RPA and CC theory & some recent results including range separation Gustavo E. Scuseria



RPA Workshop

Paris

26-29 Jan 10

Outline

- The connection between RPA and coupled cluster (CC) theory
- <u>Range-separation</u>: a powerful tool for mixing wavefunction and density functional theories
- Some recent results for RPA ground-state correlations
- A model for strong correlations based on HFB

Coupled Cluster (CC) Theory

- Arguably the most successful quantum chemistry theory
- Borrowed from nuclear physics (Coester-Kümmel) in the 60s by Cizek & Paldus
- Single-reference version can treat metals but not stronglycorrelated systems (requires multi-reference [MR])
- MRCC theory has "issues" and has never taken off...
- Size extensivity (CC) is more important than variational energy bound (CI). This was settled in mid to late 80s and changed the QC paradigm
- Of course, 2nd and larger revolution in QC was DFT in early 90s (hybrid functionals)

Coupled Cluster (CC) Theory

 $|\Psi \rangle = e^{T} |\Phi \rangle$ Exponential (as opposed to Linear) T ansatz T = T₁ + T₂ + T₃ + ... Electronic structure: T₂ dominates!

In simplest form, CCD: $T_2 = \Sigma t_{ik}^{ab} a_a^{\dagger} a_b^{\dagger} a_i a_k$

CCD equations: $E = \langle \Phi | e^{-T} H e^{T} | \Phi \rangle = \text{linear eqn in } T_2$ $O = \langle \Phi | T_2^{+} H e^{T} | \Phi \rangle = \text{quadratic}$ algebraic eqn in T_2

THE JOURNAL OF CHEMICAL PHYSICS 129, 231101 (2008) The ground state correlation energy of the random phase approximation from a ring coupled cluster doubles approach

Gustavo E. Scuseria, Thomas M. Henderson, and Danny C. Sorensen

RPA problem for excitations is:

RPA ground-state correlation energy: difference in ZPVE between two harmonic oscillators

(1) is mathematically equivalent to a ring CCD equation

Energies are mathematically identical with **CC** amplitudes

$$\begin{pmatrix} \mathbf{A} & \mathbf{B} \\ -\mathbf{B} & -\mathbf{A} \end{pmatrix} \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} = \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} \boldsymbol{\omega}, \qquad (1)$$

$$E_c^{\text{RPA}} = \frac{1}{2} \text{Tr}(\boldsymbol{\omega} - \mathbf{A}).$$
 (2)

$$\mathbf{B} + \mathbf{AT} + \mathbf{TA} + \mathbf{TBT} = 0, \qquad (3)$$

$$E_c^{\text{rCCD}} = \frac{1}{2} \text{Tr}(\mathbf{BT}) = \frac{1}{2} \text{Tr}(\boldsymbol{\omega} - \mathbf{A}).$$
(4)

 $\mathbf{T} = \mathbf{Y}\mathbf{X}^{-1},\tag{13a}$

RPA for ground-state correlations

- Q-Chem perspective from the late 70s & early 80s:
 - Plagued by triplet instability problem
 - When no instabilities are present (~Re),
 E_c is off by a factor 2
 - Ouch! Not good
- Green's function & DFT perspective:
 - Forget about **exchange**, just keep only the "direct" (Coulomb) part of the interaction!
- Perdew & Langreth (80s): long-range RPA correlation energy is ok, short-range is bad

A hierarchy of DFT approximations

The quest for accurate exchange-correlation potentials

- LSDA : functional of electron density
- **GGA**: adds gradient of electron density
- meta-GGA: adds kinetic energy density



- hybrids: add <u>nonlocal</u> HF-type exchange (Generalized Kohn-Sham scheme)
- This talk: range-separated hybrids

 + long-range RPA correlation
 + strong static correlation via "HFB"

Hybrid Functional example

- **PBE** : Perdew, Burke, Ernzerhof, **PRL** (1996)
 - GGA (depends on e-density + gradient e-density)
 - No empirical parameters
- **PBEh** : add HF-type exchange

 $E(PBEh) = a E_{x}(HF) + (1-a) E_{x}(PBE) + E_{c}(PBE)$

a = 0.25 based on good theory Uses nonlocal HF-type potential (Generalized Kohn-Sham scheme)

Perdew, Ernzerhof & Burke (1997); Ernzerhof & Scuseria (1999); Adamo & Barone (1999)

Range-separation

$$\frac{1}{r_{12}} = \frac{1 - erf(\omega r_{12})}{r_{12}} + \frac{erf(\omega r_{12})}{r_{12}} \qquad \omega \ge 0$$

short-range long-range

(a) determines the splitting between short- and long-range

- Andreas Savin (1985) proposed to use
 - DFT for short-range
 - Wavefunction theory for long-range
- Keep the **best** of both worlds: $SR \rightarrow local$ $LR \rightarrow nonlocal$
- Rigorous mathematical support (ask Andreas!)

LC-wPBE : a long-range hybrid

Vydrov & Scuseria, JCP 125, 234109 (2006)

 $E(LC-\omega PBE) = E_{x}(\omega PBE, SR) + E_{x}(\omega HF, LR) + E_{c}(PBE)$

100% PBE exchange in SR

100% HF exchange in LR

Only one parameter: w

w=0.40 (fitted)

No range-separation for correlation (we will add RPA ground state correlation later)

LC-wPBE: Mean Absolute Errors (Molecules)

| Functional | Atomic E | G2/148 | HTBH38 | |
|------------|----------|------------|------------|--|
| | (mH) | (kcal/mol) | (kcal/mol) | |
| LDA | 68.0 | 83.3 | 18.9 | |
| PBE | 7.6 | 16.9 | 9.7 | |
| PBEh | 6.3 | 4.9 | 4.6 | |
| LC-wPBE | 3.4 | 3.7 | 1.3 | |

Excellent HTBH38 activation barriers and G2 heats of formation

Vydrov & Scuseria, JCP 125, 234109 (2006)

One missing ingredient:

van der Waals

important for biological applications

RPA as ring-CCD for correlation

- Using the CC connection we can range-separate DFT+ RPA like
 DFT + WF (wavefunction method)
- Ansatz: $E_{xc} = E_{xc}(SR-DFT) + E_{x}(LR-HF) + E_{c}(LR-RPA)$
- Alternative approach: Toulouse et al. (PRL 2009)

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$$A_{ia,jb} = (\epsilon_a - \epsilon_i)\delta_{ij}\delta_{ab} + \langle ib|aj \rangle,$$
(7a)
$$B_{ia,jb} = \langle ij|ab \rangle.$$
(7b)

"Single-bar" integrals define *direct* RPA (dRPA) No instability issues!

Two-electron integrals are positive definite

$$\langle ib|aj\rangle = \langle ij|ab\rangle = u^A_{ia}u^A_{jb}, \qquad (18a)$$

In *direct* ring-CCD T are negative definite

$$-t^{ab}_{ij} = \theta^A_{ia} \theta^A_{jb}, \tag{18b}$$

RPA correlation can be evaluated very efficiently:

$$t_{ij}^{ab} = \frac{1}{\Delta \epsilon_{ij}^{ab}} \left(u_{ia}^A u_{jb}^A - u_{ia}^A u_{kc}^A \theta_{kc}^B \theta_{jb}^B - \theta_{ia}^A \theta_{kc}^A u_{kc}^B u_{jb}^B + \theta_{ia}^A \theta_{kc}^A u_{kc}^B u_{ld}^B \theta_{ld}^C \theta_{jb}^C \right)$$
(19)

Why not MP2 ?

- MP2 = 2nd order perturbation theory
 = 1st term of RPA equations
- Problems:
 - MP2 cannot treat metals (diverges)
 - **RPA** remains finite for metals
 - MP2 is OK for van der waals (but not as good as RPA)
- RPA resolves the conundrum* of molecules versus solids
 - Molecules: LR-HFx is good; LR-RPA adds vdw
 - Small band gap solids: LR-HFx is pathological LR-RPA fixes it!

(HSE neglects both LR-HF and LR-RPA)

Some RPA results

 Adding range-separated RPA correlation to LC-wPBE... not ready yet...

 No range-separated PBE correlation (coming soon)

Instead, use LC-wLDA

THE JOURNAL OF CHEMICAL PHYSICS 130, 081105 2009 Long-range-corrected hybrids including random phase approximation correlation

Benjamin G. Janesko, Thomas M. Henderson, and Gustavo E. Scuseria

Approximation tested is based on LSDA

$$E_{\mathbf{x}\mathbf{c}} = E_{\mathbf{x}\mathbf{c}}^{\mathrm{SR-LSDA}} + E_{\mathbf{x}}^{\mathrm{LR-HF}} + c_{\mathrm{RPA}}E_{\mathbf{c}}^{\mathrm{LR-RPA}}.$$
 (6)

LC-wLDA + dRPA

SR-LDA correlation from Gori-Giorgi et al.

dRPA orbitals from LC-wLDA

Two parameters: w and C_{RPA}

Two models tested: 1. For $C_{\text{RPA}} = 1.00 \rightarrow \omega$ (opt) = 0.70 2. Both optimized: $\omega = 1.20$, $C_{\text{RPA}} = 1.50$

LC-wLDA + dRPA model chemistry

Mean Absolute Errors (kcal/mol)

| Method | w(opt) | AE6 | G2 | BH6 | HT | NHT |
|------------------|--------|-----|-----|-----|-----|-----|
| LC-wPBE | 0.4 | 5.5 | 4.2 | 1.2 | 1.3 | 2.0 |
| LC-wLDA | 0.5 | 5.8 | 7.0 | 2.3 | 3.0 | 4.4 |
| LC-wLDA + dRPA | 0.7 | 5.6 | 6.2 | 1.8 | 2.3 | 3.5 |
| LC-wLDA + C dRPA | 1.2 | 4.0 | 4.4 | 1.2 | 1.6 | 3.5 |

C=1.5

6-311+G(2d,2p)

Janesko, Henderson, Scuseria, JCP 130, 081105 (2009)







LC-wLDA + dRPA "Biological" noncovalent interactions

Mean Absolute Errors (kcal/mol)

| Method | HB6 | CT7 | DI6 | WI9 | MARE |
|------------------|------|------|------|------|------|
| MP2 | 0.44 | 0.29 | 0.06 | 0.05 | 13 |
| CCSD(T) | 0.38 | 0.47 | 0.29 | 0.05 | 12 |
| LC-wLDA + dRPA | 0.33 | 1.41 | 0.72 | 0.24 | 34 |
| LC-wLDA + C dRPA | 0.49 | 0.35 | 0.06 | 0.02 | 8 |

w=1.2 C=1.5

aug-cc-pVTZ basis CCSD(T)/CBS, MARE=0 HB: Hydrogen Bonds

CT: Charge Transfer Complexes

- DI: Dipole-Dipole Complexes
- **WI: Weak Interactions**

Janesko, Henderson, Scuseria, JCP (2009)

Why C=1.5 ? (for ω=1.2)

Correction for:

- 1. Basis set effects?
- 1. Beyond-RPA correlation?
- 1. Orbitals?
- 1. Large w?

Answer:

3 and **4**

The role of the reference state in long-range RPA, B. G. Janesko and G. E. Scuseria, *J. Chem. Phys.* **131**, 154106 (2009)

A short story about RPA

Scuseria, Henderson, and Sorensen, JCP **129**, 231101 (2008)

$$\langle ij|ab \rangle + (\varepsilon_c - \varepsilon_k) \delta_{ac} \delta_{ik} t_{kj}^{cb} + \langle ic|ak \rangle t_{kj}^{cb}$$

"direct ring-CCD"

$$+ t_{ik}^{ac} (\varepsilon_c - \varepsilon_k) \delta_{bc} \delta_{jk} + t_{kj}^{cb} \langle ic|ak \rangle$$
$$+ t_{ik}^{ac} \langle kl|cd \rangle t_{lj}^{db} = 0$$

B + AT + TA + TBT = 0 "Riccati Equation"

A short story about dRPA

direct RPA correlation energy:

$$\mathsf{E}_{\mathsf{corr}} = \frac{1}{2} \sum_{ijab} \langle ij | ab \rangle t_{ij}^{ab} \equiv \frac{1}{2} \mathsf{Tr}(\mathbf{BT})$$

$$\begin{split} \Phi_{dRPA} &= \exp(\hat{T}_2) |\Psi_0^{\mathsf{HF}(\mathsf{PBE})} \rangle \\ &= (1 + \hat{T}_2 + \frac{1}{2} \hat{T}_2^2 + \dots) |\Psi_0^{\mathsf{HF}(\mathsf{PBE})} \rangle \end{split}$$

Definition of SOSEX (David Freeman, Phys Rev 1977)

$$E_{\text{corr}} = \frac{1}{2} \sum_{ijab} (\langle ij|ab \rangle - \langle ij|ba \rangle) t_{ij}^{ab}$$
$$= \frac{1}{2} \operatorname{Tr}(\mathbf{B}_{\text{as}}\mathbf{T}), \quad \text{with} \quad \mathbf{B}_{\text{as}} = \langle ij||ab \rangle$$

Exchange term corrects for all self-interaction terms <ii|aa>, i.e. i = j and a = b.





Last few bullets on RPA

- Long-range HF + dRPA looks very promising combined with short-range semilocal xc DFT
- Molecules: RPA is better than MP2 but not as accurate as CCSD(T)
- Solids: neither MP2 nor (T) are useful for small band gap systems
- Cost of dRPA correlation with ring-CCD is between O(N⁵)→O(N³) and asymptotically O(N)
- Better (self-consistent) orbitals are crucial for ground-state RPA correlation

One important additional ingredient

Strong/static correlation

Static/Strong Correlation

New Method <u>Desideratum</u>:

- Should preserve space and spin symmetries (avoid spatial symmetry breaking and spin contamination)
- 2. Should have low-computational cost (meanfield instead of CASSCF exponential blowup)
- 3. Should cleanly separate static & dynamic correlation

How do we accomplish this?

- **Break** a different (new!) symmetry in mean-field: electron number conservation
- Introduce electron number fluctuations

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- The theory that accomplishes this is **BCS (HFB**) but for Coulombic Fermionic systems, HFB always reduces to HF (Lieb 1994)
 - We introduce an attractive pairing interaction (-1/r₁₂) in a "<u>hybrid</u>" HFB scheme with HF (2/r₁₂)
 - Is mean-field (roughly same computational cost as KS or HF)
- Is "exact" at dissociation (Σ ROHF atomic energies)

Our CPMFT model

- Yields a "*definition*" for static correlation from the **2pdm** ansatz
- Reduces to **HF** in the **absence** of strong correlation
- Yields correct symmetry-adapted orbitals
- BCS-type (HFB) wavefunction w/ determinants of different N_e
- Has correct number of electrons on average $\langle N \rangle = N_e$
- Has correct number of correlated electron pairs
- Recent papers:
 - I. J. Chem. Phys. 131, 121102 (2009)
 - II J. Chem. Phys. 131, 164119 (2009)
 - III. J. Chem. Phys. 132, 024111(2010)

Constrained-Pairing Mean-Field Theory

We **constrain** pairings to the entangled region of the orbital space:



 H_2 cc-pV5Z



N₂ 6-311++G**



Dynamical Correlation

- CPMFT orbitals and density are symmetryadapted
- How do we add dynamical correlation only?
- <u>Answer</u>: use "alternative" densities derived from the CPMFT 2pdm ansatz
- Use <u>total</u> and <u>on-top</u> densities as fundamental variables
- Feed alternative densities into regular DFT correlation subroutines

Alternative Densities

We define alternative densities (χ) from total (ρ) and on-top (Γ) densities:

$$\chi_{\alpha}(\mathbf{r}) = \frac{1}{2} \left(\rho(\mathbf{r}) + \sqrt{\rho^2(\mathbf{r}) - 2\Gamma(\mathbf{r})} \right)$$
$$\chi_{\beta}(\mathbf{r}) = \frac{1}{2} \left(\rho(\mathbf{r}) - \sqrt{\rho^2(\mathbf{r}) - 2\Gamma(\mathbf{r})} \right)$$

where $\chi_{\alpha}(\mathbf{r}) + \chi_{\beta}(\mathbf{r}) = \rho(\mathbf{r})$ and $\Gamma(\mathbf{r}) = 2\rho_{\alpha}(\mathbf{r})\rho_{\beta}(\mathbf{r}) - \left(\kappa_{\alpha\beta}^{2}(\mathbf{r}) + \kappa_{\beta\alpha}^{2}(\mathbf{r})\right)$

Example: H₂ molecule

(1) At
$$\mathbf{R}_{e}$$
, $\chi_{\alpha}(\mathbf{r}) = \chi_{\beta}(\mathbf{r}) = \rho_{\alpha}(\mathbf{r}) = \rho_{\beta}(\mathbf{r})$ (RKS density)

(2) At dissociation, $\chi_{\alpha}(\mathbf{r}) = \rho(\mathbf{r})$ and $\chi_{\beta}(\mathbf{r}) = 0$ (UKS-like density)

We change variables of
$$E_{xc}$$
:
 $E_{XC}[\rho_{\alpha}, \rho_{\beta}, \cdots] \rightarrow E_{XC}[\chi_{\alpha}, \chi_{\beta}, \cdots]$

T.Tsuchimochi, G.E.Scuseria, and A.Savin, JCP (2010)

 H_2 cc-pV5Z



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