



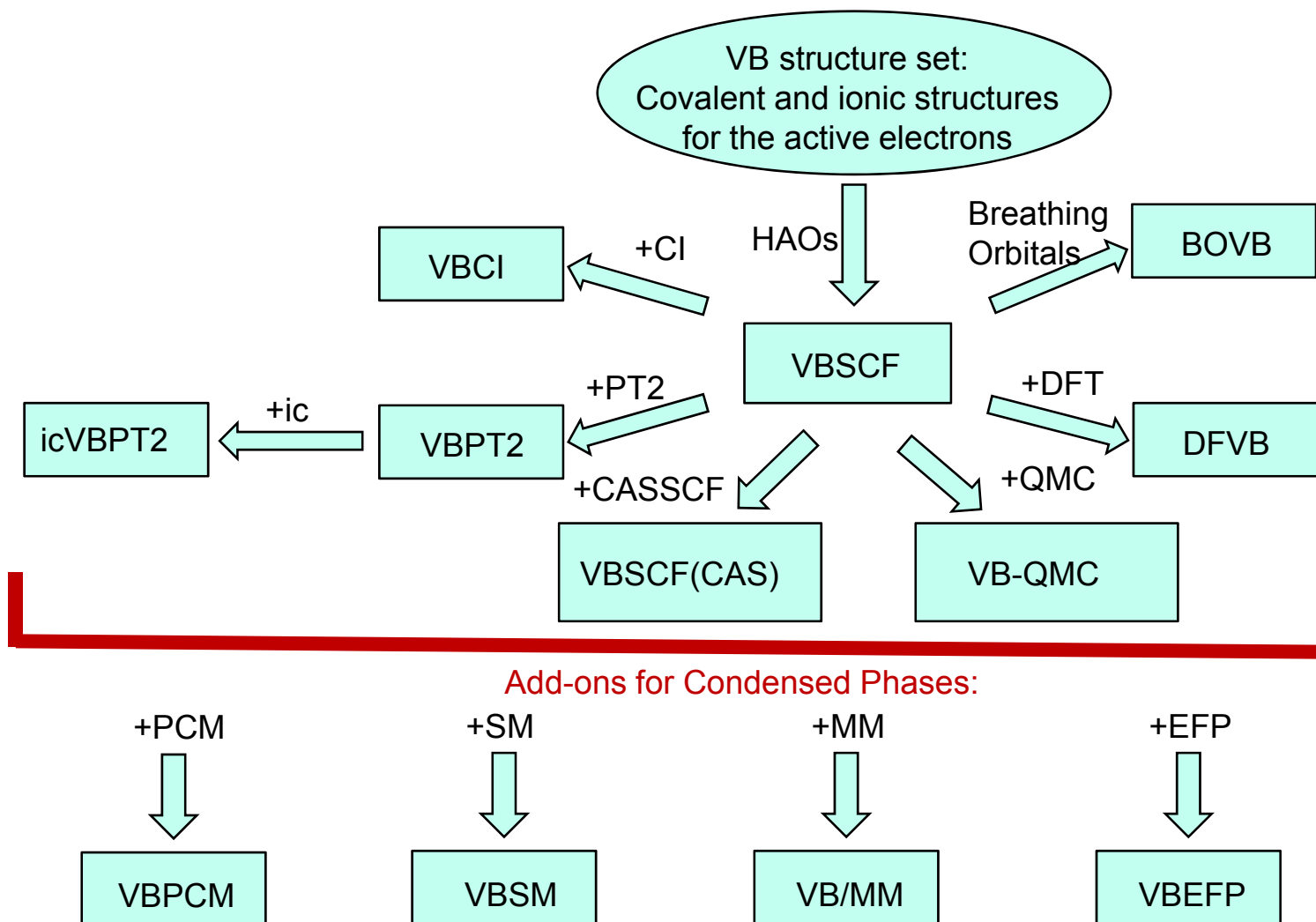
# Post-VBSCF Methods & VB Methods for Solvation Effect

Wei Wu

Department of Chemistry  
Xiamen University

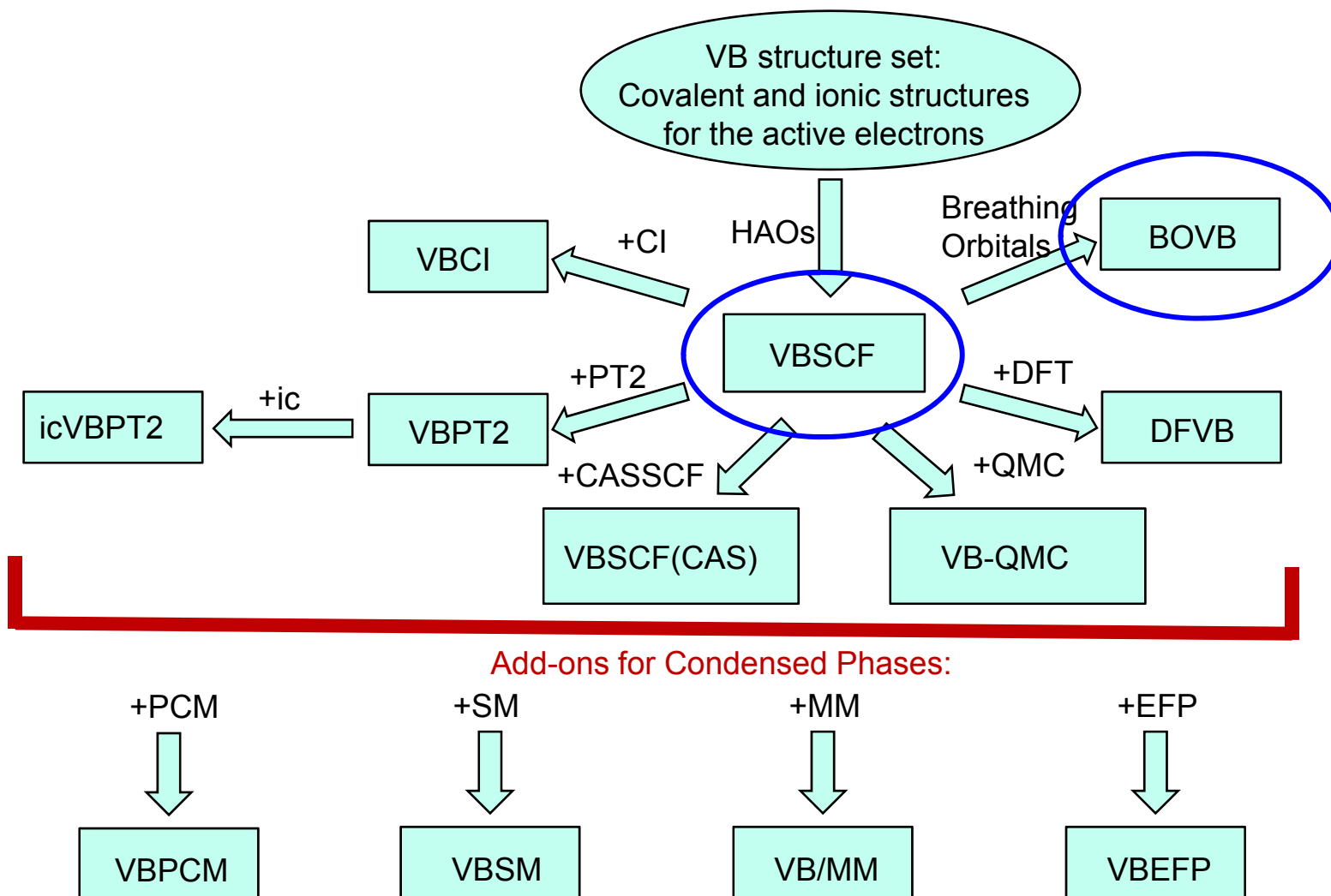
July 16, 2012 Paris





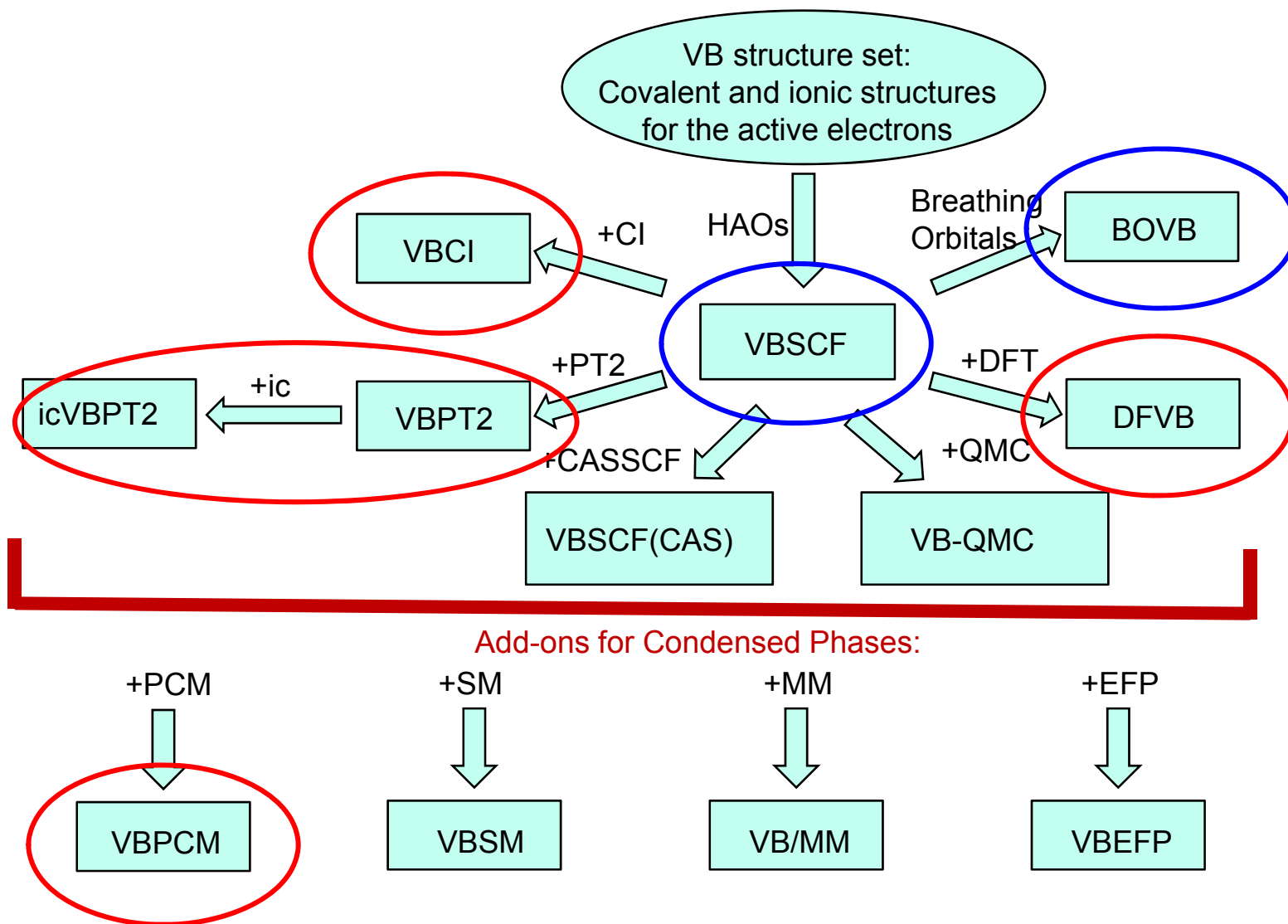
The tree of VB methods that are based on **classical** VB theory.

Wu; Su; Shaik; Hiberty, **Chem. Rev.** 2011. 111, 7557-7593.



The tree of VB methods that are based on **classical** VB theory.

Wu; Su; Shaik; Hiberty, **Chem. Rev.** 2011, 111, 7557-7593.



The tree of VB methods that are based on **classical** VB theory.

Wu; Su; Shaik; Hiberty, **Chem. Rev.** 2011. 111, 7557-7593.

# Outlines

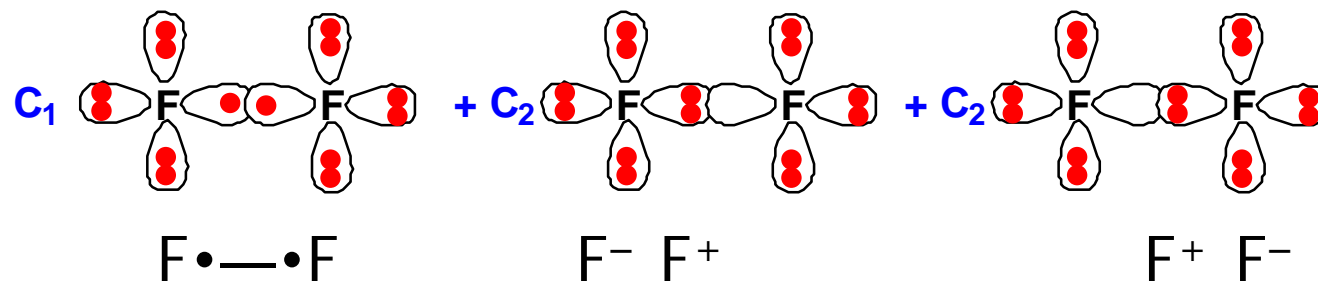
---

- VBCI
- VBPT2 (icVBPT2)
- DFVB
- VBPCM

# Valence Bond Self-Consistent Field (VBSCF)

---

Both structure **coefficients** and VB **orbitals** are optimized simultaneously to minimize the total energy.



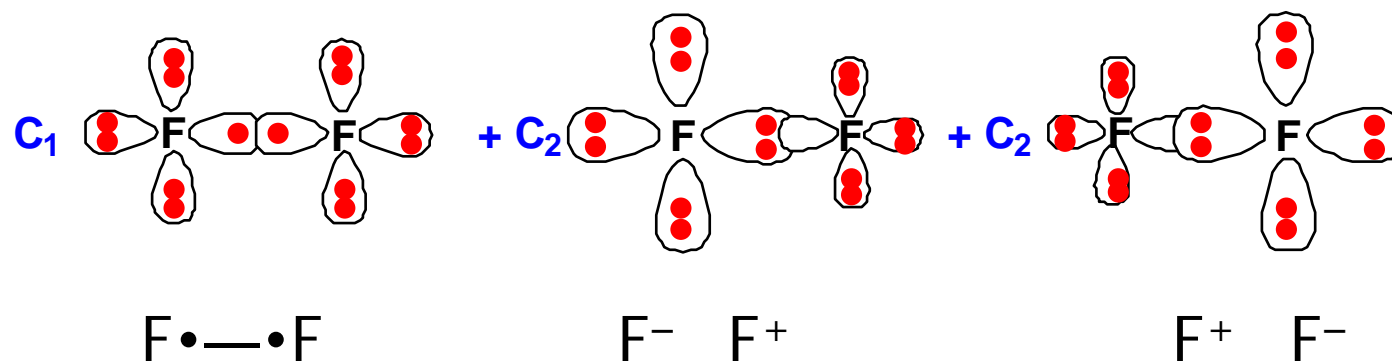
VBSCF provides qualitative correct description for bond breaking/forming, but its accuracy is still wanting.

VBSCF takes care of the **static** correlation, but lacks **dynamic** correlation.

van Lenthe; Balint-Kurti, *J. Chem. Phys.* 1983, 78, 5699.

# Breathing Orbital Valence Bond (BOVB)

- Different orbital sets for different VB structures



The degree of freedom introduces dynamic correlation with very compact wave function. It improves considerably the accuracy of the results.

Levels: L-BOVB; D-BOVB; SL-BOVB; SD-BOVB.

Hiberty, et al. *Chem. Phys. Lett.* **1992**, 189, 259.

TABLE 1: Bond Energies ( $D$ ) with Various Methods(kcal/mol)

molecule	$D^{\text{HF}}$	$D^{\text{B3LYP}}$	$D^{\text{CCSD}}$	$D^{\text{VBSCF}}$	$D^{\text{L-BOVB}}$	$D^{\text{VBCIS}}$	$D^{\text{VBCISD}}$
H <sub>2</sub>	84.6	111.7	105.9	95.8	96.0	96.0(11)	105.9(55)
LiH	32.5	57.2	49.5	42.4	43.0	42.8(27)	49.6(118)
HF	94.9	132.4	127.2	105.1	115.9	125.0(40)	126.0(274)
HCl	77.6	103.1	99.1	85.8	89.9	92.0(40)	98.0(274)
F <sub>2</sub>	-33.1	41.4	28.3	10.9	31.5	40.4(81)	33.9(1089)
Cl <sub>2</sub>	14.5	48.4	41.6	26.2	35.6	38.9(81)	42.1(1089)

Note that BOVB values may be improved by using higher BOVB levels.



## Post-HF in MO theory

For single-reference, post Hartree-Fock methods, such as CI, MP2, and CC;

For multi-reference, CASPT2, MRCI, etc.

**Can we have post-VBSCF methods?**

VBCI & VBPT2

## Post-VBSCF Wave Function

$$\left| \Psi^{\text{Post-VBSCF}} \right\rangle = \left| \Psi^{\text{VBSCF}} \right\rangle + \sum_R C_R \left| \Phi_R \right\rangle$$

$\Phi_R$  Excited VB structures

## ***How to define localized VB orbitals (HAOs)?***

$$\{\chi_{\alpha}\} = \{\chi_1^A, \chi_2^A, \dots, \chi_{m_A}^A; \chi_1^B, \chi_2^B, \dots, \chi_{m_B}^B; \chi_1^C, \chi_2^C, \dots, \chi_{m_C}^C; \dots\}$$

A: atom or fragment

Localized occupied VB orbitals

$$\phi_i^A = \sum_{\beta} c_{\beta i}^A \chi_{\beta}^A$$

*Occupied VB orbitals are obtained by VBSCF calculations.  
Virtual orbitals are not necessary in VBSCF.*

# How to define virtual orbitals?

---

Scheme 1:

Schmidt Orthogonalization  
to occupied orbitals on their  
own fragments

$$\chi'_{\mu,A} = \left( \chi_{\mu} - S_{\mu\nu} T_n^{\mu} S^{nm} \phi_m \right)_{\mu,\nu \in A}$$

Localized

Used in VBCI

Scheme 2:

Schmidt Orthogonalization  
to all occupied orbitals

$$\chi'_{\mu} = \chi_{\mu} - S_{\mu\nu} T_n^{\mu} S^{nm} \phi_m$$

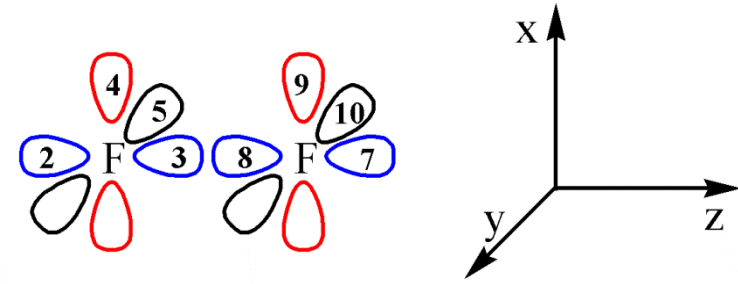
Delocalized

Used in VBPT2

Note:

- Virtual basis functions are not linearly independent.
- To remove the redundant basis functions, one may diagonalize the overlap matrix of the transformed basis functions.

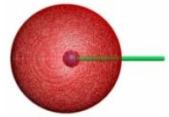
# Occupied orbitals of F<sub>2</sub>



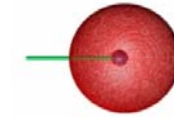
$$1\phi_{\sigma}^A = 1s$$



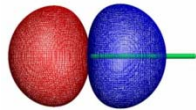
$$6\phi_{\sigma}^B = 1s$$



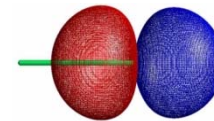
$$2\phi_{\sigma}^A = 2s$$



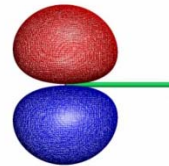
$$7\phi_{\sigma}^B = 2s$$



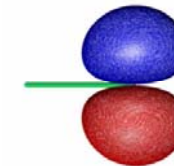
$$3\phi_{\sigma}^A = 2p_z$$



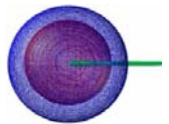
$$8\phi_{\sigma}^B = 2p_z$$



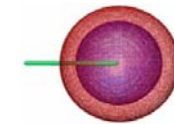
$$4\phi_{\pi}^A = 2p_x$$



$$9\phi_{\pi}^B = 2p_x$$

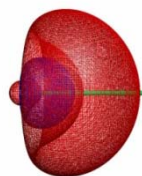


$$5\phi_{\pi}^A = 2p_y$$

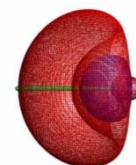


$$10\phi_{\pi}^B = 2p_y$$

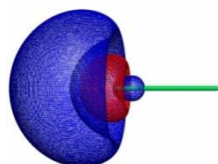
## Virtual Orbitals from Scheme 1



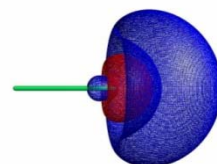
$11\phi_{\sigma}^A$



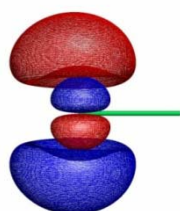
$15\phi_{\sigma}^B$



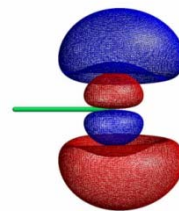
$12\phi_{\sigma}^A$



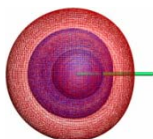
$16\phi_{\sigma}^B$



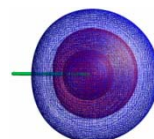
$13\phi_{\pi}^A$



$17\phi_{\pi}^B$



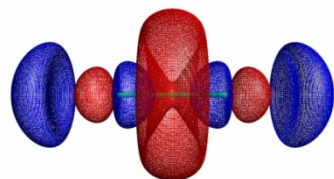
$14\phi_{\pi}^A$



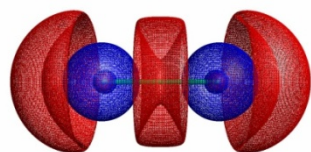
$18\phi_{\pi}^B$



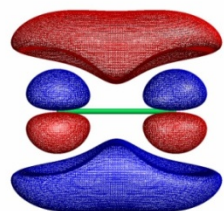
## Virtual Orbitals from Scheme 2



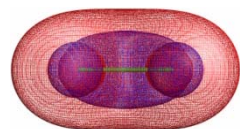
$11\phi_\sigma$



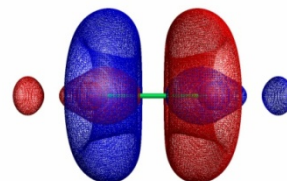
$12\phi_\sigma$



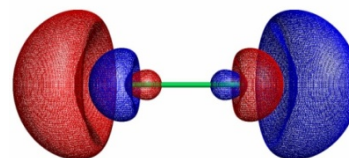
$13\phi_\pi$



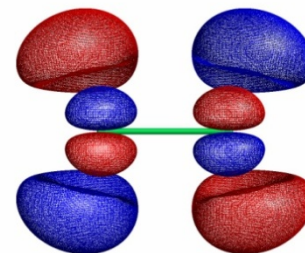
$14\phi_\pi$



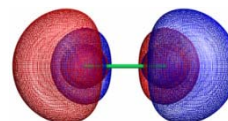
$15\phi_\sigma$



$16\phi_\sigma$



$17\phi_\pi$



$18\phi_\pi$

# **1. Valence Bond Configuration Interaction (VBCI)**

# Excited VB structures

---

$$\Psi^{\text{VBSCF}} = \sum_K C_K^{\text{SCF}} \Phi_K^{\text{SCF}}$$

An excited VB structure  $\Phi_K^i$  is built from  $\Phi_K^{\text{SCF}}$  by replacing occupied  $\phi_j^A$  with virtual orbital  $\phi_\alpha^A$ .

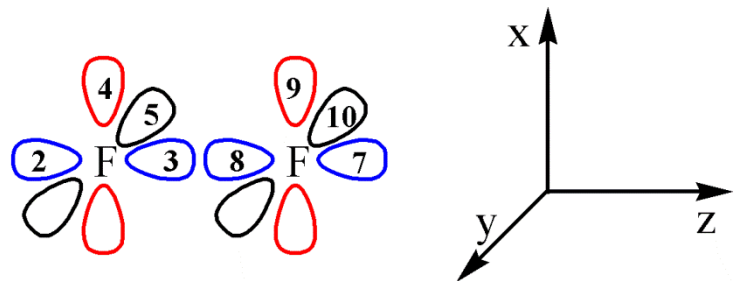
$\Phi_K^i$  and  $\Phi_K^{\text{SCF}}$  describe the **same** classical VB structure.

By collecting all  $\Phi_K^i$ , we have a wave function corresponding to VB structure  $K$ .

$$\Phi_K^{\text{CI}} = \sum_i C_{Ki}' \Phi_K^i$$

**Corresponding to a VB structure.**

# VBCI Wave Function



CI wave function

Covalent

Ionic

Fundamental  
(VBSCF)

$$3\phi_{\sigma}^A \text{ — } 8\phi_{\sigma}^B$$

$$3\phi_{\sigma}^A \ 3\phi_{\sigma}^A$$

+

Single Excitation

$$3\phi_{\sigma}^A \text{ — } 15\phi_{\sigma}^B$$

$$3\phi_{\sigma}^A \ 11\phi_{\sigma}^A$$

+

Double Excitation

$$12\phi_{\sigma}^A \text{ — } 15\phi_{\sigma}^B$$

$$12\phi_{\sigma}^A \ 11\phi_{\sigma}^A$$

+

....

$$\begin{aligned}\Psi^{VBCI} &= \sum_K C_K^{CI} \Phi_K^{CI} \\ &= \sum_K \sum_i C_{Ki} \Phi_K^i\end{aligned}$$

$$\Psi^{VBSCF} = \sum_K C_K^{SCF} \Phi_K^{SCF}$$

$$E^{VBCI} = \frac{\sum_{K,L} \sum_{i,j} C_{Ki} C_{Lj} \langle \Phi_K^i | H | \Phi_L^j \rangle}{\sum_{K,L} \sum_{i,j} C_{Ki} C_{Lj} \langle \Phi_K^i | \Phi_L^j \rangle}$$

## Condensed matrix elements and weights

$$H_{KL}^{CI} = \sum_{i,j} C_{Ki} C_{Lj} \langle \Phi_K^i | H | \Phi_L^j \rangle$$

$$M_{KL}^{CI} = \sum_{i,j} C_{Ki} C_{Lj} \langle \Phi_K^i | \Phi_L^j \rangle$$

$$W_K = \sum_i W_{Ki}$$

$$W_{Ki} = \sum_{L,j} C_{Ki} \langle \Phi_K^i | \Phi_L^j \rangle C_{Lj}$$

***All formulas are similar to those of VBSCF, and compact.***

## Levels of VBCI Method

---

VBCI(A,I), A= D, S; I = D, S

A = Active electrons that are involved in a chemical process

I = Inactive electrons that are NOT involved in ...

VBCI(D,D) = VBCISD

VBCI(S,S) = VBCIS

VBCI(D,S)

The “inactive” electrons play an indirect role in a chemical process, and the dynamic correlation of inactive electrons is quasi constant during the process.

## Davidson Correction of VBCISD

---

Size inconsistency problem is one of the most undesirable drawbacks in truncated CI methods.

$$\Delta E_Q = (1 - \sum_K W_K) \Delta E_D$$

to estimate the contribution of quadruple excitations that are product of double excitations



TABLE 1: Bond Energies ( $D$ ) with Various Methods(kcal/mol)

molecule	$D^{\text{HF}}$	$D^{\text{B3LYP}}$	$D^{\text{CCSD}}$	$D^{\text{VBSCF}}$	$D^{\text{L-BOVB}}$	$D^{\text{VBCIS}}$	$D^{\text{VBCISD}}$
H <sub>2</sub>	84.6	111.7	105.9	95.8	96.0	96.0(11)	105.9(55)
LiH	32.5	57.2	49.5	42.4	43.0	42.8(27)	49.6(118)
HF	94.9	132.4	127.2	105.1	115.9	125.0(40)	126.0(274)
HCl	77.6	103.1	99.1	85.8	89.9	92.0(40)	98.0(274)
F <sub>2</sub>	-33.1	41.4	28.3	10.9	31.5	40.4(81)	33.9(1089)
Cl <sub>2</sub>	14.5	48.4	41.6	26.2	35.6	38.9(81)	42.1(1089)

Table 2. Weights of the VB structures for various VB methods

	H2	LiH	HF	HCl	F <sub>2</sub>	Cl <sub>2</sub>
VBSCF	0.8074	0.8919	0.5829	0.6708	0.7933	0.6949
	0.0963	0.0028	0.0106	0.0725	0.1033	0.1531
	0.0963	0.1053	0.4066	0.2567	0.1033	0.1520
BOVB	0.7606	0.8756	0.5182	0.6486	0.7012	0.6513
	0.1197	-0.0062	0.0637	0.0998	0.1494	0.1744
	0.1197	0.1306	0.4181	0.2516	0.1494	0.1744
VBCISD	0.7727	0.8799	0.5515	0.6547	0.7395	0.6637
	0.1137	-0.0137	0.0191	0.0847	0.1305	0.1681
	0.1137	0.1338	0.4293	0.2606	0.1300	0.1682

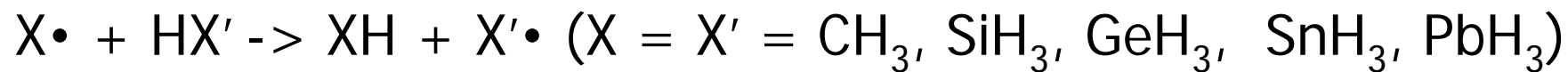
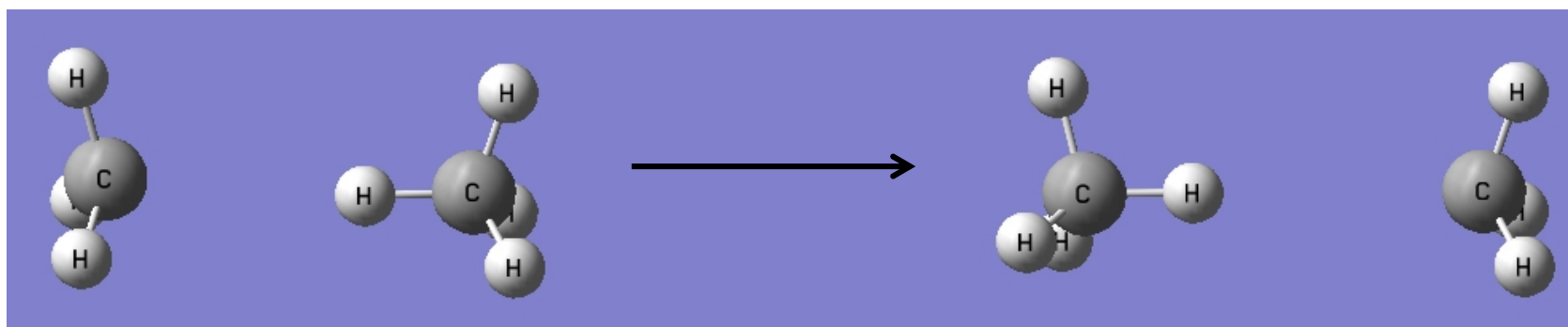
VBCI method provides significant improvement in accuracy from VBSCF method but still preserves the important feature of VB method.

**Table 1. Bond Dissociation Energies Calculated with Valence Bond Methods<sup>180</sup>**

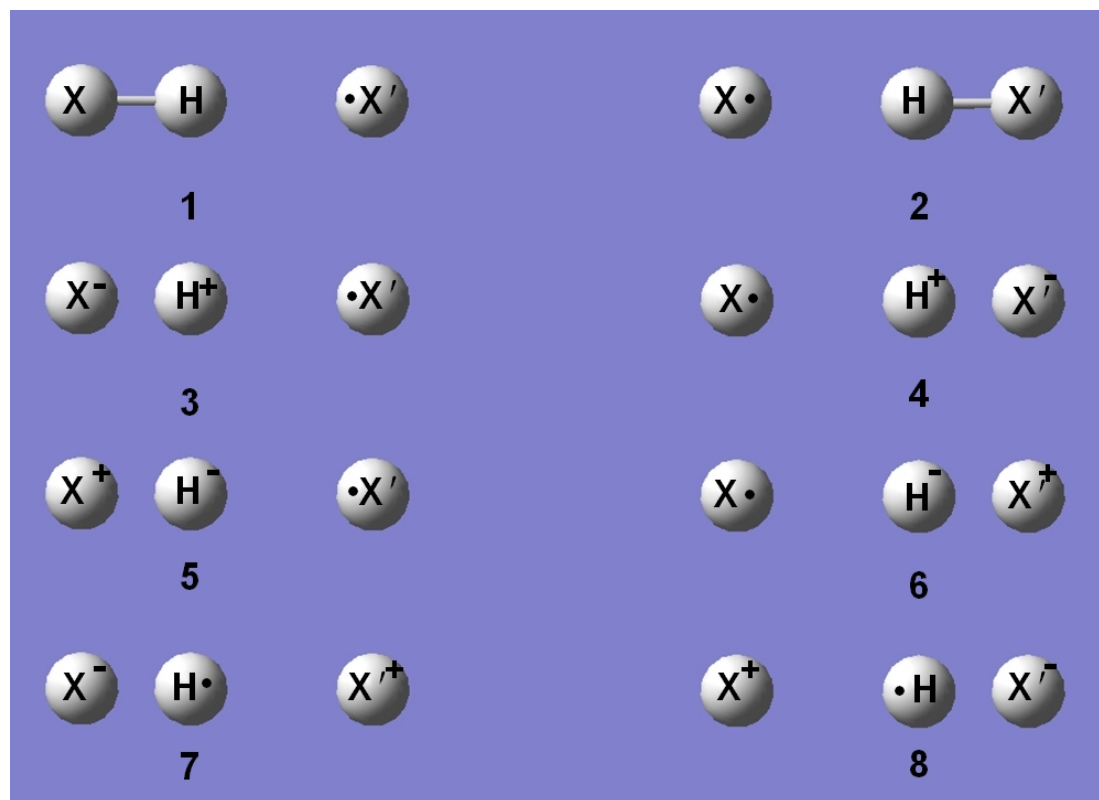
bond	basis set	$D_e$ (kcal/mol)			
		BOVB	VBCISD <sup>a</sup>	CCSD(T)	exptl
F–F	6-31G*	36.2	32.3	32.8	
	cc-pVTZ	37.9	36.1	34.8	38.3
Cl–Cl	6-31G*	40.0	41.6	40.5	
	cc-pVTZ	50.0	56.1	52.1	58.0
Br–Br	6-31G*	41.3	44.1	41.2	
	cc-pVTZ	44.0	50.0	48.0	45.9
F–Cl	6-31G*	47.9	49.3	50.2	
	cc-pVTZ	53.6	58.8	55.0	60.2
H–H	6-31G**	105.4	105.4	105.9	109.6
Li–Li	6-31G*	20.9	21.2	21.1	24.4
H <sub>3</sub> C–H	6-31G**	105.7	113.6	109.9	112.3
H <sub>3</sub> C–CH <sub>3</sub>	6-31G*	94.7	90.0	95.6	96.7
HO–OH	6-31G*	50.8	49.8	48.1	53.9
H <sub>2</sub> N–NH <sub>2</sub>	6-31G*	68.5	70.5	66.5	75.4 ± 3
H <sub>3</sub> Si–H	6-31G**	93.6	90.2	91.8	97.6 ± 3
H <sub>3</sub> Si–F	6-31G*	140.4 <sup>b</sup>	151.1	142.6	160 ± 7
H <sub>3</sub> Si–Cl	6-31G*	102.1	101.2	98.1	113.7 ± 4

<sup>a</sup> With Davidson correction. <sup>116</sup> <sup>b</sup> Two-structure calculations ( $\text{H}_3\text{Si}^-\text{F}^+$  is omitted).

## Hydrogen Abstraction Reactions



VB structures (3 electrons/3 orbitals system)



( $X = X' = CH_3, SiH_3, GeH_3, SnH_3, PbH_3$ )

Table 4. 6-31G\* Barriers for the Hydrogen Eaxhange Reactions,  $X\cdot + X'H \rightarrow XH + X'\cdot$  ( $X = CH_3, SiH_3, GeH_3, SnH_3, PbH_3$ ) (kcal/mol)

Molecule	HF	CCSD	VBSCF	BOVB	VBCIS	VBCISD	VBCI(D,S)	VBCIPT
CH <sub>3</sub>	35.1	26.5	33.0	23.1	26.7(266)	25.8(4156)	28.8(2788)	25.5(765)
SiH <sub>3</sub>	25.2	19.3	25.5	19.1	19.2(266)	19.7(4156)	19.4(2788)	19.0(794)
GeH <sub>3</sub>	22.0	16.6	25.5	18.0	18.9(266)	18.1(4156)	17.9(2788)	17.0(791)
SnH <sub>3</sub>	18.5	13.5	20.5	14.9	15.5(266)	15.3(4156)	15.1(2788)	14.1(565)
PbH <sub>3</sub>	15.2	13.0	17.3	12.3	12.7(266)	12.5(4156)	12.3(2788)	11.5(573)
CH <sub>3</sub> <sup>a</sup>							19.7(17.3)[17.7]	

The seven  $\sigma$  valence electrons are set into CI window, where three are active, and four are inactive.

This barrier was calculated with the mixed basis set and no valence orbitals are frozen. The result 19.7 is obtained with the set of 8 fundamental VB structures (Scheme 3). In parentheses is CCSD(T) barrier. In the square brackets is the result obtained by adding the VB structures that can be generated by admitting all the s-electrons (seven) into the active space. There are additional 32 structures and their effect was evaluated without orbital optimization.

(aug-cc-pVDZ for C, cc-pVDZ for the exchanged H atom and 6-31G for H atoms)

# Summary 1

---

- VBCI method covers dynamic correlation with a minimal number of effective structures .
- The VBCIS results are in good agreement with those of BOVB.
- The VBCISD results match CCSD results.

## **2. Valence Bond Second Order Perturbation Theory (VBPT2)**



As one of post-VBSCF methods, VBCI incorporates dynamic correlation without further SCF procedure.

However, though the VBCI space is much smaller than those of MO-based methods. **VBCI method is computational demanding.**

*Can we have a VB method that is accurate and cheap?*

# Valence Bond Second Order Perturbation Theory (VBPT2)

---

$$|\Psi\rangle = |\Psi^{(0)}\rangle + |\Psi^{(1)}\rangle \quad |\Psi^{(0)}\rangle = |\Psi^{SCF}\rangle = \sum_K C_K^{SCF} |\Phi_K\rangle$$

$$|\Psi^{(1)}\rangle = \sum_{R \in V_{SD}} C_R^{(1)} |\Phi_R\rangle$$

$$\langle \Psi^{(0)} | \Psi^{(1)} \rangle = 0 \quad \langle \Psi^{(0)} | \Psi \rangle = 1$$

## Excited VB structures:

Excited structures are generated from the VBSCF structures by replacing occupied orbitals with virtual orbitals.

Scheme 2 is used for virtual orbitals, which are not localized. Thus, excited structures are not “classical”.

## The zeroth-order Hamilton:

$$\hat{H}_0 = \hat{P}_0 \hat{F} \hat{P}_0 + \hat{P}_K \hat{F} \hat{P}_K + \hat{P}_{SD} \hat{F} \hat{P}_{SD} + \dots$$

$$\hat{F} = \sum_i \hat{f}(i)$$

$$\hat{f}(i) = \hat{h}(i) + \sum_{m,n} D_{mn}^{SCF} (\hat{J}_{nm}(i) - \hat{K}_{nm}(i))$$

$$f_{pq} = h_{pq} + \sum_{m,n} D_{mn}^{SCF} \left[ (pq|mn) - \frac{1}{2} (pm|qn) \right]$$

The first-order wave function:  $\Psi^{(1)} = \sum_{R \in V_{SD}} C_R^{(1)} |\Phi_R\rangle$

$$\mathbf{C}^{(1)} = (\mathbf{H}_0^{11} - E^{(0)} \mathbf{M}^{11})^{-1} \mathbf{H}^{10} \mathbf{C}^{(0)}$$

$$(H_0^{11})_{RS} = \langle \Phi_R | \hat{H}_0 | \Phi_S \rangle$$

$$(H^{01})_{KR} = \langle \Phi_K | \hat{H} | \Phi_R \rangle$$

$$(H^{10})_{RK} = \langle \Phi_R | \hat{H} | \Phi_K \rangle$$

$$(M^{11})_{RS} = \langle \Phi_R | \Phi_S \rangle$$

The second-order energy:

$$E^{(2)} = \mathbf{C}^{(0)\dagger} \mathbf{H}^{01} (\mathbf{H}_0^{11} - E^{(0)} \mathbf{M}^{11})^{-1} \mathbf{H}^{10} \mathbf{C}^{(0)}$$

The most time-consuming part:

$$(\mathbf{H}_0^{11} - E^{(0)}\mathbf{M}^{11})^{-1}$$

which is block diagonalized, due to the block-orthogonality between different orbital blocks.

If the occupation numbers of inactive or virtual orbitals are different in the two excited structures, the corresponding matrix element is zero.

VBPT2 is much cheaper than VBCI.

The structure weights in VBPT2 method:

$$\Phi_K^{PT} = N_K (\Phi_K + \sum_R X_{RK} \Phi_R) \quad \mathbf{X} = (\mathbf{H}_0^{11} - E^{(0)} \mathbf{M}^{11})^{-1} \mathbf{H}^{10}$$

$$\Psi^{VBPT} = \sum_K C_K^{PT} \Phi_K^{PT}$$

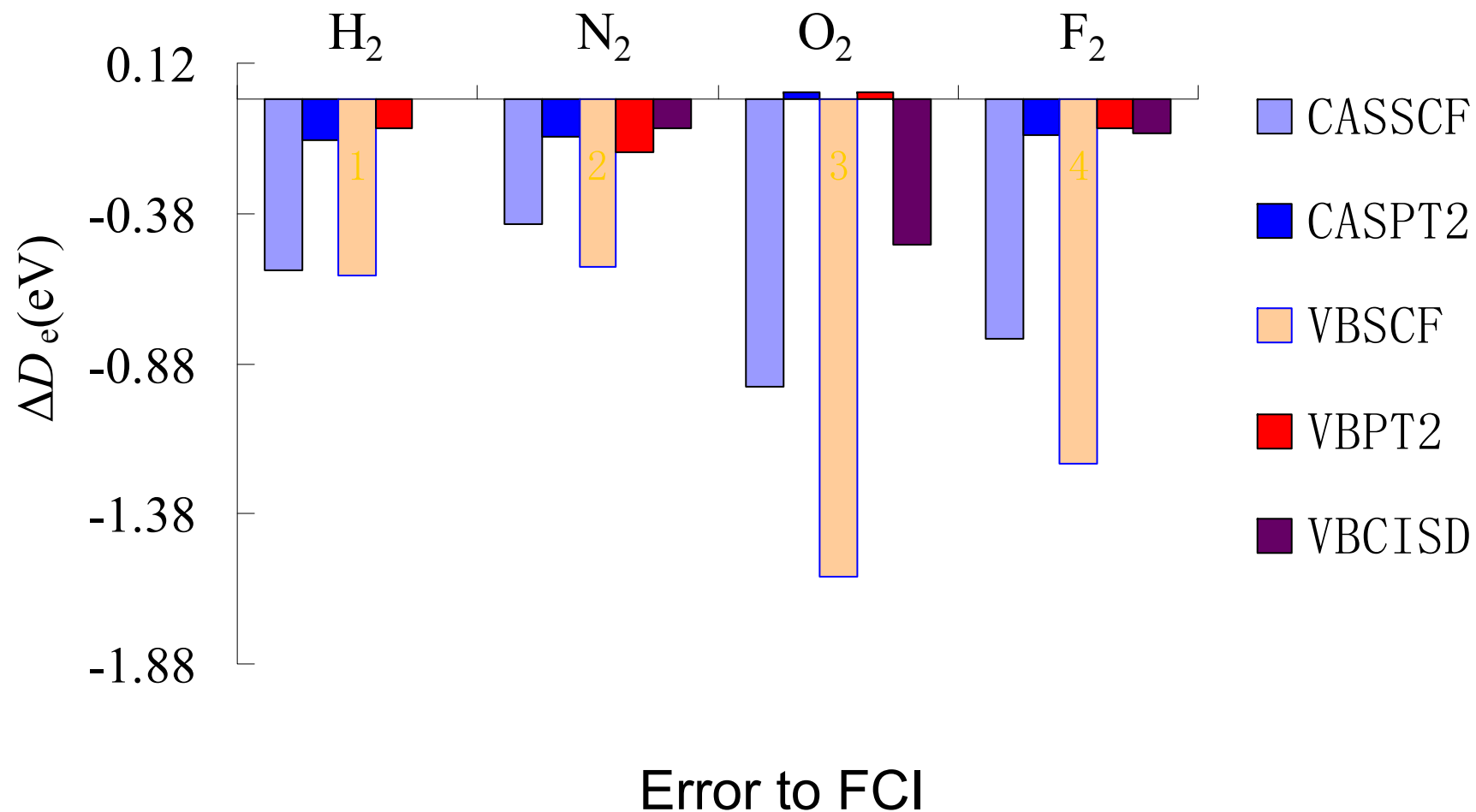
$$C_K^{PT} = C_K^{(0)} / N_K \quad N_K = (1 + \sum_{R,S} X_{RK} M_{RS}^{11} X_{SK})^{-1/2}$$

$$W_K^{PT} = \sum_L C_K^{PT} M_{KL}^{PT} C_L^{PT}$$

$$E_{KL}^{(2)} = \sum_{R,S} H_{KR}^{01} (\mathbf{H}_0^{11} - E^{(0)} \mathbf{M}^{11})_{RS}^{-1} H_{SL}^{10}$$

$$H_{KL}^{PT} = H_{KL}^{SCF} + E_{KL}^{(2)}$$

# Bond Dissociation Energy (Error to FCI)



**Table 6.** Weights of the VB structures and resonance energy (RE) of  $F_2$ , at an F-F distance of  $2.70 a_0$

Method	$W_{cov}$	$W_{ion1}/W_{ion2}$	RE (kcal/mol)
VBSCF	0.785	0.108	46.28
L-BOVB	0.695	0.152	71.46
VBCISD	0.730	0.135	60.68
VBPT2	0.778	0.111	60.02



The VBPT2 results are in very good agreement with CASPT2.

The VBPT2 results are in very good agreement with CASPT2.

### Disadvantage of VBPT2

The dimension of excited VB structures increases dramatically, when the number of VB structures in VBSCF is large.

### Solution

Internally contraction strategy

$$|\Psi\rangle = |\Psi^{(0)}\rangle + |\Psi^{(1)}\rangle$$

$$|\Psi^{(0)}\rangle = |\Psi^{\text{SCF}}\rangle = \sum_K C_K^{(0)} |\Phi_K\rangle$$

The first-order wave function: a linear combination of the internally constructed excited configurations.

$$|\Psi^{(1)}\rangle = \sum_{pqrs} t_{pr}^{qs} |\Psi_{qs}^{pr}\rangle = \sum_{pqrs} t_{pr}^{qs} E_s^{\bullet r} E_q^{\bullet p} |\Psi^{(0)}\rangle$$

Spin averaged orbital excitation operator

$$E_q^{\bullet p} = a_{q\alpha}^+ a^{p\alpha} + a_{q\beta}^+ a^{p\beta}$$

The expansion coefficients of first-order wave function

$$\mathbf{t} = (\mathbf{H}_0^{11} - E^{(0)}\mathbf{M}^{11})^{-1} \mathbf{V}$$

The second-order energy

$$E^{(2)} = \mathbf{V}^+ (\mathbf{H}_0^{11} - E^{(0)}\mathbf{M}^{11})^{-1} \mathbf{V}$$

$$(H_0^{11})_{\alpha\beta} = \langle \Psi_\alpha^{(1)} | \hat{H}_0 | \Psi_\beta^{(1)} \rangle$$

$$V_\alpha = \langle \Psi_\alpha^{(1)} | \hat{H} | \Psi^{(0)} \rangle$$

$$(M^{11})_{\alpha\beta} = \langle \Psi_\alpha^{(1)} | \Psi_\beta^{(1)} \rangle$$

## Various types of excitations

$$|\Psi_{xy}^{ij}\rangle = E_y^{\bullet j} E_x^{\bullet i} |\Psi^{(0)}\rangle$$

$$|\Psi_{ta}^{ij}\rangle = E_a^{\bullet j} E_t^{\bullet i} |\Psi^{(0)}\rangle$$

$$|\Psi_{ab}^{ij}\rangle = E_b^{\bullet j} E_a^{\bullet i} |\Psi^{(0)}\rangle$$

$$|\Psi_{xy}^{ti}\rangle = E_y^{\bullet i} E_x^{\bullet t} |\Psi^{(0)}\rangle$$

$$|\Psi_{xa}^{ti}\rangle = E_a^{\bullet i} E_x^{\bullet t} |\Psi^{(0)}\rangle$$

$$|\Psi_{ab}^{ti}\rangle = E_b^{\bullet i} E_a^{\bullet t} |\Psi^{(0)}\rangle$$

$$|\Psi_{ax}^{tu}\rangle = E_x^{\bullet u} E_a^{\bullet t} |\Psi^{(0)}\rangle$$

$$|\Psi_{ax}^{tu}\rangle = E_x^{\bullet u} E_a^{\bullet t} |\Psi^{(0)}\rangle$$

Singlet excitation operator

$$|\Psi_{ax}^{ti}\rangle = E_x^{\bullet i} E_a^{\bullet t} |\Psi^{(0)}\rangle$$

Orbitals:

Inactive.  $i, j, k, l$  (orthogonal)

Active.  $t, u, v, w, x, y, z$  (non-orthogonal)

Virtual.  $a, b, c$

Any.  $p, q, r, s$

Any occupied.  $m, n$

$$\begin{aligned}
& \langle \Psi_{x_1 y_1}^{ij} | \hat{H}_0 | \Psi_{x_2 y_2}^{ij} \rangle = (E_{inact}^{(0)} - \varepsilon_i - \varepsilon_j) \langle \Psi_{x_1 y_1}^{ij} | \Psi_{x_2 y_2}^{ij} \rangle \\
& + \left\{ (4 s_{x_2 x_1} s_{y_2 y_1} - 2 s_{y_2 x_1} s_{x_2 y_1}) E_{act}^{(0)} \right. \\
& + 4 f_{x_2 x_1} s_{y_2 y_1} + 4 f_{y_2 y_1} s_{x_2 x_1} - 2 f_{x_2 y_1} s_{y_2 x_1} - 2 f_{y_2 x_1} s_{x_2 y_1} \\
& + \left[ (f_{tx_1} s_{x_2 y_1} s_{y_2 u} + f_{ty_1} s_{y_2 x_1} s_{x_2 u} - 2 f_{tx_1} s_{y_2 y_1} s_{x_2 u} - 2 f_{ty_1} s_{x_2 x_1} s_{y_2 u}) \right. \\
& + (f_{tx_2} s_{y_2 x_1} s_{y_1 u} + f_{ty_2} s_{x_2 y_1} s_{x_1 u} - 2 f_{tx_2} s_{y_2 y_1} s_{x_1 u} - 2 f_{ty_2} s_{x_2 x_1} s_{y_1 u}) \\
& + (f_{y_2 x_1} s_{y_1 u} s_{x_2 t} + f_{x_2 y_1} s_{x_1 u} s_{y_2 t} - 2 f_{x_2 x_1} s_{y_1 u} s_{y_2 t} - 2 f_{y_2 y_1} s_{x_1 u} s_{x_2 t}) \left. \right] D^{tu} \\
& + \left[ f_u (s_{y_2 x_1} s_{x_2 v} s_{y_1 w} + s_{x_2 y_1} s_{y_2 v} s_{x_1 w} - 2 s_{y_2 y_1} s_{x_2 v} s_{x_1 w} - 2 s_{x_2 x_1} s_{y_2 v} s_{y_1 w}) \right. \\
& + (f_{tx_1} s_{x_2 u} s_{y_1 v} s_{y_2 w} + f_{ty_1} s_{y_2 u} s_{x_1 v} s_{x_2 w} + f_{tx_2} s_{y_2 v} s_{x_1 u} s_{y_1 w} + f_{ty_2} s_{x_2 v} s_{y_1 u} s_{x_1 w}) \left. \right] \Pi^{tv, uw} \\
& + f_u s_{y_2 v_1} s_{x_2 w_1} s_{y_1 v_2} s_{x_1 w_2} \Gamma^{t v_1 w_1, uv_2 w_2} \left. \right\} \\
& + \delta^{ij} \left\{ (4 s_{x_2 y_1} s_{y_2 x_1} - 2 s_{x_2 x_1} s_{y_2 y_1}) E_{act}^{(0)} \right. \\
& + 4 f_{x_2 y_1} s_{y_2 x_1} + 4 f_{y_2 x_1} s_{x_2 y_1} - 2 s_{x_2 x_1} f_{y_2 y_1} - 2 f_{x_2 x_1} s_{y_2 y_1} \\
& + \left[ (f_{tx_2} s_{y_2 y_1} s_{x_1 u} + f_{ty_2} s_{x_2 x_1} s_{y_1 u} - 2 f_{ty_2} s_{x_2 y_1} s_{x_1 u} - 2 f_{tx_2} s_{y_2 x_1} s_{y_1 u}) \right. \\
& + (f_{ty_1} s_{x_2 x_1} s_{y_2 u} + f_{tx_1} s_{y_2 y_1} s_{x_2 u} - 2 f_{ty_1} s_{y_2 x_1} s_{x_2 u} - 2 f_{tx_1} s_{x_2 y_1} s_{y_2 u}) \\
& + (f_{y_2 y_1} s_{x_1 u} s_{x_2 t} + f_{x_2 x_1} s_{y_1 u} s_{y_2 t} - 2 f_{x_2 y_1} s_{x_1 u} s_{y_2 t} - 2 f_{y_2 x_1} s_{y_1 u} s_{x_2 t}) \left. \right] D^{tu} \\
& \left[ + f_u (s_{y_2 y_1} s_{x_2 v} s_{x_1 w} + s_{x_2 x_1} s_{y_2 v} s_{y_1 w} - 2 s_{x_2 y_1} s_{y_2 v} s_{x_1 w} - 2 s_{y_2 x_1} s_{x_2 v} s_{y_1 w}) \right. \\
& + (f_{ty_1} s_{x_2 u} s_{x_1 v} s_{y_2 w} + f_{tx_1} s_{y_2 u} s_{y_1 v} s_{x_2 w} + f_{tx_2} s_{y_1 u} s_{y_2 v} s_{x_1 w} + f_{ty_2} s_{x_1 u} s_{x_2 v} s_{y_1 w}) \left. \right] \Pi^{tv, uw} \\
& + f_u s_{y_2 v_1} s_{x_2 w_1} s_{x_1 v_2} s_{y_1 w_2} \Gamma^{t v_1 w_1, uv_2 w_2} \left. \right\}
\end{aligned}$$

The expressions for matrix elements are complicated due to use of non-orthogonal orbitals.

In implementation, the following techniques are applied

1. Second Quantization Technique for non-orthogonal orbitals;
2. Tensor Analysis;
3. Automatic Formula/Code Generator.

## Example 1. The Spectroscopic Constants of H<sub>2</sub>

Method	$r_e$ ( $a_0$ )	$\omega_e$ (cm <sup>-1</sup> )	$D_e$ (kcal/mol)
FCI	1.405	4421	108.56
CASSCF <sup>a</sup>	1.427	4255	95.48
CASPT2N <sup>a</sup>	1.410	4407	105.40
VBSCF	1.429	4193	95.04
VBPT2	1.408	4376	106.30
icVBPT2	1.407	4379	106.09
VBCISD	1.405	4421	108.56

- 
- a. J. Phys. Chem., 1990, 94, 5483., where ANO(4s3p2d) basis set was used and orbitals  $1\sigma_g$  and  $1\sigma_u$  are taken as active orbitals.



## Example 2. The Spectroscopic Constants of N<sub>2</sub>

Method	$r_e$ ( $a_0$ )	$\omega_e$ (cm <sup>-1</sup> )	$D_e$ (kcal/mol)
FCI <sup>a</sup>	2.123	2341	201.75
VBSCF(17) <sup>b</sup>	2.109	2388	186.49
VBPT2(17) <sup>b</sup>	2.115	2373	194.21
icVBPT2(17) <sup>b</sup>	2.115	2367	195.25
VBSCF(175) <sup>c</sup>	2.114	2364	188.88
VBPT2(175) <sup>c</sup>	2.120	2344	197.72
icVBPT2(175) <sup>c</sup>	2.120	2345	197.65
VBCISD(17) <sup>b</sup>	2.121	2330	199.52
CASSCF <sup>d</sup>	2.119	2337	192.18
CASPT2N <sup>d</sup>	2.122	2342	198.82

a. J. Chem. Phys. 86, 5595 (1987).

b. 17 fundamental VB structures are used.

c. 175 fundamental VB structures are used, but the orbitals are optimized using 17 VB structures.

d. J. Chem. Phys. 96, 1218 (1992).

### Example 3. The Spectroscopic Constants of F<sub>2</sub>

Method	$r_e$ ( $a_0$ )	$\omega_e$ (cm <sup>-1</sup> )	$D_e$ (kcal/mol)
VBSCF	2.784	552.7	10.05
VBPT2	2.683	905.3	35.72
icVBPT2	2.708	900.0	33.67
VBCISD	2.689	886.8	35.35
L-BOVB	2.700	892.7	33.32
CASSCF <sup>a</sup>	2.755	803.1	19.62
CASPT2 <sup>a</sup>	2.691	899.1	35.19
MRCI <sup>a</sup>	2.680	889.5	32.81
Expt <sup>b</sup>	2.668	917	38.0

- a. The MOLPRO package is used for the CASSCF(2,2), CASPT2 and MRCI calculations. MOLPRO, version 2006.1, a package of ab initio programs, Werner, H.-J., etc.
- b. Chase, M. W. Jr. NIST-JANAF Thermochemical table, 4th ed.; J. Phys. Chem. Ref. Data, Monogr. 1998, 9, pp1-1951.

Example 4. Total energies (a.u.) of H<sub>2</sub>O as a function of the symmetric stretching from the equilibrium OH distance  $R_e$

Method	$R_e$	$1.5*R_e$	$2.0*R_e$	$2.5*R_e$	$3.0*R_e$	$8.0*R_e$
FCI <sup>a</sup>	-76.2419	-76.0723	-75.9517	-75.9180	-75.9119	-75.9103
CASSCF(8,6) <sup>a</sup>	0.1640	0.1500	0.1336	0.1263	0.1247	
CASPT2 <sup>a</sup>	0.0128	0.0108	0.0081	0.0080	0.0083	
VBSCF <sup>b</sup>	0.1857	0.1626	0.1380	0.1272	0.1250	0.1242
VBPT2(occ) <sup>c</sup>	0.0353	0.0374	0.0381	0.0394	0.0397	0.0397
VBPT2(all) <sup>d</sup>	0.0274	0.0266	0.0268	0.0283	0.0287	0.0287
icVBPT2	0.0201	0.0121	0.0093	0.0097	0.0099	0.0099

- a. Olsen, J., Jorgensen, P., Koch, H., Balkova, A. Bartlett, R. J. J. Chem. Phys. 1996, 104, 8007.
- b. 20 structures are used in VB calculations.
- c. Only excitations from occupied orbitals to virtual ones are included.
- d. All excitations are included, involving excitations from inactive orbitals to active ones.

# Summary 2

---

- icVBPT2 provides a cheaper ab initio VB tool to cover dynamical correlation.
- Test calculation shows that both VBPT2 and icVBPT2 results are in good agreement with CASPT2 ones.
- Due to the use of delocalized virtual orbitals, excited structures are not classically physical any more.

### **3. DFVB: A VB Method with Density Functional Energy Correction**

**DFT: Single reference based method**

**VBSCF: lack of dynamic correlation**

**VB + DFT**

## From KS-DFT to DFVB

### KS-DFT

$$E^{\text{KS-DFT}}(\rho) = T_S[\rho] + J[\rho] + E_{\text{Ne}}[\rho] + E_{\text{XC}}[\rho]$$

$$E_{\text{XC}} = (T[\rho] - T_S[\rho]) + (E_{\text{ee}}[\rho] - J[\rho])$$

### DFVB

$$E^{\text{DFVB}}[\rho^{\text{VB}}] \approx E^{\text{VBSCF}} + E_{\text{C}}[\rho^{\text{VB}}]$$

Approximate

Problem:

How to handle non-orthogonal VB orbitals in KS-DFT?



Problem:

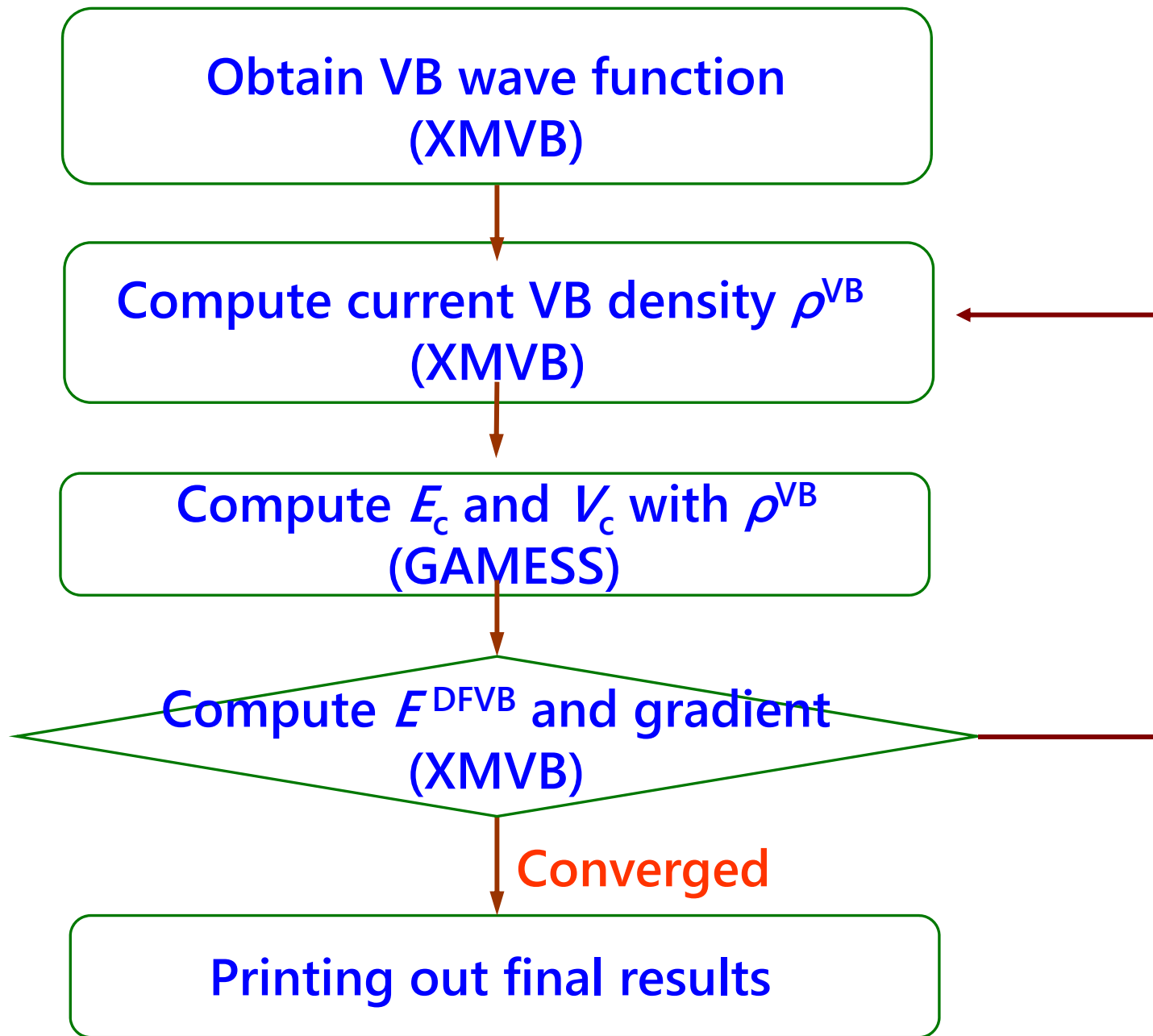
How to handle non-orthogonal VB orbitals in KS-DFT?

Non-orthogonal VB orbitals

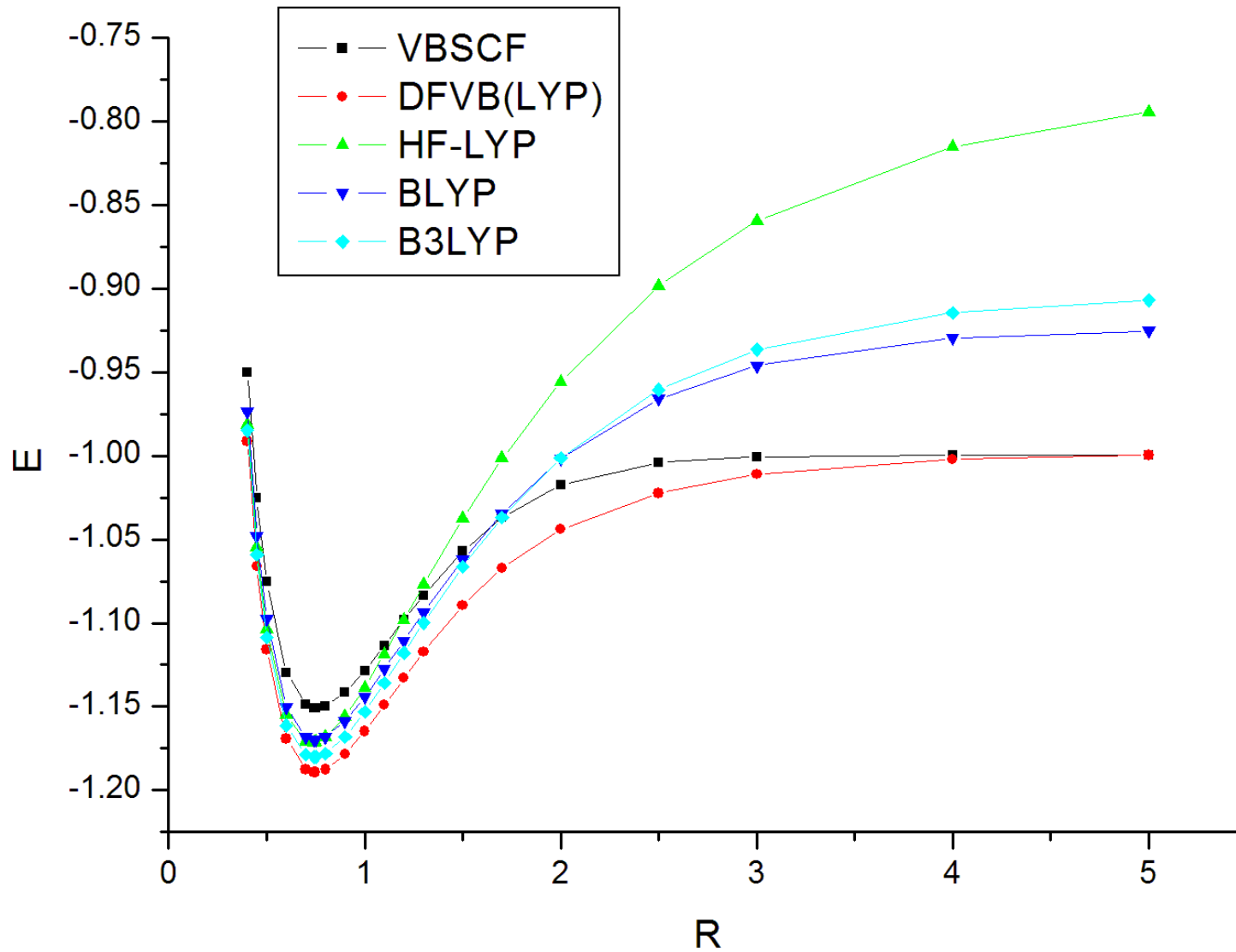
```
graph TD; A[Non-orthogonal VB orbitals] --> B[VB density]; B --> C[Orthogonal Natural Orbitals]
```

VB density

Orthogonal Natural Orbitals



# PEC of H<sub>2</sub> with Different Methods with cc-pVTZ



# $R_0$ of diatomic molecules

Mol	Exp	B3LYP	BLYP	PW91	VBSCF	DFVB	
						LYP	PW
H <sub>2</sub>	0.741	0.002	0.006	0.008	0.015	0.008	0.010
HF	0.917	0.005	0.016	0.012	-0.003	-0.010	-0.013
HCl	1.275	0.006	0.015	0.012	0.011	-0.001	-0.003
Li <sub>2</sub>	2.673	0.028	0.038	0.054	0.272	0.157	0.197
N <sub>2</sub>	1.098	-0.007	0.005	0.003	-0.003	-0.012	-0.014
O <sub>2</sub>	1.208	-0.004	0.023	0.011	0.035	0.019	0.009
F <sub>2</sub>	1.412	-0.016	0.019	0.000	0.065	0.016	-0.011
Cl <sub>2</sub>	1.988	0.023	0.052	0.019	0.050	0.009	-0.016
ClF	1.628	0.012	0.044	0.022	0.031	-0.005	-0.023
NF	1.317	0.001	0.026	0.010	0.043	0.014	0.003
<b>MAE</b>		<b>0.010</b>	<b>0.024</b>	<b>0.015</b>	<b>0.053</b>	<b>0.025</b>	<b>0.030</b>

## Frequency ( $\omega_e$ ) of diatomic molecules (Errors to experimental values)

Mol	Exp	B3LYP	BLYP	PW91	VBSCF	DFVB	
						LYP	PW
H <sub>2</sub>	4401	18	-55	-69	-179	-92	-92
HF	4138	-13	-172	-116	15	122	177
HCl	2991	-20	-113	-69	-13	-4	63
Li <sub>2</sub>	351	-7	-18	-18	-88	-53	-59
N <sub>2</sub>	2359	91	-23	-1	14	73	95
O <sub>2</sub>	1580	56	-91	-31	-76	1	60
F <sub>2</sub>	917	139	51	90	-277	-135	-56
Cl <sub>2</sub>	560	-18	-55	-20	57	-15	15
ClF	786	3	-56	-19	-95	-35	6
NF	1141	16	-59	-8	-86	10	52
<b>MAE</b>		<b>38</b>	<b>69</b>	<b>44</b>	<b>90</b>	<b>54</b>	<b>68</b>

## BDE of diatomic molecules (Errors to experimental values)

Mol	$D_e$	B3LYP	BLYP	PW91	VBSCF	DFVB	
						LYP	PW
H <sub>2</sub>	109.5	3.7	0.0	-4.3	-14.6	9.4	6.0
HF	141.3	-2.6	-1.2	0.5	-30.4	-4.9	-5.0
HCl	106.4	-0.3	-1.1	1.0	-17.0	5.2	6.1
Li <sub>2</sub>	24.4	-3.7	-3.9	-4.5	-14.6	1.7	-1.2
N <sub>2</sub>	228.5	3.1	13.8	16.1	-67.3	-5.8	-10.1
O <sub>2</sub>	120.3	4.3	17.7	25.4	-45.3	-11.2	-9.3
F <sub>2</sub>	38.2	2.3	14.4	18.7	-25.7	-6.9	-7.4
Cl <sub>2</sub>	58.0	-0.2	2.2	10.0	-21.0	-2.5	2.3
ClF	60.2	2.4	8.9	14.5	-29.0	-8.6	-6.9
NF	70.4	14.0	23.6	27.5	-32.0	-2.5	-3.3
Au <sub>2</sub>	53.0	-10.6	-7.2	-3.0	-32.9	-18.6	-18.2
<b>MAE</b>		<b>4.3</b>	<b>8.5</b>	<b>11.4</b>	<b>30.0</b>	<b>7.0</b>	<b>6.9</b>

# S-T energy gaps of diradical systems (Errors to experimental values)

	Expt	Refs	CASPT2	B3LYP	BLYP	VBSCF	DFVB	
							LYP	PW
<b>C</b>	29.1 <sup>[1]</sup>	1.3 <sup>[5]</sup>	0.9	11.3	10.1	5.4	-4.5	-1.3
<b>O</b>	45.4 <sup>[1]</sup>	-1.6 <sup>[1]</sup>	0.9	17.0	15.5	4.9	-5.1	-1.7
<b>O<sub>2</sub></b>	22.7 <sup>[2]</sup>	1.0 <sup>[6]</sup> /0.7 <sup>[7]</sup>	0.9	13.1	12.4	-2.4	-5.7	-5.1
<b>NF</b>	32.7 <sup>[3]</sup>	2.6 <sup>[3]</sup>	0.1	10.8	8.9	5.9	-3.0	-0.8
<b>CH<sub>2</sub></b>	32.9 <sup>[1]</sup>	1.6 <sup>[5]</sup> /-1.3 <sup>[1]</sup>	-6.0	-25.5	-25.6	6.3	-7.0	-1.6
<b>TMM</b>	18.1 <sup>[4]</sup>	-0.3 <sup>[3]</sup>	2.0	25.7	16.6	3.4	-3.4	-0.4
<b>MAE</b>			<b>1.8</b>	<b>17.2</b>	<b>14.9</b>	<b>4.7</b>	<b>4.8</b>	<b>1.8</b>

[1] D.H. Ess, E.R. Johnson, X. Hu, et al., *J Phys Chem A*(2011) **115**(1), 76

[2] K.P. Huber and G. Herzberg, *Constants of Diatomic Molecules*, in *Molecular Spectra and Molecular Structure*, V.N. Reinhold, Editor. 1979, New York

[3] G.S. Harbison, *J Am Chem Soc*(2001) **124**(3), 366

[4] XZ Li and J. Paldus, *J Chem Phys* (2008) **129**(17), 174101

[5] L.V. Slipchenko and A.I. Krylov, *J Chem Phys* (2002) **117**(10), 4694

[6] J. Pittner, P. Čársky and I. Hubač, *Int. J. Quantum Chem*(2002) **90**(3), 1031

[7] P. Su, L. Song, W. Wu, et al., *J. Comput. Chem*(2007) **28**(1), 185.

# Summary 3

---

- DFVB method incorporates VB method to density functional theory.
- Test calculations show that DFVB improves VBSCF results, and overcomes the difficulty of DFT for diradical systems.
- Double-counting problem still exists and need to be corrected.



## **4. Valence Bond Polarizable Continuum Model (VBPCM) Method**

# VBPCM Method

---

Solvation effects play a very important role in molecular energy, structures, and properties.

PCM is an efficient and economical tool for describing solvation problems.

The solvent is usually represented as a homogeneous medium that is characterized by a single dielectric constant.

The QM packages of PCM are widely applied to study solvent effect in ab initio level.

# The PCM Model

---

$$H^0 \Psi^0 = E^0 \Psi^0$$

In Vacuum:

$$(H^0 + V_R) \Psi = E \Psi$$

In Solution:

$$V_R = V_{el} + V_{dis} + V_{rep}$$

Minimizing

$$G = \langle \Psi | H^0 | \Psi \rangle + \langle \Psi | V_R'' | \Psi \rangle + \frac{1}{2} \langle \Psi | V_R'(\Psi) | \Psi \rangle + V_{NN} + G_{nel}$$

$$\mathbf{F}^S \mathbf{C} = \mathbf{S} \mathbf{C} \varepsilon$$

As one of QM method, essentially, PCM can be applied to VB method by the same strategy.

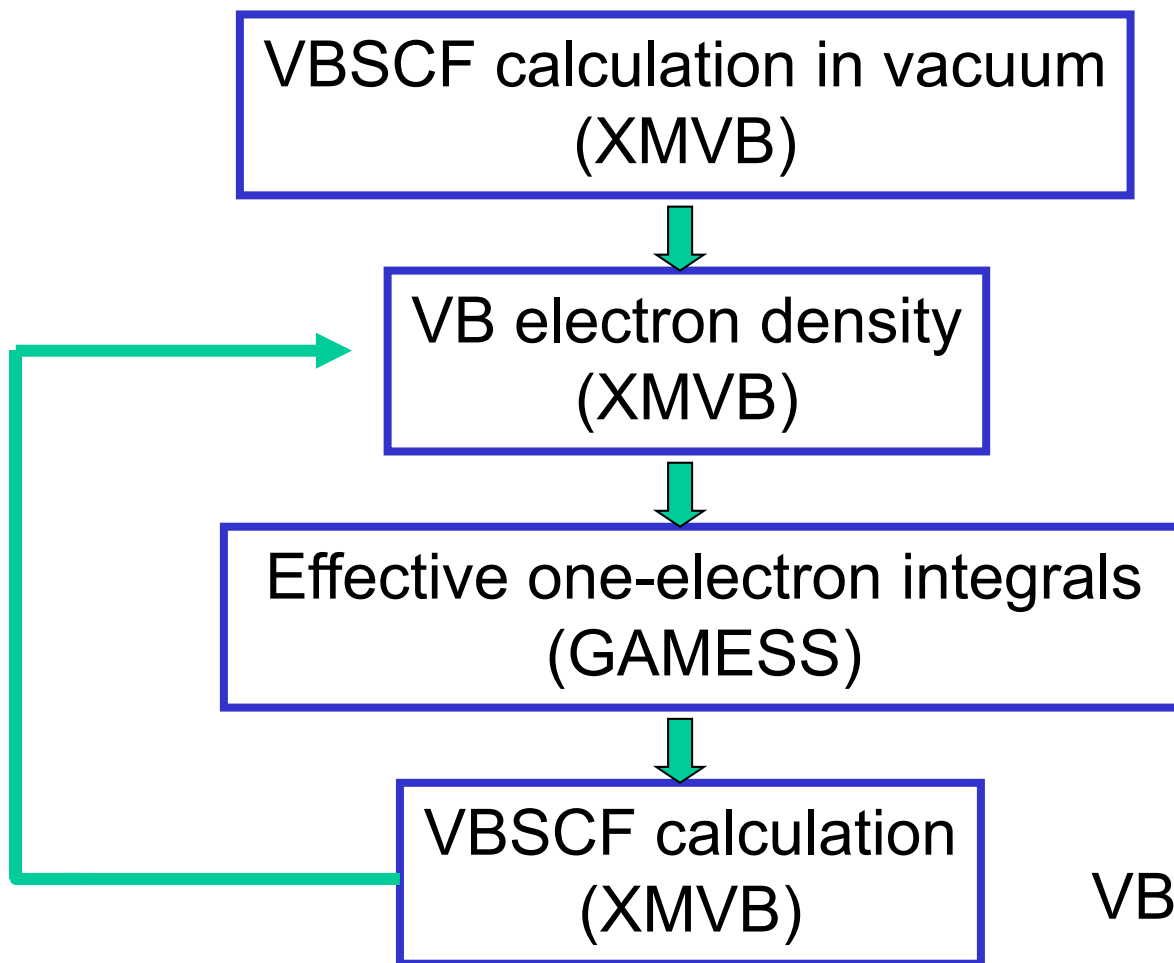
The difference between VB and MO-based methods is the form of wave function.

VBPCM method incorporates PCM into VB method by using **VB density** for the solvent-solute interaction potential, instead of MO-based density.

J. Phys. Chem. A, 108, 6017-6024 (2004).

# VBPCM Calculation

---



VBPCM//VBSCF

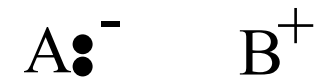
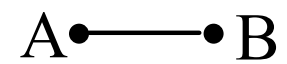
$$E = \langle \Psi | H^0 + \frac{1}{2} V_R | \Psi \rangle$$

# Dissociation of LiF

---

VBPCM//VBSCF

6-31G\* and 6-31+G\*

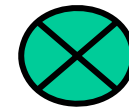


1

2

3

A=Li, B=F



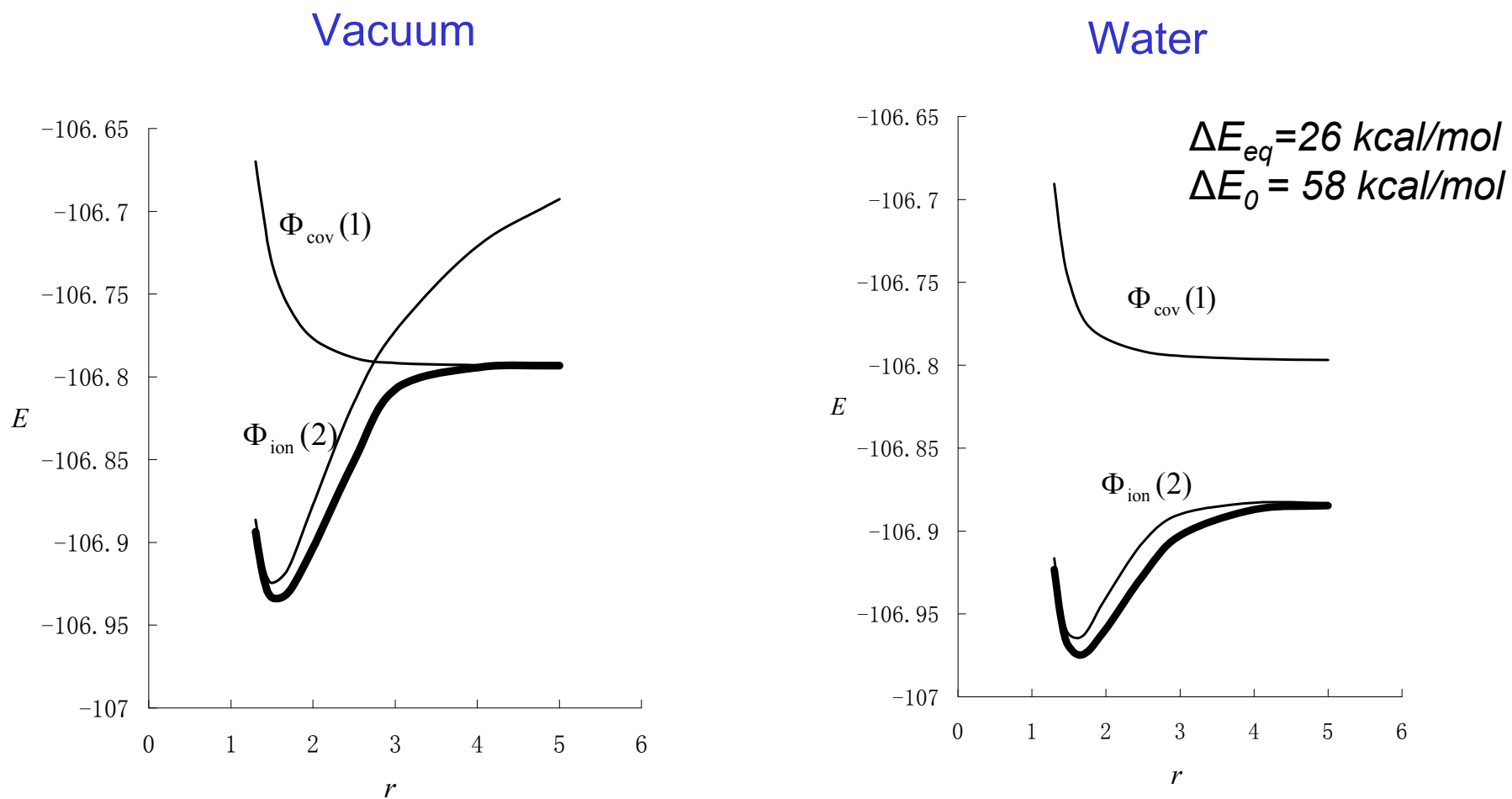


Figure 1a. VBSCF/6-31G\* dissociation energy profiles of LiF in vacuum. Adiabatic potentials are shown in bold curves. 1b. VBPCM//VBSCF/6-31G\* dissociation energy profiles of LiF in H<sub>2</sub>O.

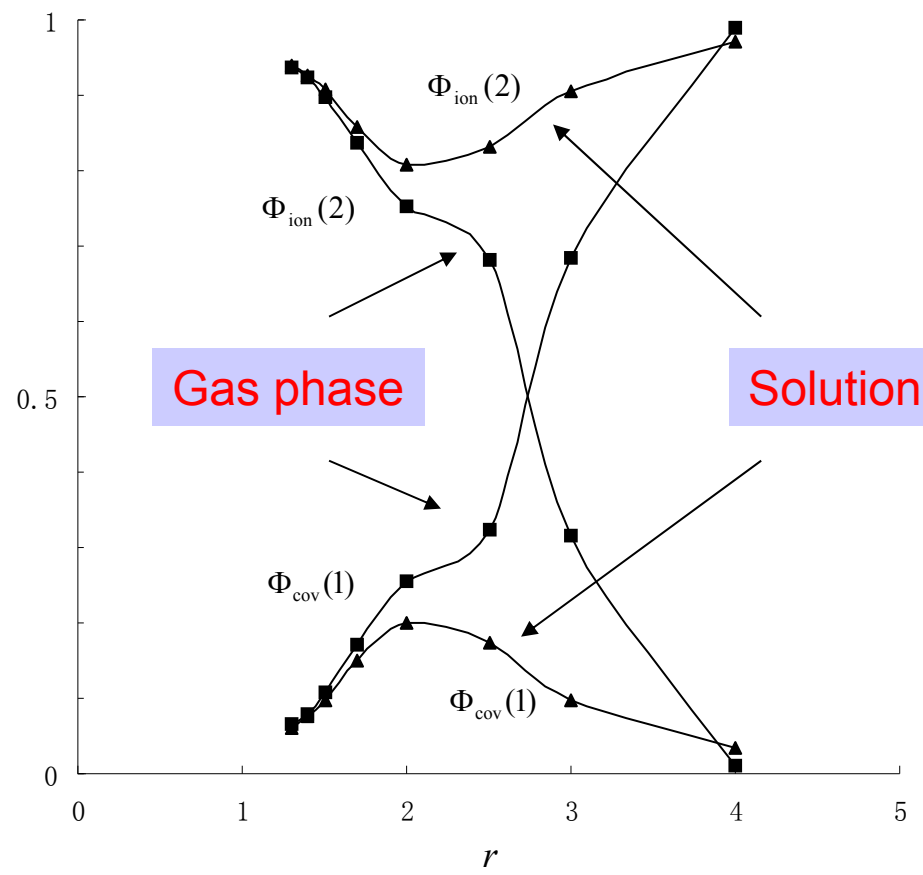
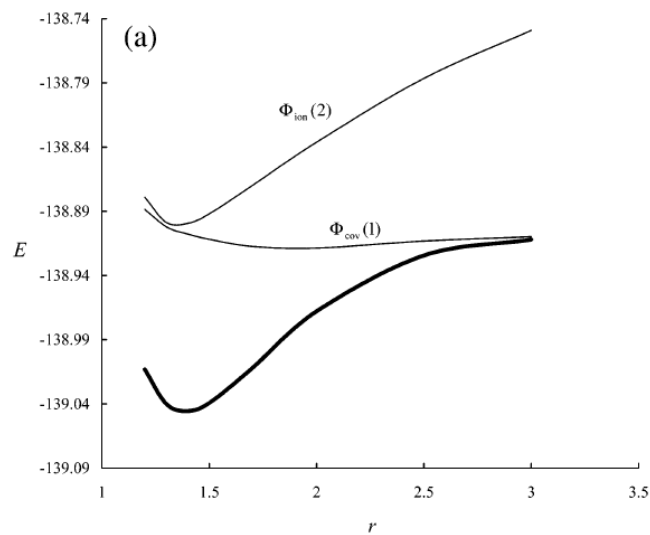


Figure 3. Weights of VB structures of LiF. VBSCF/6-31G\* weights are annotated with bold squares and VBPCM//VBSCF/6-31G\* weights with triangles.

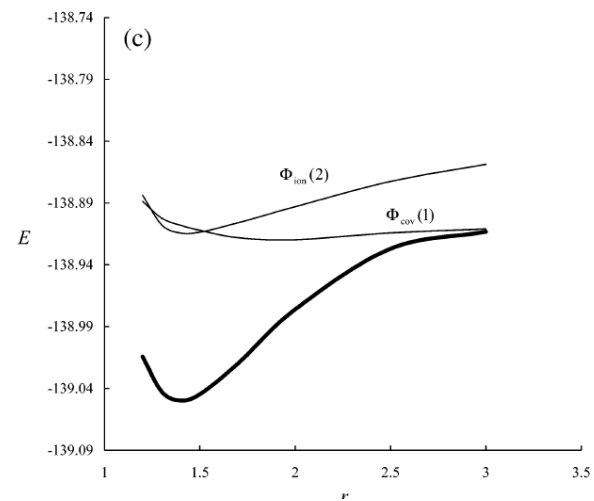
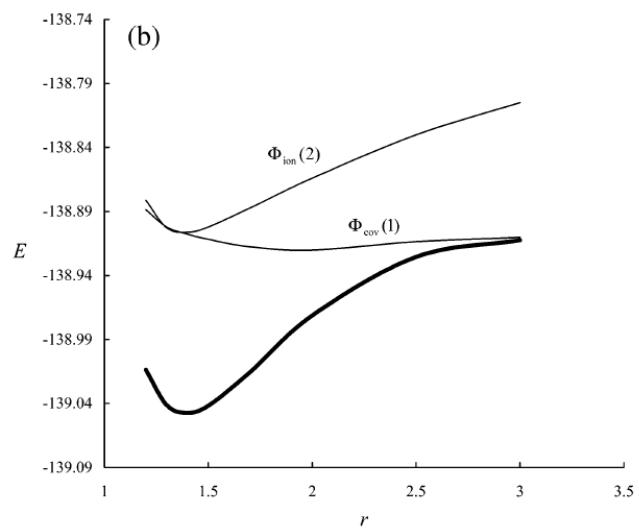


# Dissociation of $\text{CH}_3\text{F}$ (polar-covalent bond)



$\text{CCl}_4$

$\text{H}_2\text{O}$



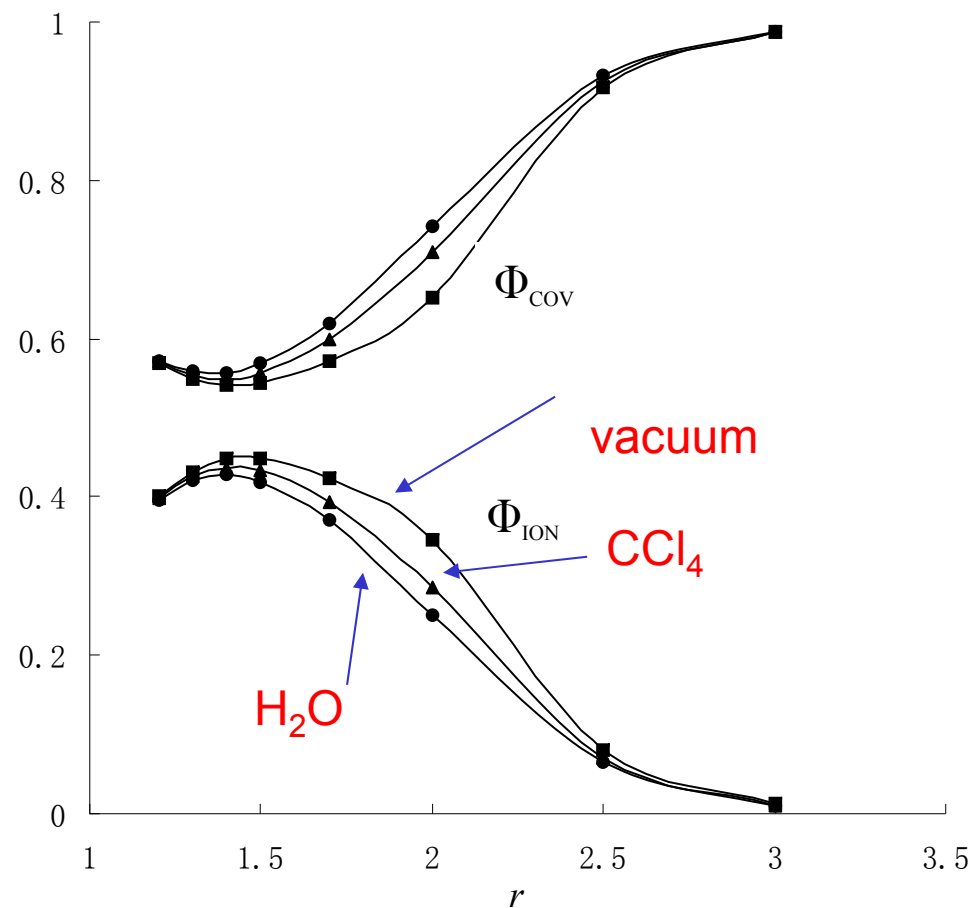


Figure 5. VBSCF/6-31G\* and VBPCM//VBSCF/6-31G\* weights of VB structures of  $\text{CH}_3\text{F}$  along the C-F dissociation coordinate. The curves in vacuum,  $\text{CCl}_4$ ,  $\text{H}_2\text{O}$  are annotated with bold squares, triangles and circles respectively.

# Summary 4

---

- VBPCM method enables us to compute the energy profile of the full state as well as of individual VB structures. Thus, it provides qualitative insights into the solvent effects of chemical problems.
- One-electron density plays a role as a bridge between the VB and the PCM methods: More general VBPCM methods can be presented by combining different VB methods and solvent models.

# More VB Methods for Solvation Effects

---

## Valence Bond Solvation Model (VBSM)

Su; Wu; Kelly; Cramer; Truhlar, *J. Phys. Chem. A*, 2008, 112, 12761.

## Valence Bond Effect Fragment Potential (VBEFP)

Ying; Chang; Su; Wu, *J. Phys. Chem., A*, 2012, 116, 1846.

## VB/EFP/PCM

Su; Ying; Wu, in preparation

## Combined VB and MM method (VB/MM)

Shurki; Cromn, *J. Phem. Chem. B*, 2005, 109, 23638;

Sharir-Ivry; Crown; Wu; Shurki, *J. Phys. Chem. A*, 2008, 112, 2489

## Explicit Solvation VB Method

Braida; Hiberty, *Int. J. Quant. Chem.* 2010, 110, 571.

# Acknowledgements

---

## Xiamen Group:

Zhenhua Chen, Peifeng Su, Fuming Ying, Jinshuai Song,  
Xun Chen, Lingchun Song, Qianer Zhang

## Collaborators:

Prof. S. Shaik (Hebrew University, Jerusalem)

Prof. P. C. Hiberty (Université de Paris-Sud, Orsay)

Prof. Y. Mo (Western Michigan University)

\$ Natural Science Foundation of China



*Thank you!*



# Total Time Consuming by Various VB Methods

---

	<b>VBSCF</b>	<b>VBCISD</b>	<b>VBPT2</b>	<b>DFVB(LYP)</b>	<b>DFVB(PW)</b>
<b>H<sub>2</sub></b>	0.04	0.05	0.05	4.96	5.76
<b>HF</b>	0.10	14.18	0.10	80.76	126.49
<b>F<sub>2</sub></b>	0.52	5217.80	0.37	354.28	465.88
<b>O<sub>2</sub></b>	0.81	35506.17	0.67	390.40	441.10
<b>N<sub>2</sub>(17str)</b>	1.02 <sup>a</sup>	147163.48	0.74	917.64	1019.32
<b>N<sub>2</sub>(11str)</b>	0.84 <sup>a</sup>	101598.66	0.73	602.10	819.08

unit: second

<sup>a</sup> with GUESS=MO

lscf=5 for VBSCF, lscf=1 for DFVB

# Computational costs of VBSCF, VBCISD, VBPT2, icVBPT2 and icVBCISD

The computational cost of integral transformation:

VBSCF Energy	VBSCF Gradient	VBSCF Hessian	VBCISD	VBPT2	icVBPT2	icVBCISD
$(N_{act})^*(N_{bas})^4$	$(N_{act})^*(N_{bas})^4$	$(N_{occ})^*(N_{bas})^4$	$(N_{bas})^5$	$(N_{occ})^*(N_{bas})^4$	$(N_{occ})^*(N_{bas})^4$	$(N_{bas})^5$

$N_{act}$ ,  $N_{occ}$ ,  $N_{bas}$  are the number of active, occupied orbitals and basis functions respectively.



VBCISD	VBPT2	icVBPT2 <sup>c</sup>	icVBCISD <sup>d</sup>
$(N_{CISD}^{VB})^2 (N_{ele})^4$	$(N_{PT2}^{VB})^2 (N_{ele})^2$	$(N_{str}^{\alpha})^2 (N_{act}^{\alpha})^8$ $+ (N_{str}^{\beta})^2 (N_{act}^{\beta})^8$ $+ (N_{det})^2 (N_{act}^{\beta})^4$	$(N_{str}^{\alpha})^2 (N_{act}^{\alpha})^{10}$ $+ (N_{str}^{\beta})^2 (N_{act}^{\beta})^{10}$ $+ (N_{det})^2 (N_{act}^{\beta})^4$

- a. string-based implementation
- b. The computational cost of RDMs up to 3<sup>th</sup>-order.
- c. The computational cost of RDMs up to 4<sup>th</sup>-order.
- d. The computational cost of RDMs up to 5<sup>th</sup>-order.