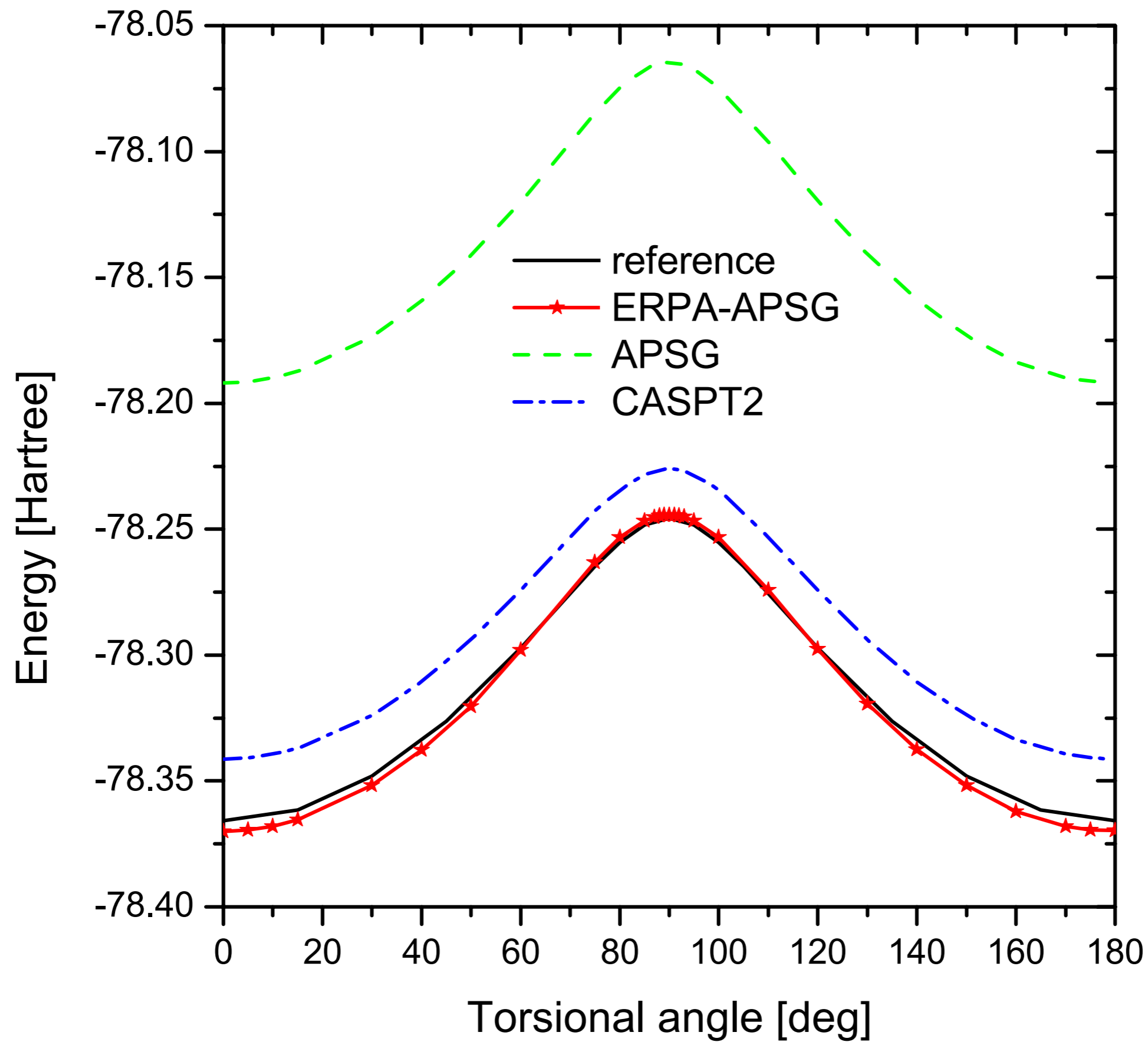
The idea of inter-pair (two-body interaction) has been generalized to include

- 1-body (intra-pair) correlation
- 2-body correlation (density fluctuations of two pairs are coupled)
- 3-body correlation (density fluctuations of three pairs are coupled)
- 4-body correlation (density fluctuations of four pairs are coupled)

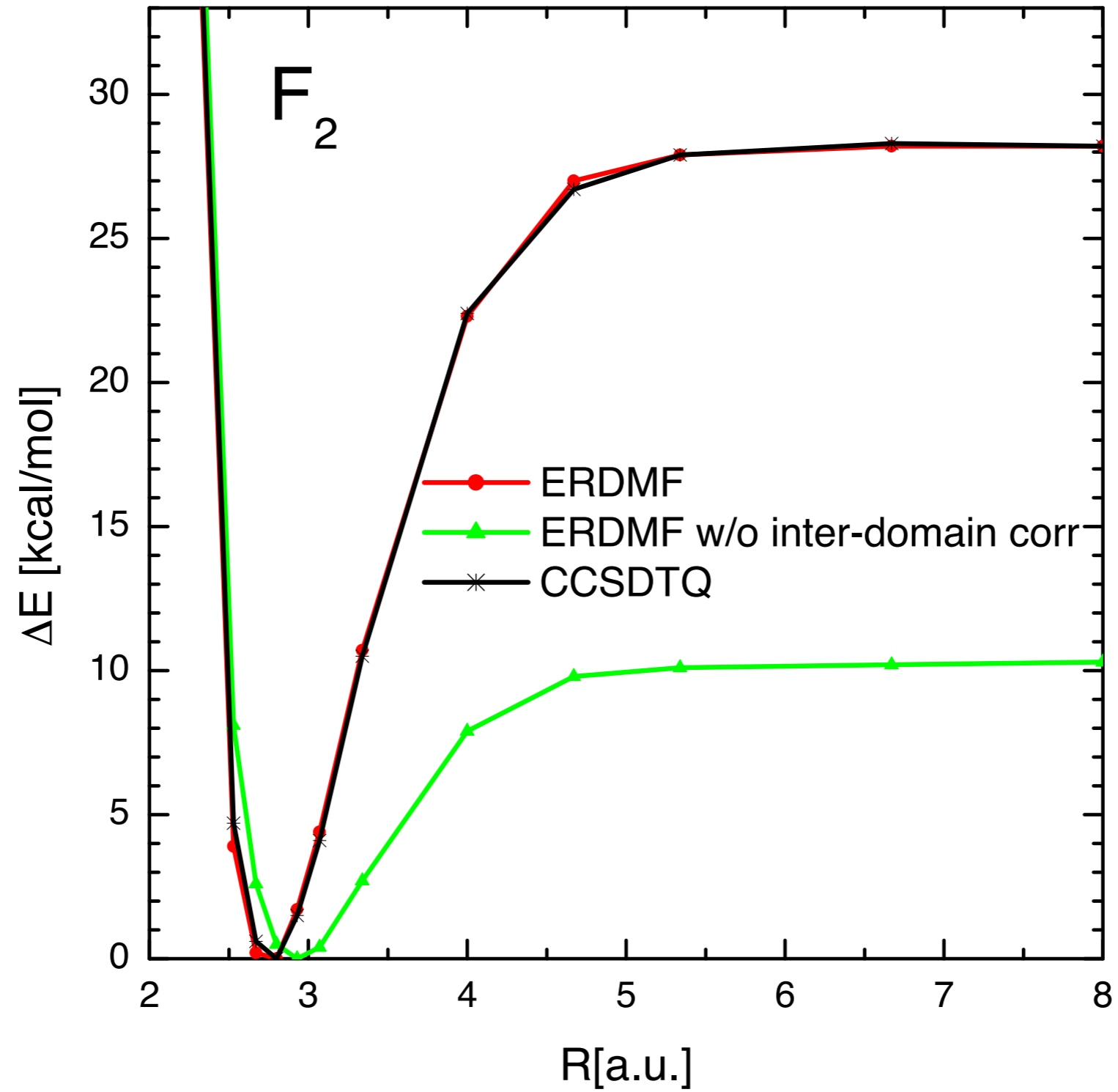
$$E_{corr}^{ERPA-GVB} = \sum_A^{1-body} E_A^{corr} + \sum_{AB}^{2-body} E_{AB}^{corr} + \sum_{ABC}^{3-body} E_{ABC}^{corr} + \sum_{ABCD}^{4-body} E_{ABCD}^{corr}$$

where A,B,C,D stand for geminals.

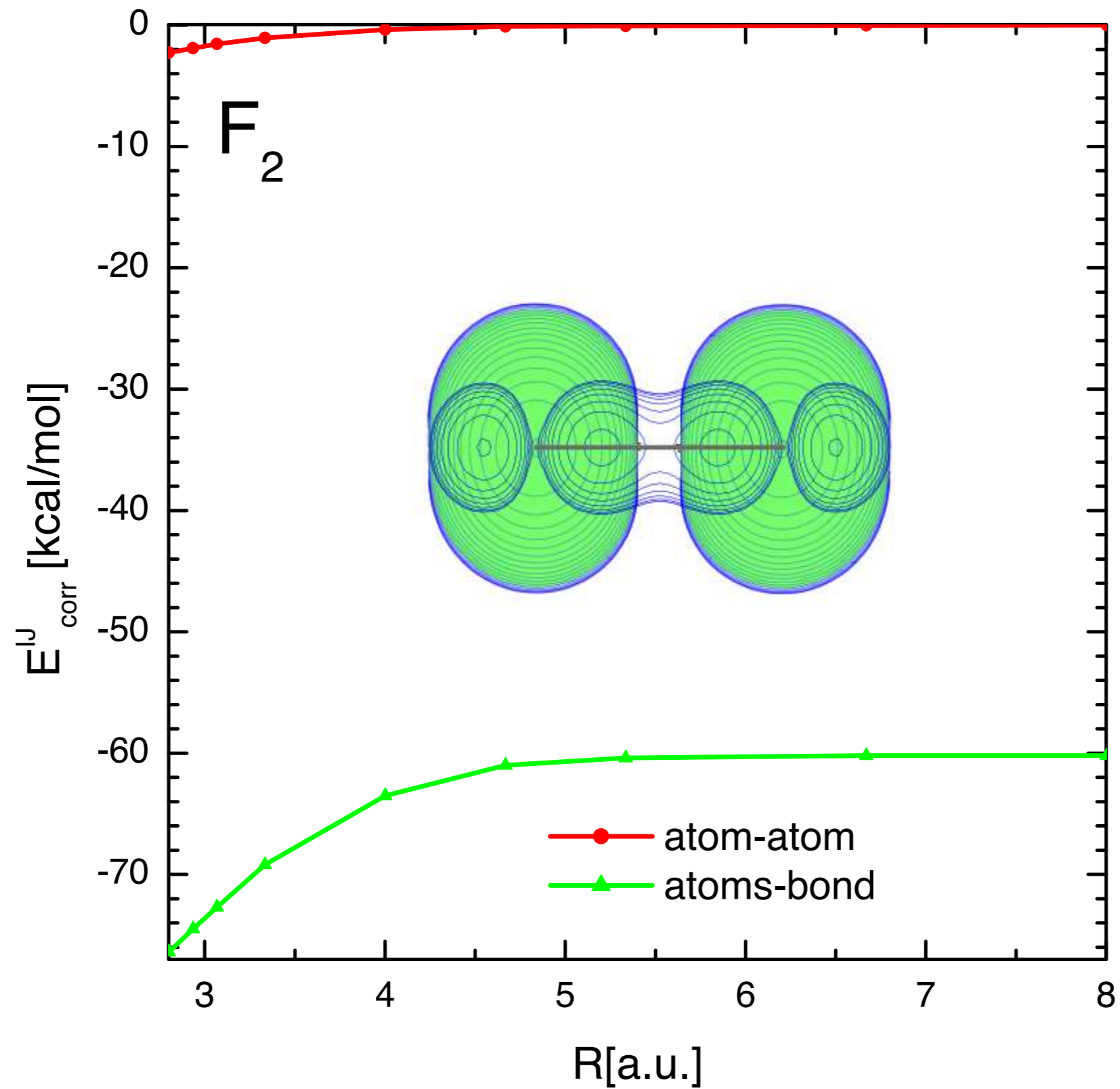
Twisting of the CC bond in ethylene



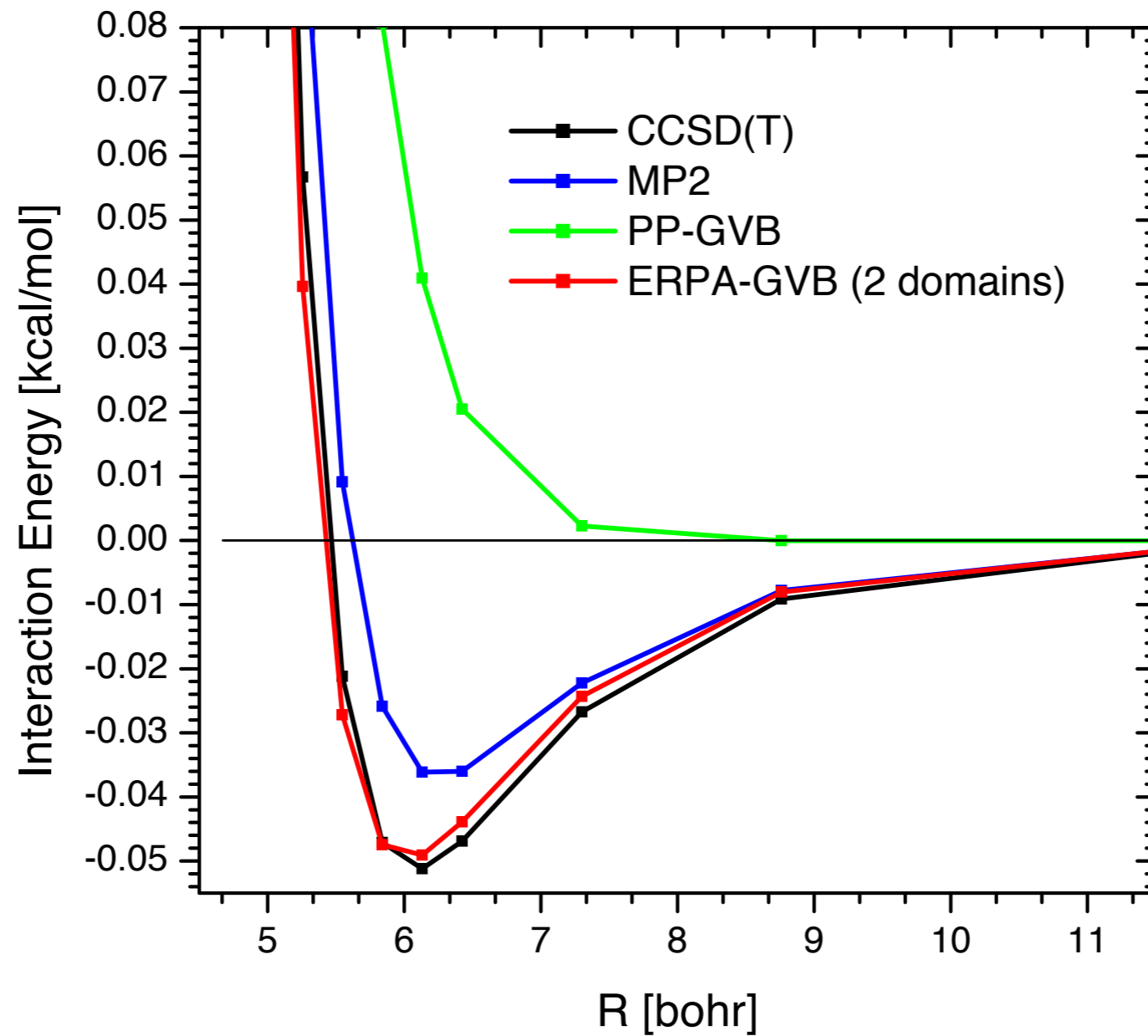
Dissociation curve of F₂ (cc-pVDZ basis set)



Dissociation curve of F₂ - inter-domain contributions



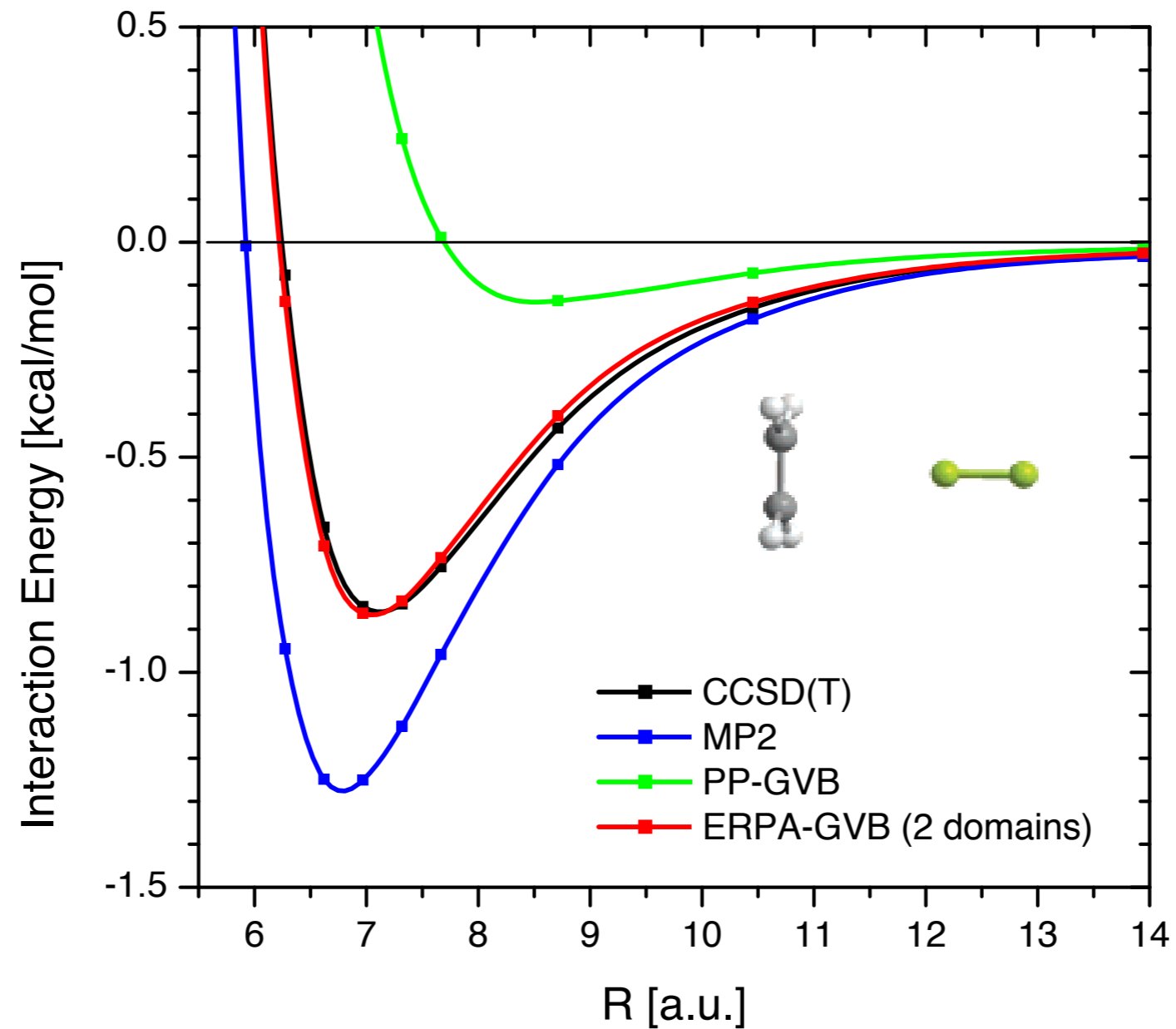
Ne₂ (aug-cc-pVDZ)



PP-GVB: no inter-domain correlation

ERPA-GVB (2 domains): intra- and inter-domain ERPA correlation added

C₂H₄--F₂ dimer (aug-cc-pVDZ)



PP-GVB: no intra and inter-domain ERPA correlation

ERPA-GVB (2 domains): intra- and inter-domain ERPA correlation added

Circular H_6 - breakdown for non-Lewis-structure molecule

